

Delocalization in the light matter interaction

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Delocalization in the light matter interaction

1. How realistic are UdW detectors as a model for atoms interacting with light?
2. Are there phenomena which the UdW detector model misses?

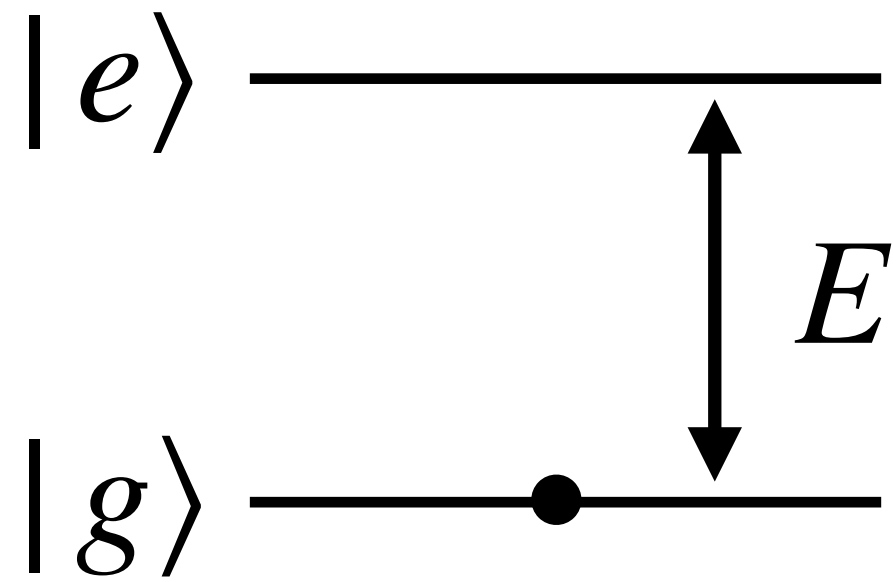
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2. Are there phenomena which the UdW detector model misses?

E.g. spontaneous emission affected by delocalization

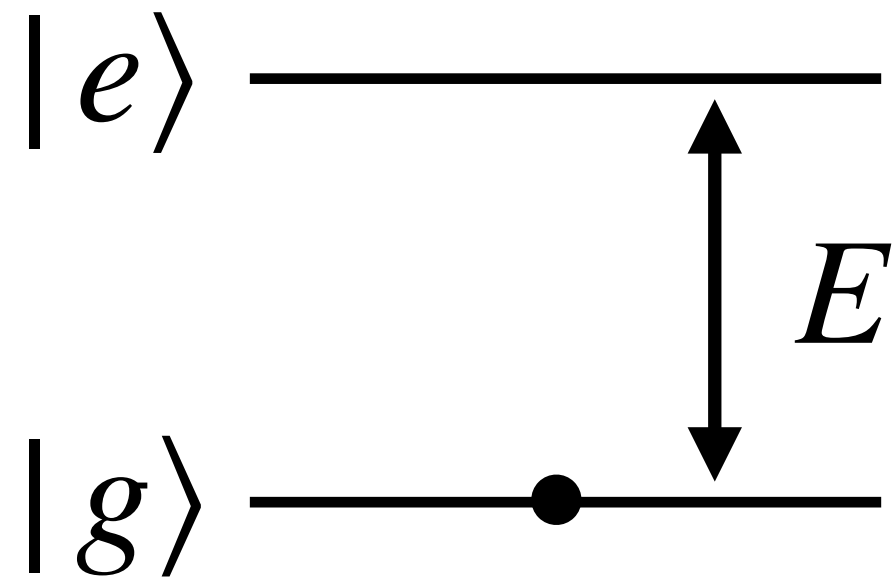
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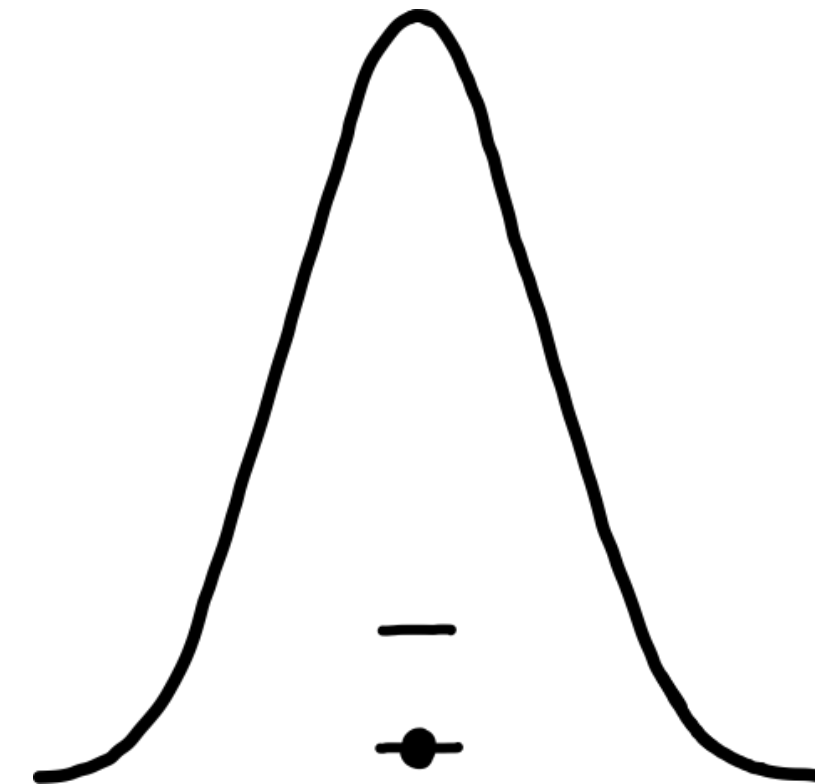


$$\hat{H}^{(internal)} = E |e\rangle\langle e|$$

Incorporate delocalization into detector model?



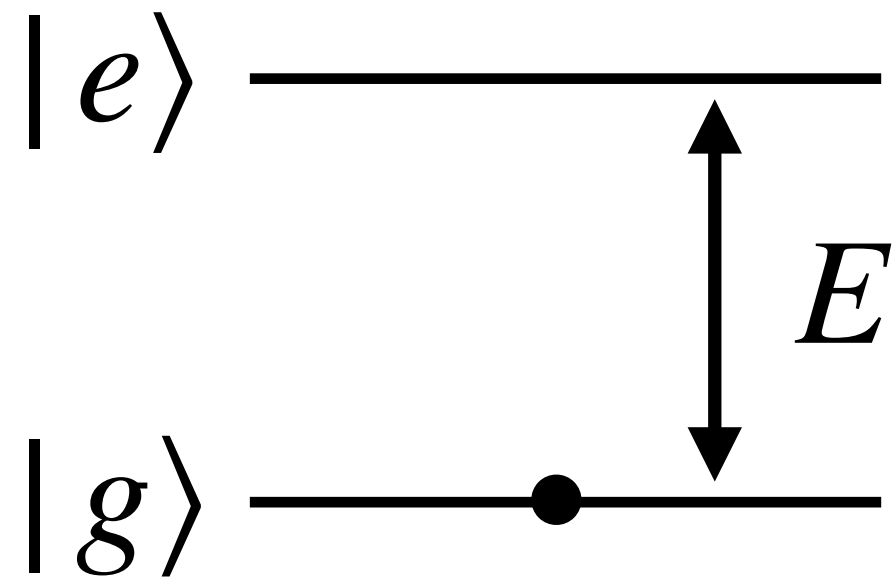
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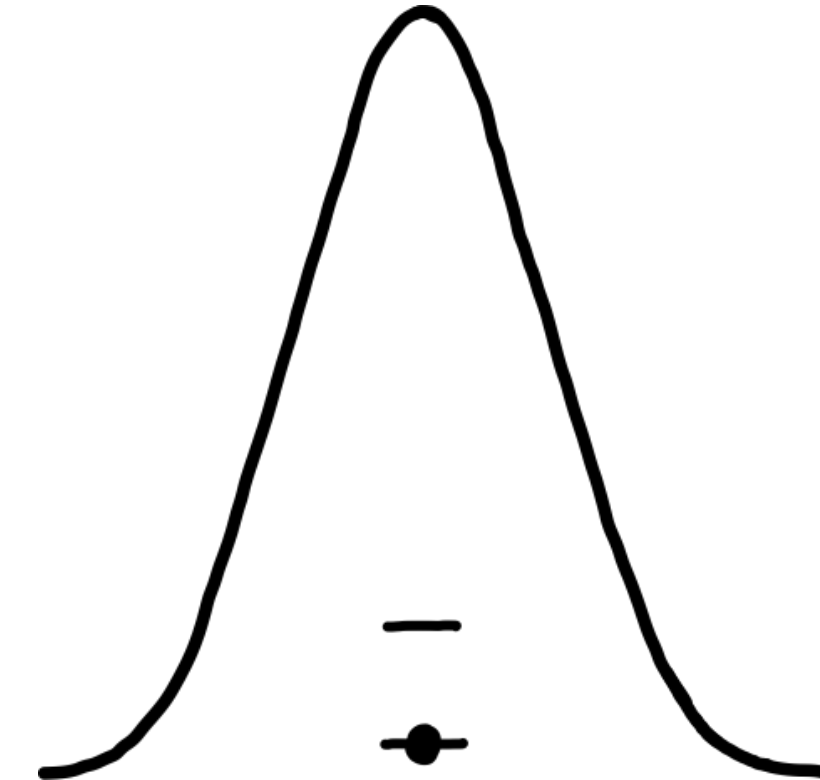
$$\hat{H}^{(external)} = \frac{\hat{\mathbf{p}}^2}{2M}$$

Incorporate delocalization into detector model?

$$\mathcal{H}_{total} = \mathcal{H}_{external} \otimes \mathcal{H}_{internal} \otimes \mathcal{H}_{field}$$



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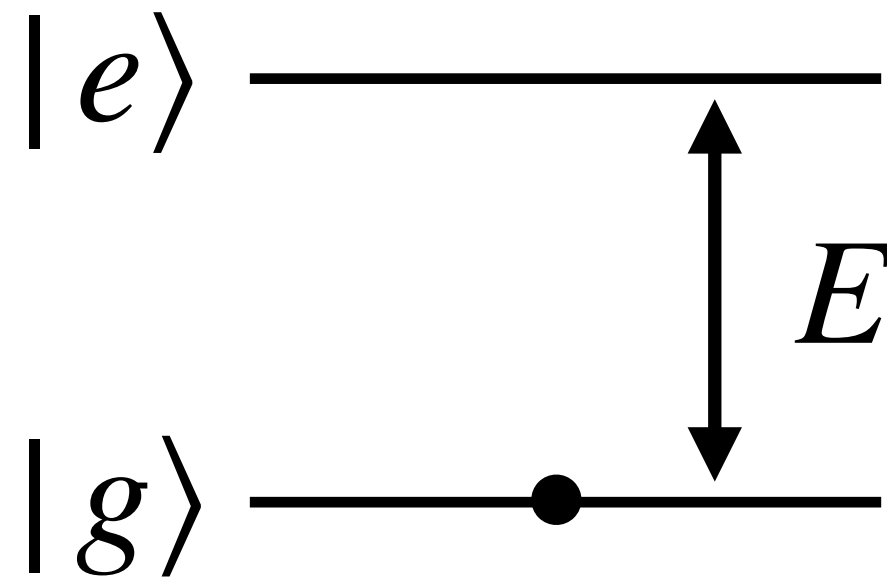


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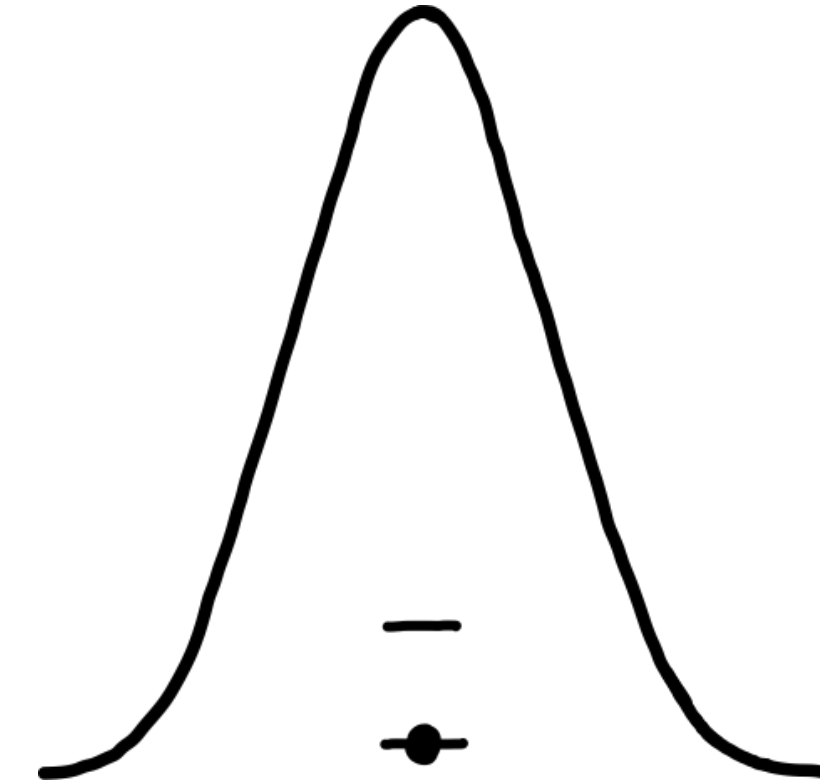
Incorporate delocalization into detector model?

$$\mathcal{H}_{total} = \mathcal{H}_{external} \otimes \mathcal{H}_{internal} \otimes \mathcal{H}_{field}$$

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Incorporate delocalization into detector model?

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Spectral theorem: $\hat{X} |x\rangle = x |x\rangle$, $f(\hat{X}) = \int dx f(x) |x\rangle \langle x|$

For operator valued functions: $\hat{f}(\hat{X}) = \int dx \hat{f}(x) \otimes |x\rangle \langle x|$

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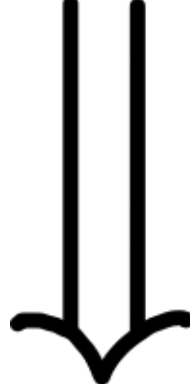
$$\hat{H}^{(int)} = \lambda \int d^3x | \mathbf{x} \rangle \langle \mathbf{x} | \otimes \hat{\mu} \otimes \hat{\phi}(\mathbf{x})$$

More realistic model:

$$\hat{H}^{(total)} = \frac{\left(\hat{\mathbf{p}}_p + e\hat{\mathbf{A}}(\hat{\mathbf{x}}_p)\right)^2}{2m_p} + \frac{\left(\hat{\mathbf{p}}_e - e\hat{\mathbf{A}}(\hat{\mathbf{x}}_e)\right)^2}{2m_e} - \frac{e^2}{4\pi\epsilon_0|\hat{\mathbf{x}}_p - \hat{\mathbf{x}}_e|} + \int \frac{d^3k}{(2\pi)^{3/2}} |\mathbf{k}| \sum_{s=1}^2 a_{\mathbf{k}}^{s\dagger} a_{\mathbf{k}}^s$$

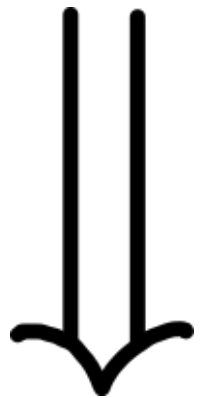
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CM +  relative d.o.f.

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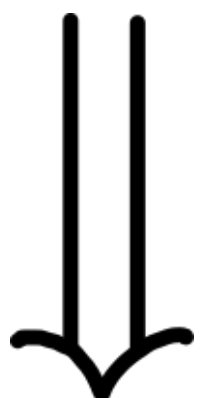
CM +  relative d.o.f.

$$\hat{H}^{(free)} := \frac{\hat{\mathbf{p}}_{CM}^2}{2M} + \frac{\hat{\mathbf{p}}_{rel}^2}{2\mu} - \frac{e^2}{4\pi\epsilon_0|\mathbf{x}_{rel}|} + \int \frac{d^3k}{\sqrt{2\pi}^3} |\mathbf{k}| \sum_{s=1}^2 a_{\mathbf{k}}^{s\dagger} a_{\mathbf{k}}^s$$

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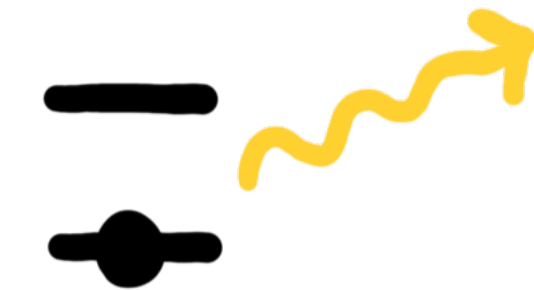
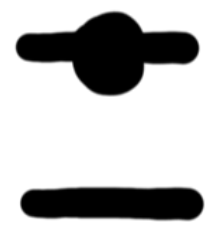
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... qualitatively similar results. Go back to detector model.

Do e.g. the spontaneous emission rates differ?

$$\hat{H}^{(int)} = \lambda \int d^3x \, |\mathbf{x}\rangle\langle\mathbf{x}| \otimes \hat{\mu} \otimes \hat{\phi}(\mathbf{x}) \quad \text{vs.} \quad \hat{H}^{(int)} = \lambda \hat{\mu} \otimes \hat{\phi}(\mathbf{x})$$

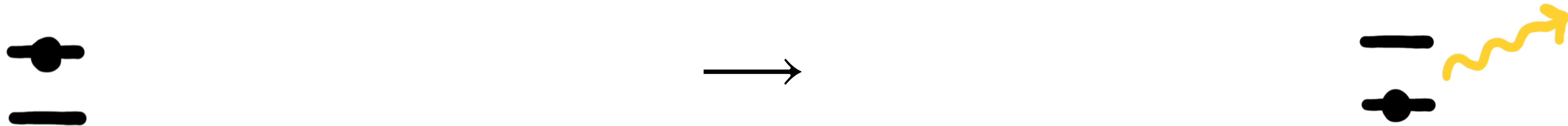
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$$|\Psi_{in}\rangle = |e\rangle \otimes |0\rangle$$

$$|\Psi_{out}\rangle = |g\rangle \otimes |1_{\mathbf{k}}\rangle$$

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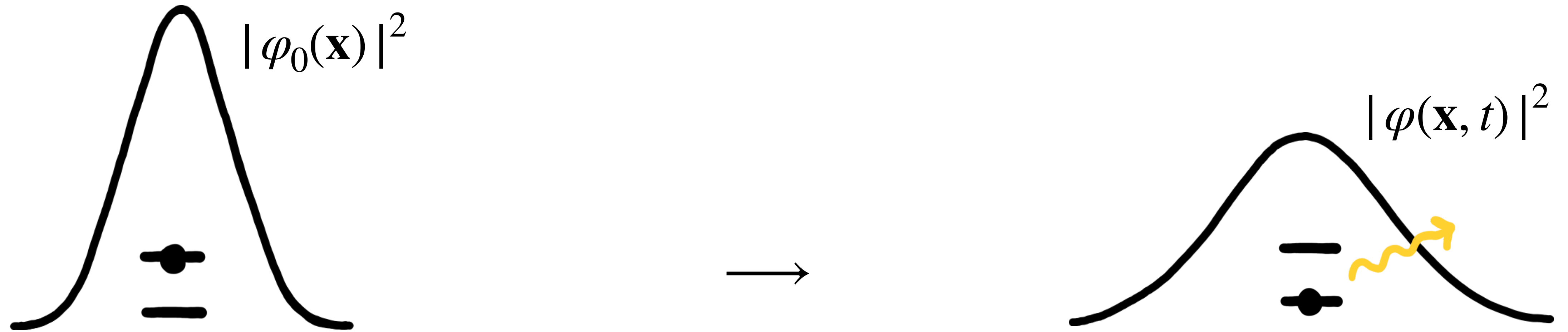


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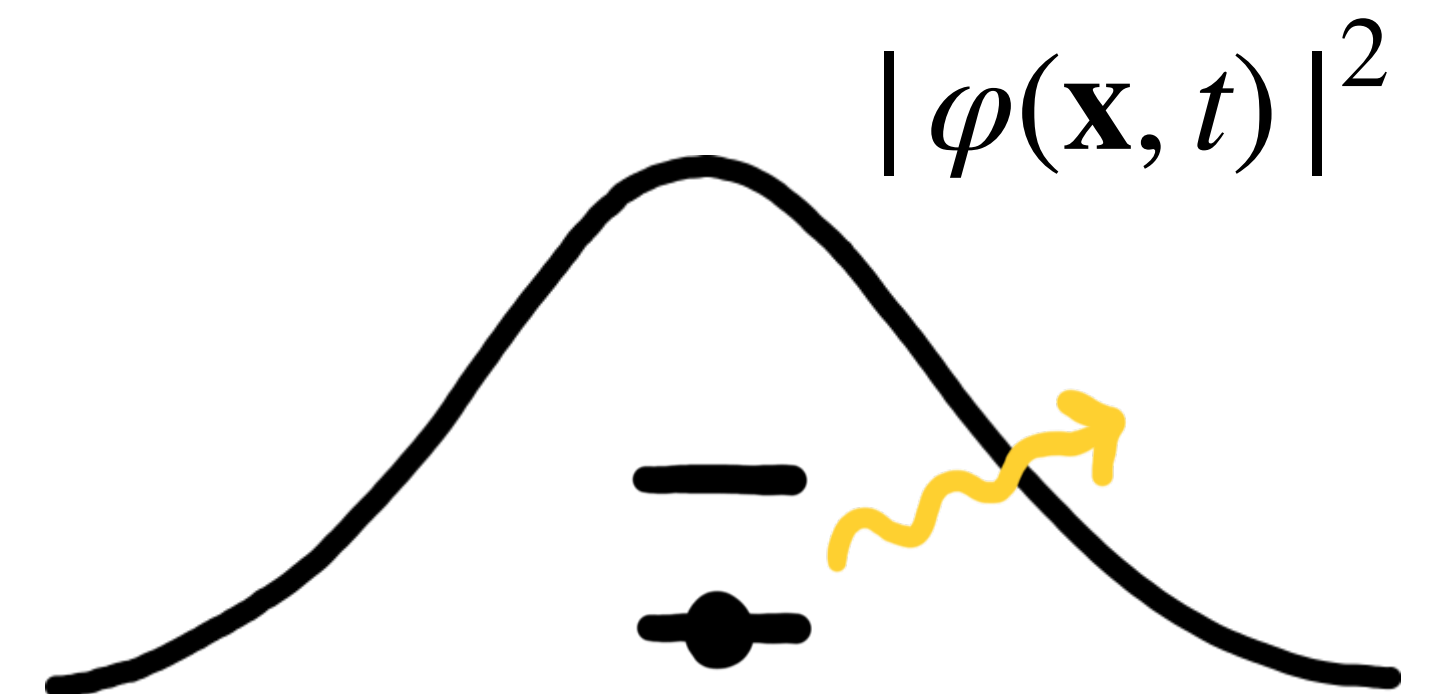
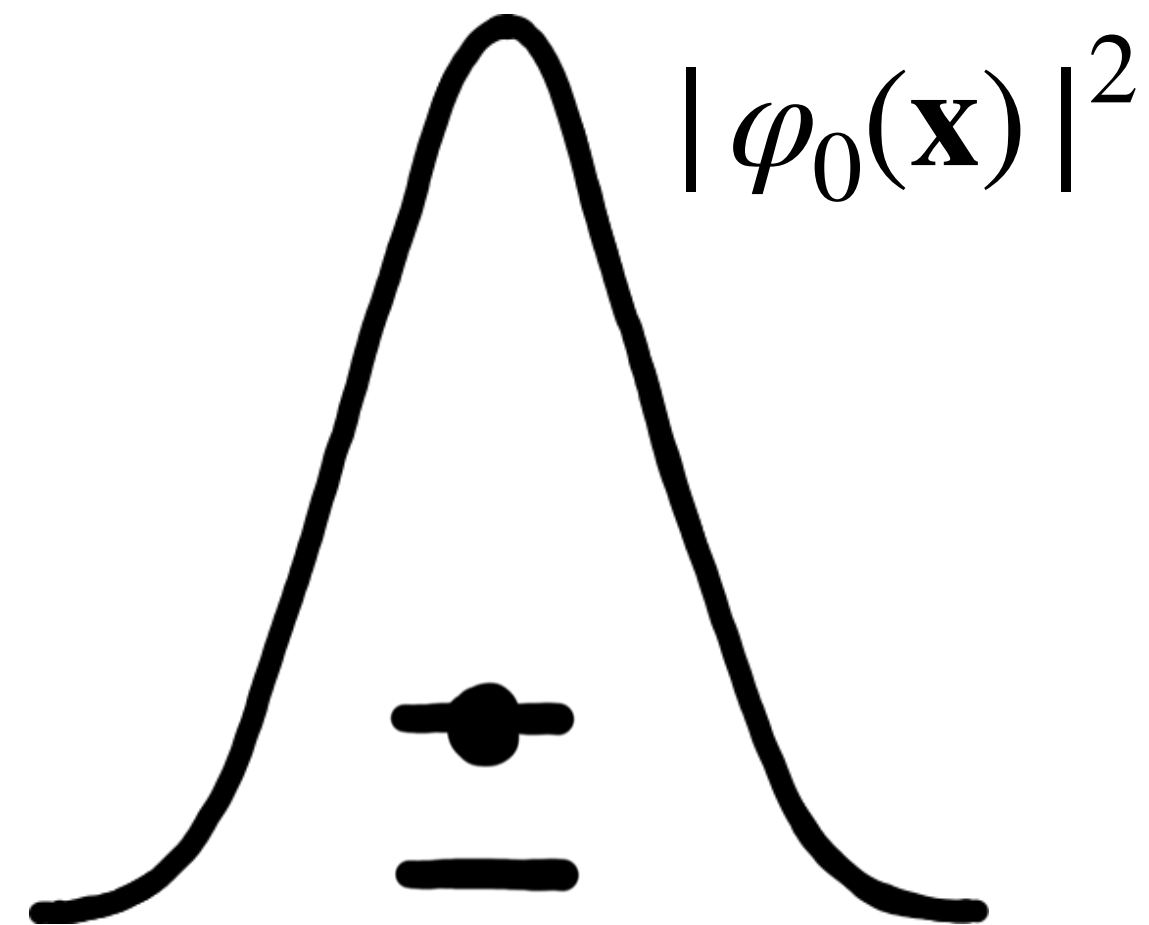
$$|\Psi_{out}\rangle = |g\rangle \otimes |1_{\mathbf{k}}\rangle$$

Spontaneous emission rate (tracing over field) : $\mathcal{R}_{e \rightarrow g} = \frac{\lambda^2 E}{c^3 \hbar^3}$

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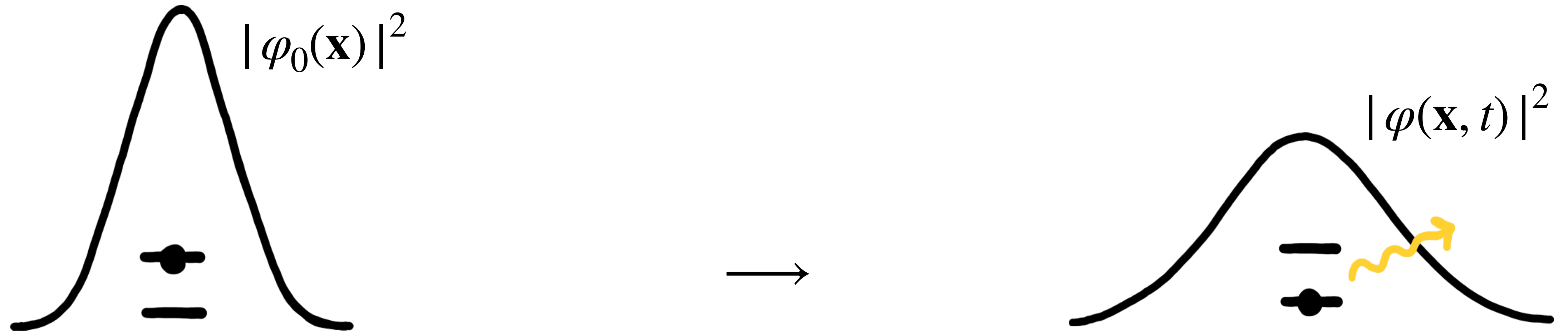
Do e.g. the spontaneous emission rates differ?



$$|\Psi_{in}\rangle = \int d^3p \, \varphi_0(\mathbf{p}) |\mathbf{p}\rangle \otimes |e\rangle \otimes |0\rangle$$

$$|\Psi_{out}\rangle = |\mathbf{p}'\rangle \otimes |g\rangle \otimes |1_{\mathbf{k}}\rangle$$

Do e.g. the spontaneous emission rates differ?



Spontaneous emission rate (tracing over field and recoil momentum):

$$\mathcal{R}_{e \rightarrow g} = \frac{\lambda^2}{2c\hbar^3} \int d^3p |\varphi_0(\mathbf{p})|^2 \left(2M - \frac{M^2}{|\mathbf{p}|} \sqrt{\left(\frac{|\mathbf{p}|}{M} + c \right)^2 + \frac{2E}{M}} + \frac{M^2}{|\mathbf{p}|} \sqrt{\left(\frac{|\mathbf{p}|}{M} - c \right)^2 + \frac{2E}{M}} \right)$$

Delocalization affects spontaneous emission!

$$\mathcal{R}_{e \rightarrow g} = \underbrace{\frac{\lambda^2}{2c\hbar^3} \int d^3p |\varphi_0(\mathbf{p})|^2}_{\text{initial CM wave function}} \underbrace{\left(2M - \frac{M^2}{|\mathbf{p}|} \sqrt{\left(\frac{|\mathbf{p}|}{M} + c \right)^2 + \frac{2E}{M}} + \frac{M^2}{|\mathbf{p}|} \sqrt{\left(\frac{|\mathbf{p}|}{M} - c \right)^2 + \frac{2E}{M}} \right)}_{\text{template function}}$$

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Limit of infinite detector mass:

$$\lim_{M \rightarrow \infty} \mathcal{R}_{e \rightarrow g} = \frac{\lambda^2 E}{c^3 \hbar^3}$$

= decay rate for standard UdW detector!

$$\mathcal{R}_{e \rightarrow g} = \underbrace{\frac{\lambda^2}{2c\hbar^3} \int d^3p \, |\varphi_0(\mathbf{p})|^2}_{\text{assume } |\mathbf{p}| \ll Mc} \left(2M - \frac{M^2}{|\mathbf{p}|} \sqrt{\left(\frac{|\mathbf{p}|}{M} + c\right)^2 + \frac{2E}{M}} + \frac{M^2}{|\mathbf{p}|} \sqrt{\left(\frac{|\mathbf{p}|}{M} - c\right)^2 + \frac{2E}{M}} \right)$$

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assume $|\mathbf{p}| \ll Mc$

$$= 2A_1 + 2A_2 |\mathbf{p}|^2 + \mathcal{O}(|\mathbf{p}|^4)$$

$$A_1 := 1 - \frac{c\sqrt{M}}{\sqrt{c^2M + 2E}}$$

$$A_2 := \frac{Ec}{\sqrt{M}} \frac{1}{(c^2M + 2E)^{5/2}}$$

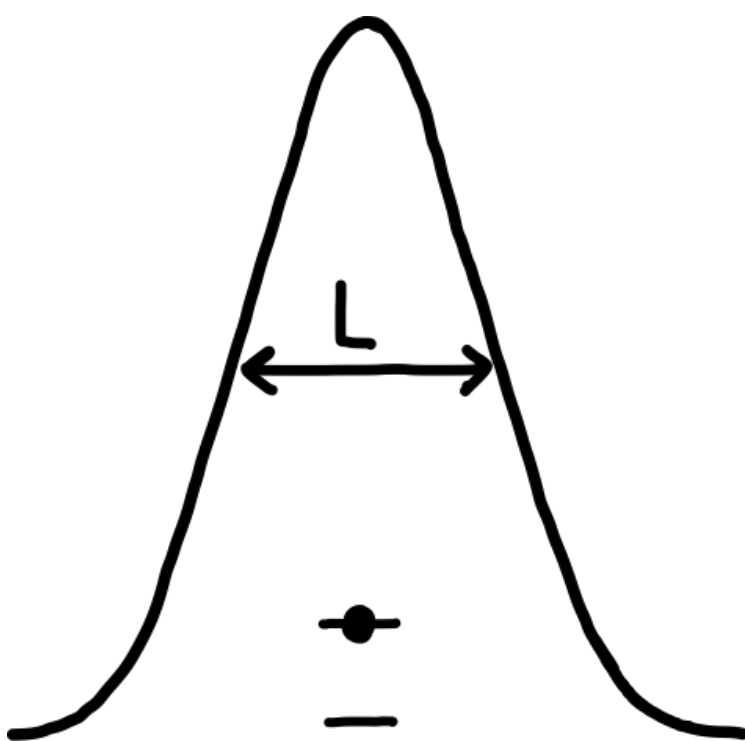
Spontaneous emission rate (with delocalization):

$$\mathcal{R}_{e \rightarrow g} = \frac{\lambda^2}{c \hbar^3} \int d^3 p \, |\varphi_0(\mathbf{p})|^2 \left(A_1 + A_2 |\mathbf{p}|^2 + \mathcal{O}(|\mathbf{p}|^4) \right)$$

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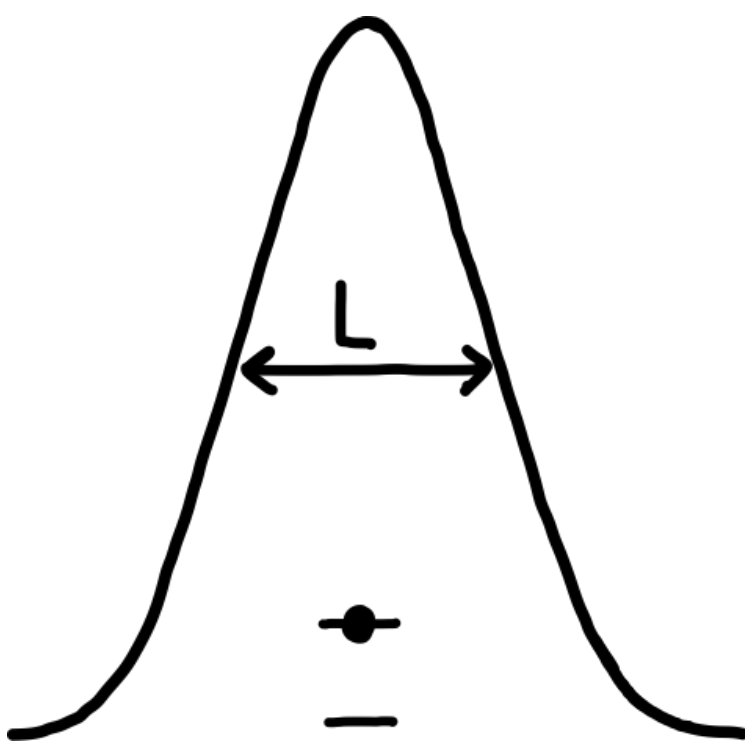
Gaussian wave packet state:

$$\varphi_0(\mathbf{p}) \sim e^{-\frac{|\mathbf{p}|^2 L^2}{4}}$$


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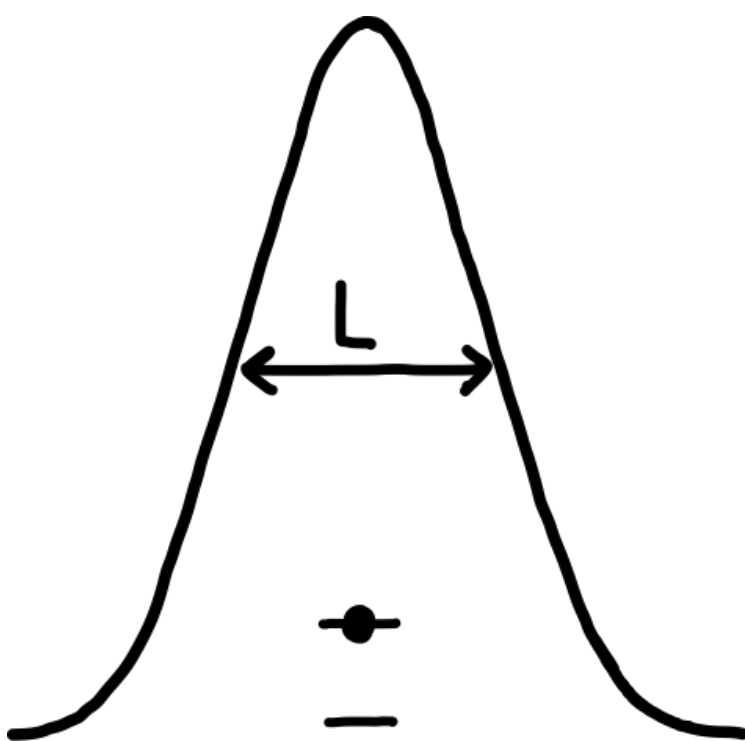
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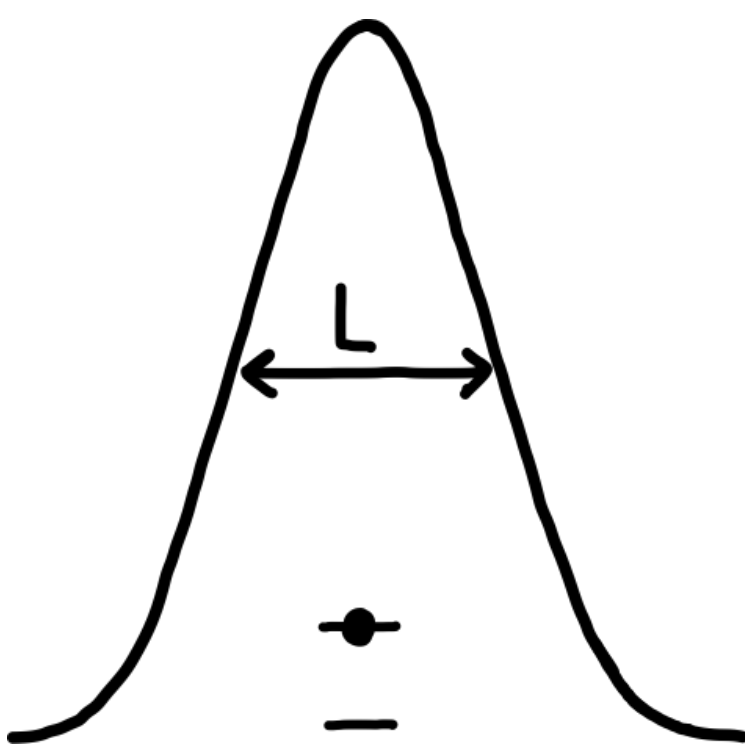
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Faster delocalization
 $\Rightarrow$ Faster decay

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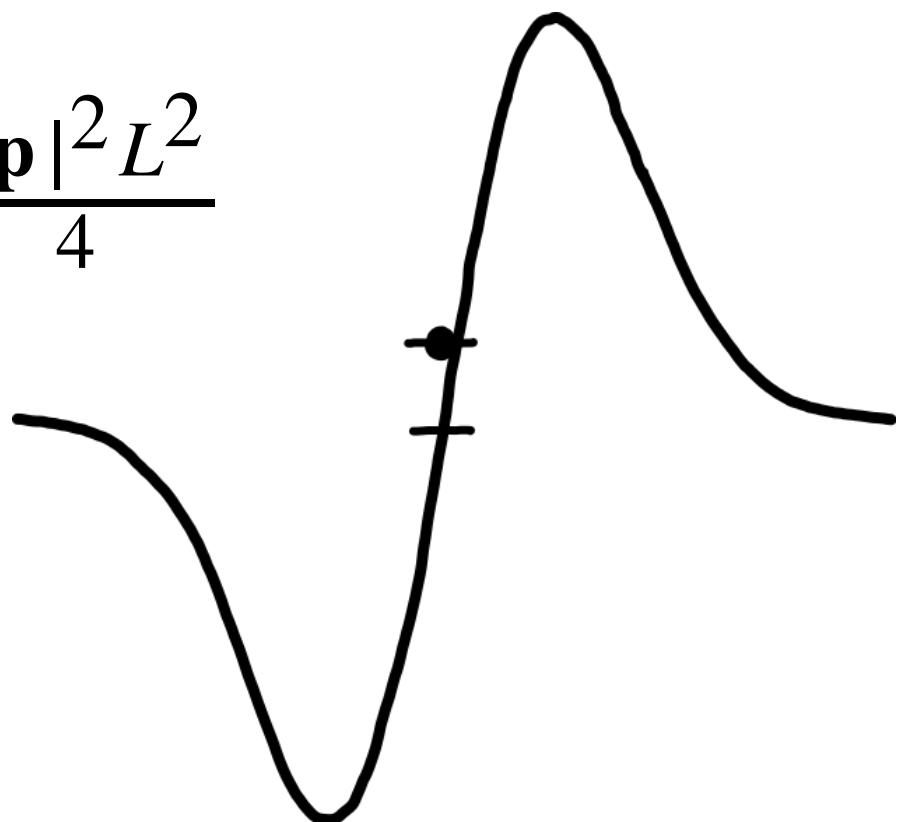
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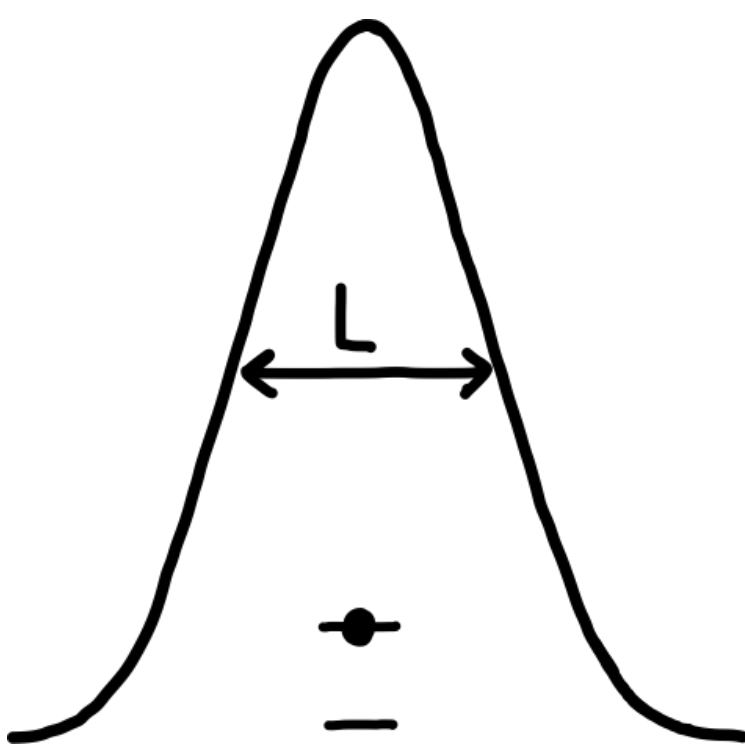
Excited state (e.g. ion trap):

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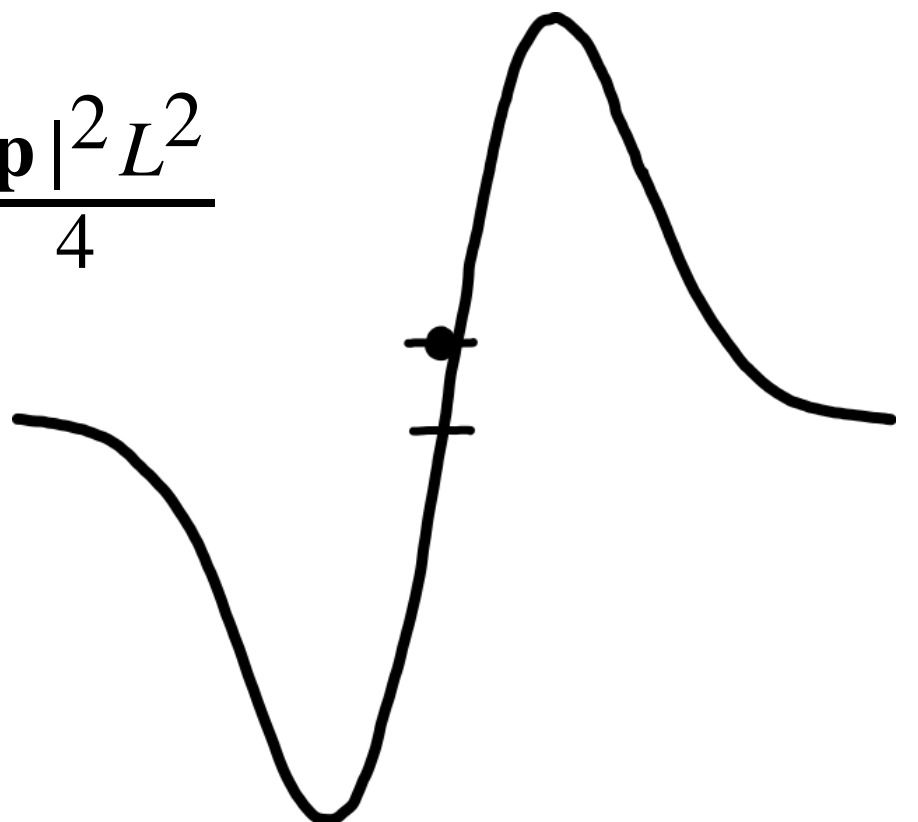
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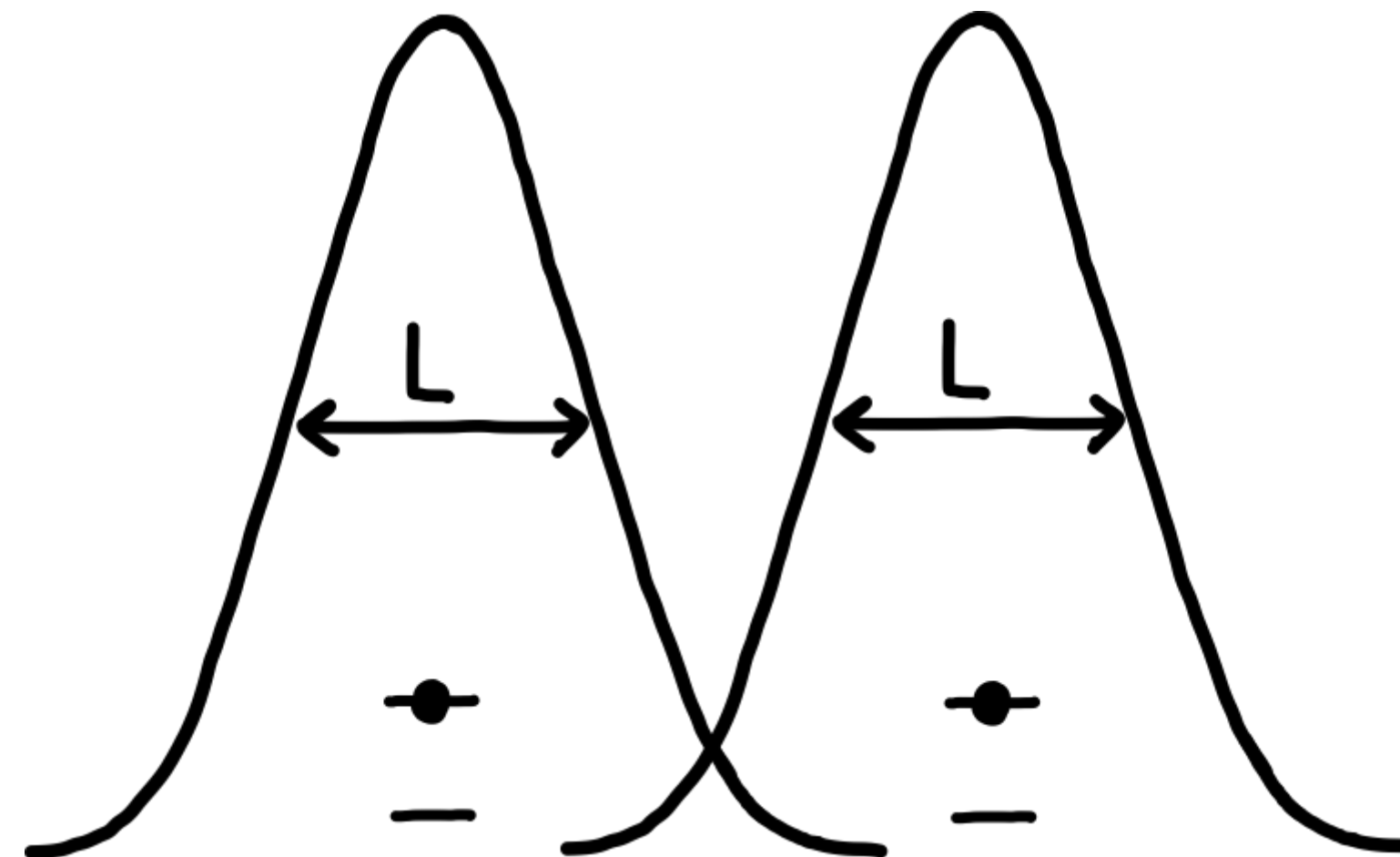
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Spontaneous emission affected by delocalization process

Decay rates affected differently by incoherent vs. coherent delocalization?

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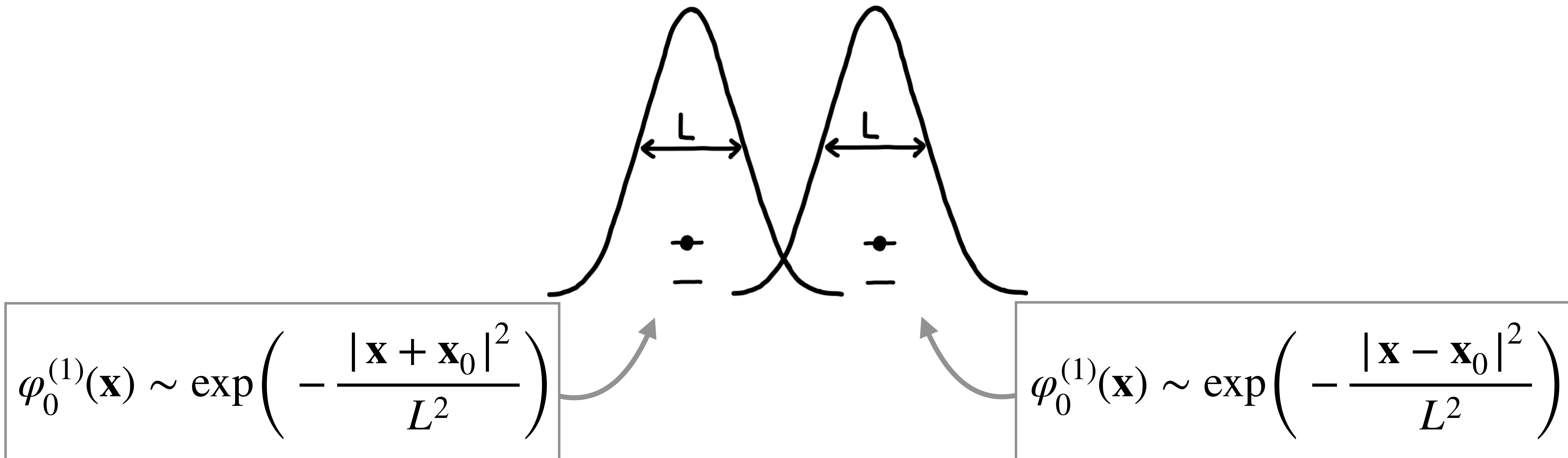
$$\varphi_0^{(1)}(\mathbf{x}) \sim \exp\left(-\frac{|\mathbf{x} + \mathbf{x}_0|^2}{L^2}\right)$$

$$\varphi_0^{(1)}(\mathbf{x}) \sim \exp\left(-\frac{|\mathbf{x} - \mathbf{x}_0|^2}{L^2}\right)$$

Decay rates affected differently by incoherent vs. coherent delocalization?

Incoherent superposition

$$\rho_0 = \frac{1}{2} (|\varphi_0^{(1)}\rangle\langle\varphi_0^{(1)}| + |\varphi_0^{(2)}\rangle\langle\varphi_0^{(2)}|)$$



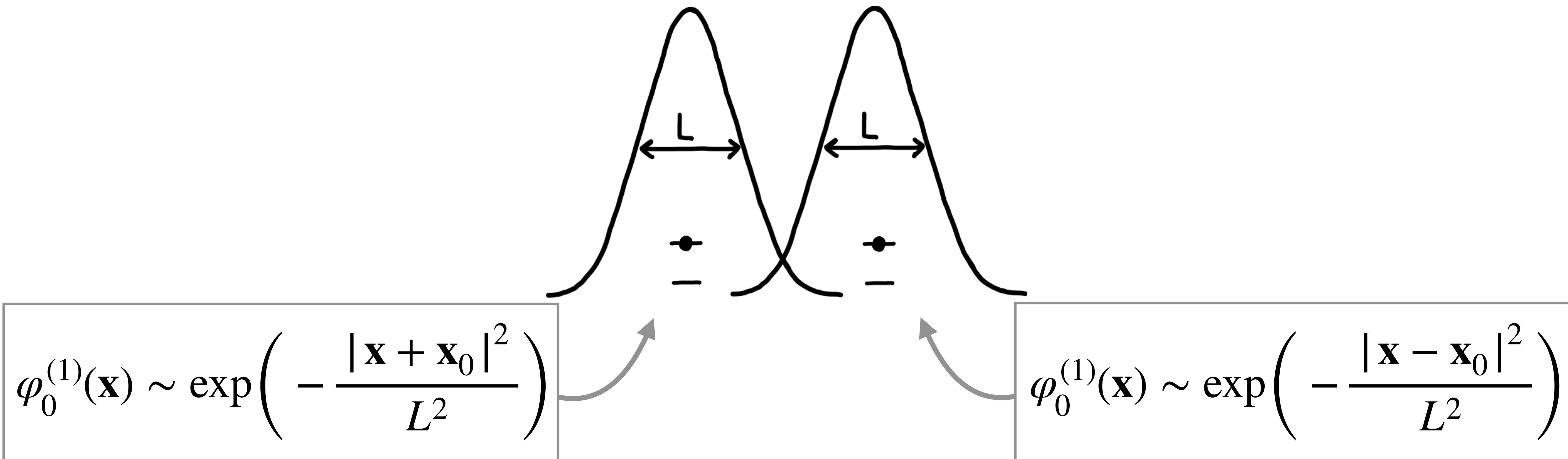
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Coherent superposition

$$\varphi_0(\mathbf{x}) \sim \varphi_0^{(1)}(\mathbf{x}) + \varphi_0^{(2)}(\mathbf{x})$$



Decay rates affected differently by incoherent vs. coherent delocalization?

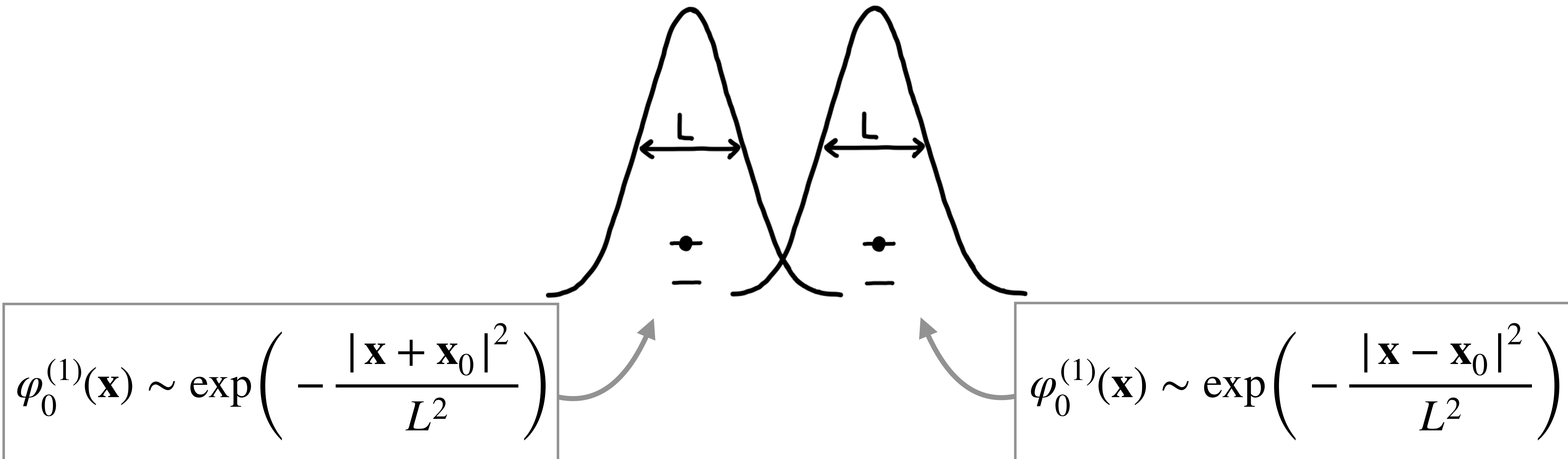
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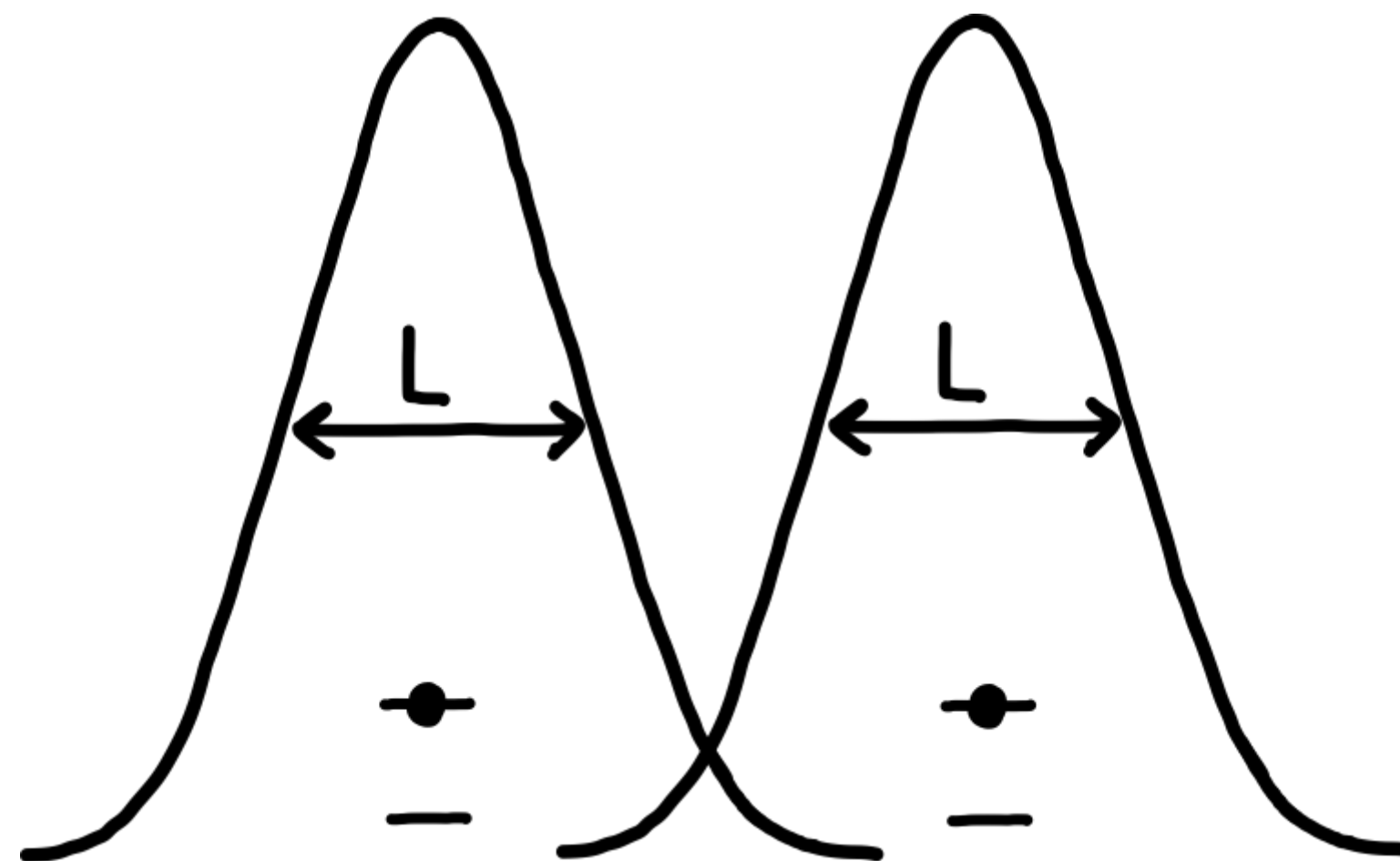
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$$\mathcal{R}_{e \rightarrow g} \approx \frac{\lambda^2 M}{c\hbar^3} \left(A_1 + \frac{3}{L^2} A_2 \left(1 - \frac{4|\mathbf{x}_0|^2}{3L^2} \frac{e^{-2|\mathbf{x}_0|^2/L^2}}{1 + e^{-2|\mathbf{x}_0|^2/L^2}} \right) \right)$$



$$\varphi_0^{(1)}(\mathbf{x}) \sim \exp\left(-\frac{|\mathbf{x} + \mathbf{x}_0|^2}{L^2}\right)$$

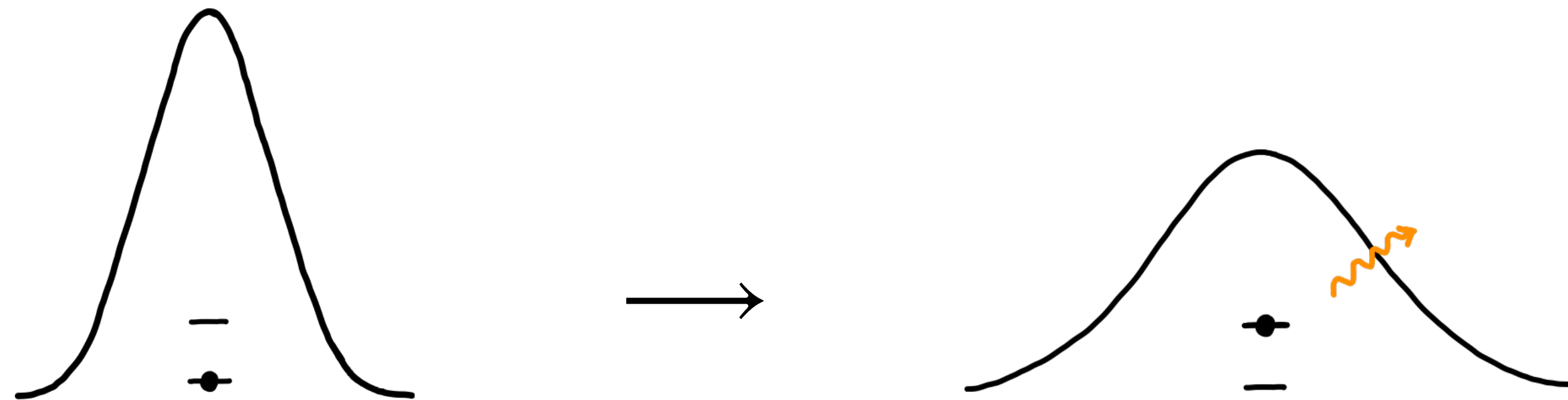
$$\varphi_0^{(2)}(\mathbf{x}) \sim \exp\left(-\frac{|\mathbf{x} - \mathbf{x}_0|^2}{L^2}\right)$$

**Decay rates affected differently by
incoherent vs. coherent delocalization!**

Can delocalization trigger excitation of atoms?

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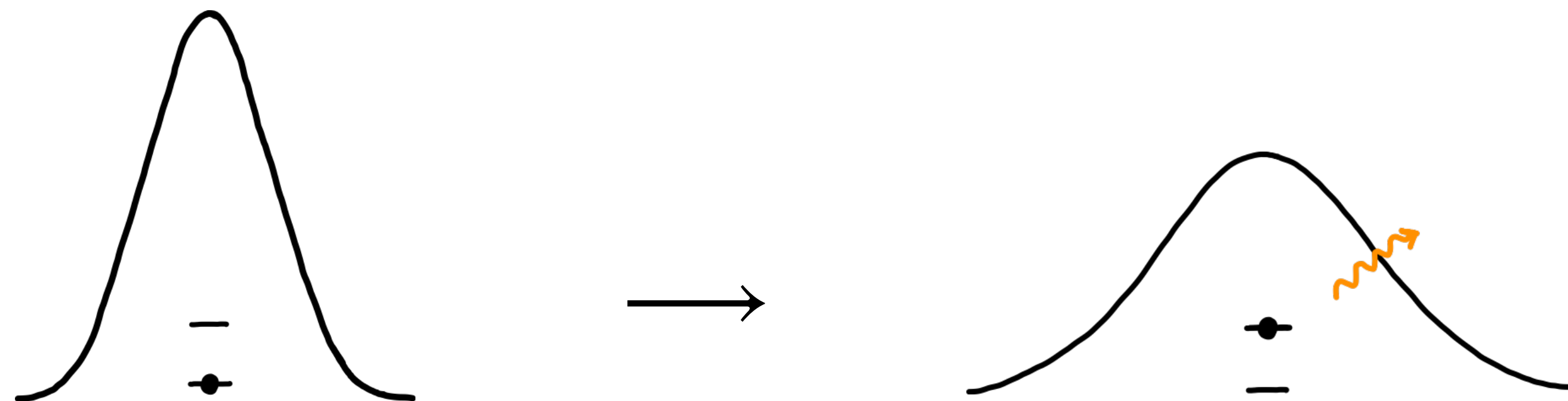
$$|\Psi_{in}\rangle = \int d^3p \varphi_0(\mathbf{p}) |\mathbf{p}\rangle \otimes |g\rangle \otimes |0\rangle \longrightarrow |\Psi_{out}\rangle = |\mathbf{p}'\rangle \otimes |e\rangle \otimes |1_{\mathbf{k}}\rangle$$



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... delocalization implies kinetic energy uncertainty!

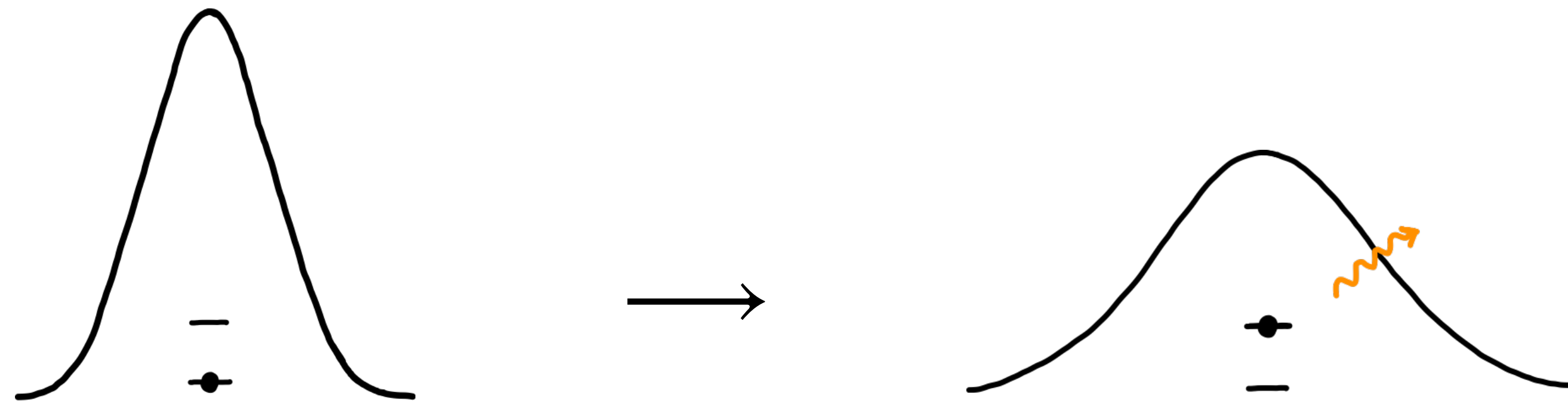
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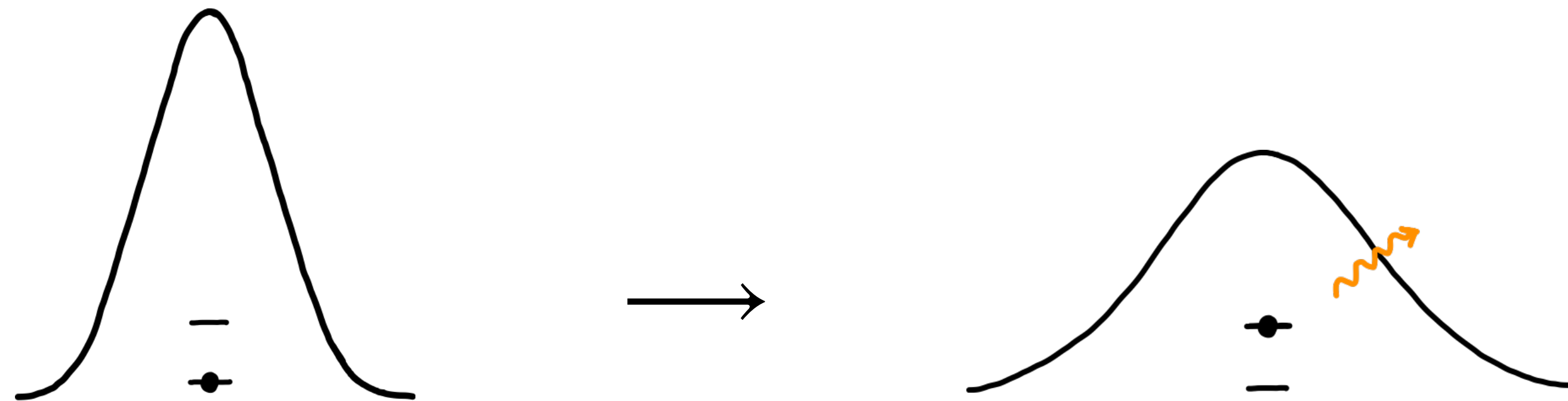
$$\mathcal{R}_{g \rightarrow e} = \frac{\lambda^2}{2c\hbar^2} \underbrace{\int d^3p |\varphi_0(\mathbf{p})|^2}_{\text{initial CM wave function}} \int_0^\infty dk \int_{-1}^1 dz \frac{M}{|\mathbf{p}|} \delta \left(z - \frac{M}{|\mathbf{p}|k} \left(\frac{\hbar k^2}{2M} + \frac{E}{\hbar} + ck \right) \right)$$

initial CM wave function



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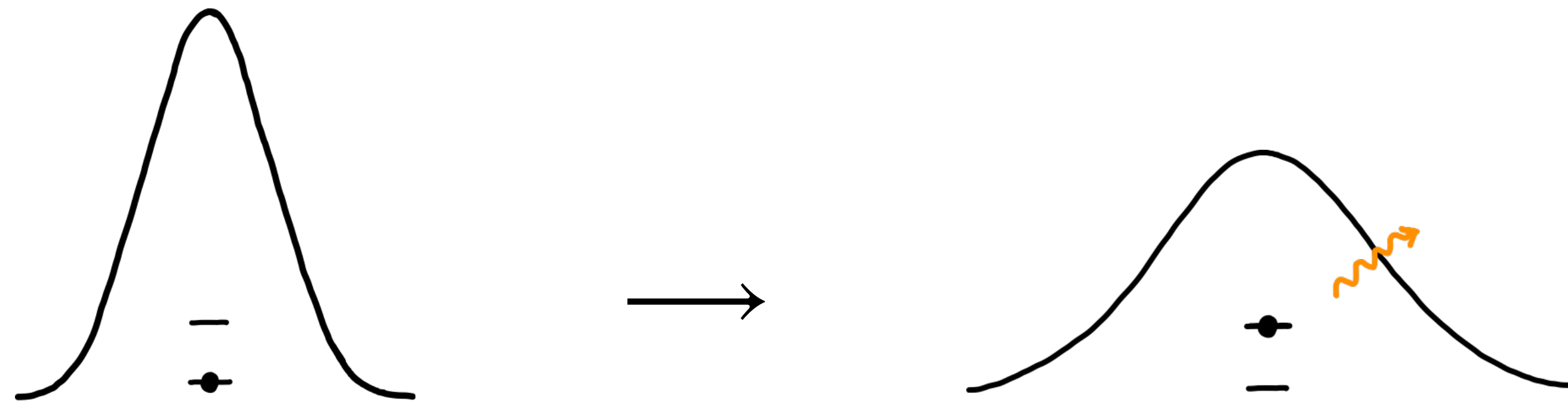
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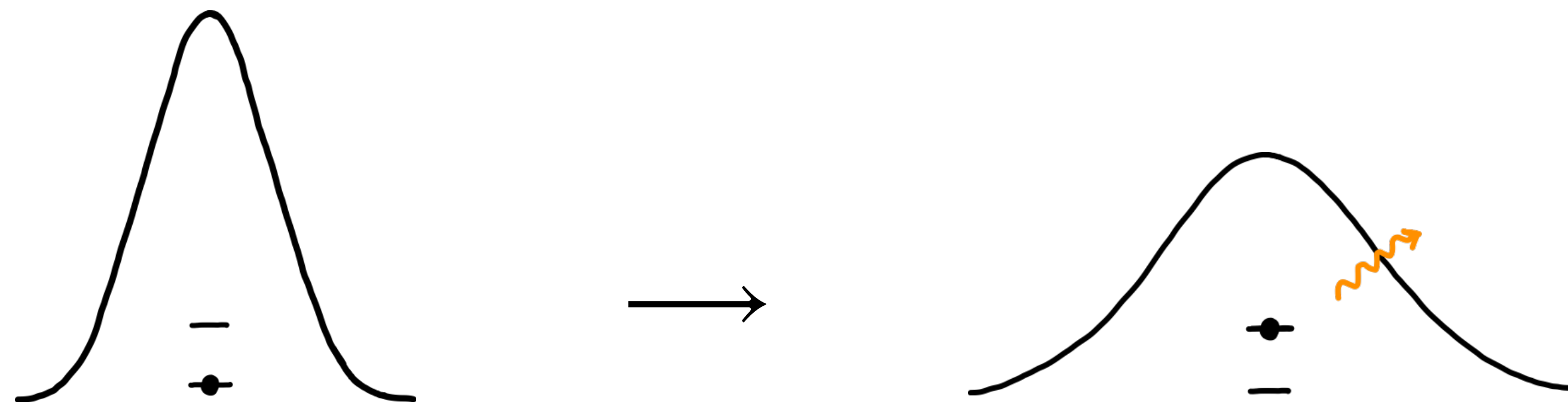
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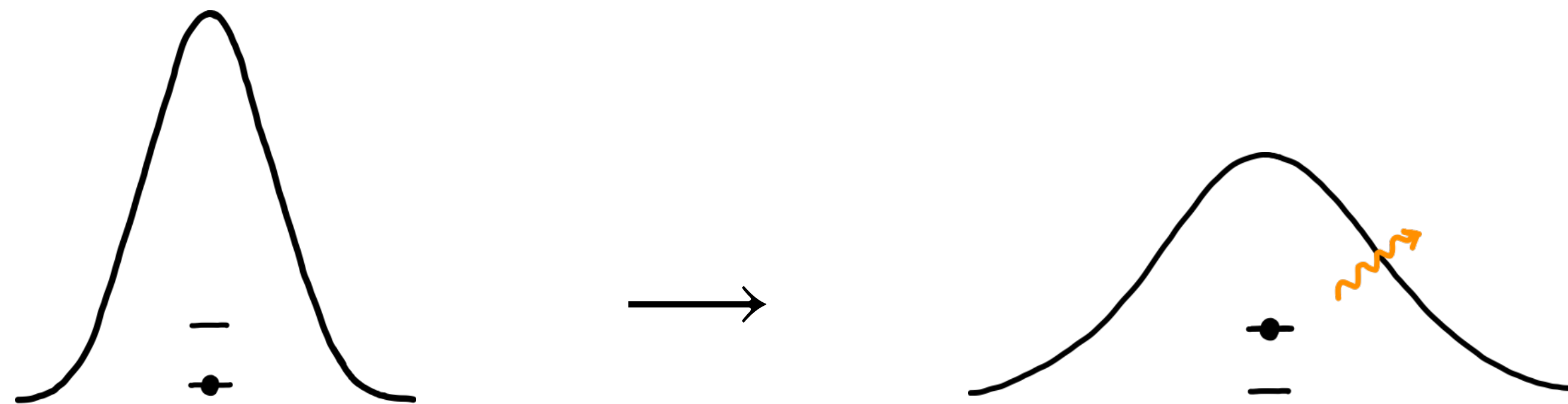
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- Virtual-motion-induced Cherenkov-like radiation

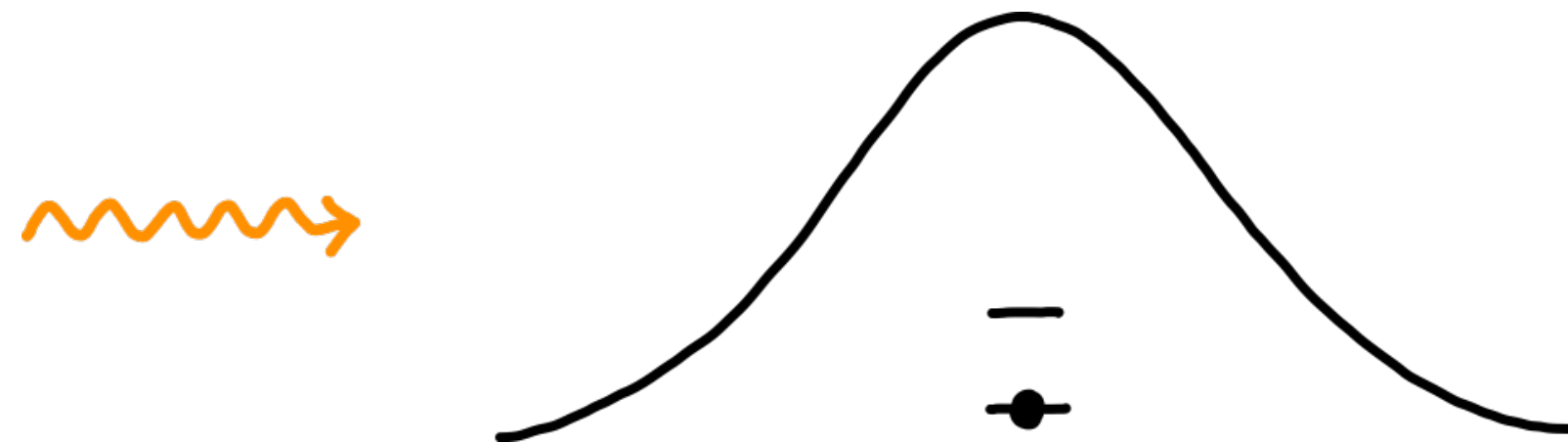


**Delocalization can trigger excitation of atoms
in a medium!**

Quantum channel capacities of the light matter interaction

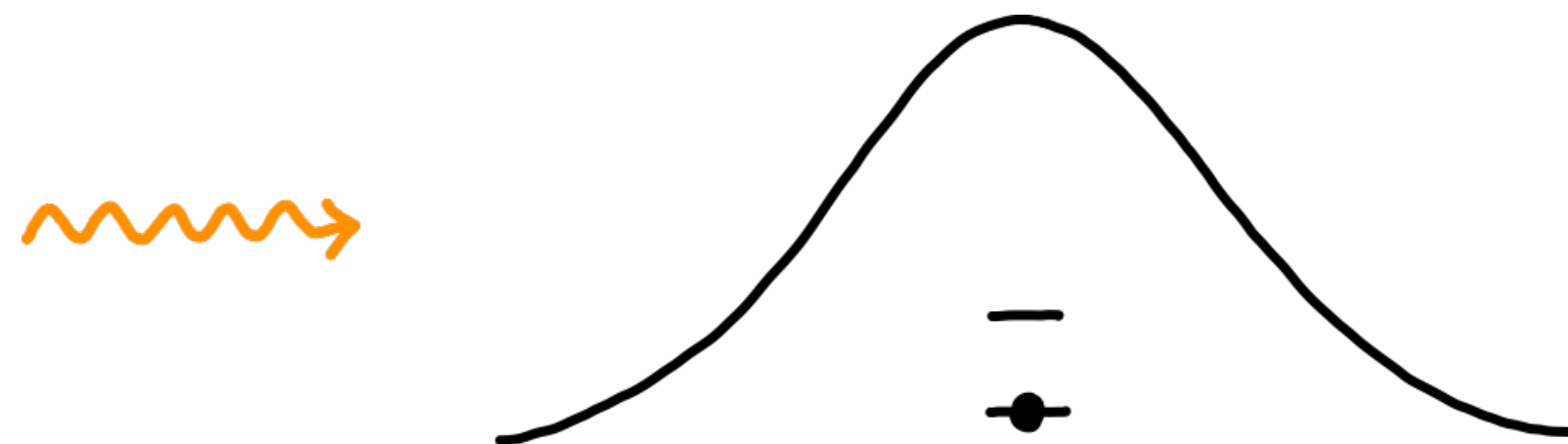
Quantum channel capacities of the light matter interaction

- Assume input (e.g. to-be-absorbed photon) entangled with ancilla



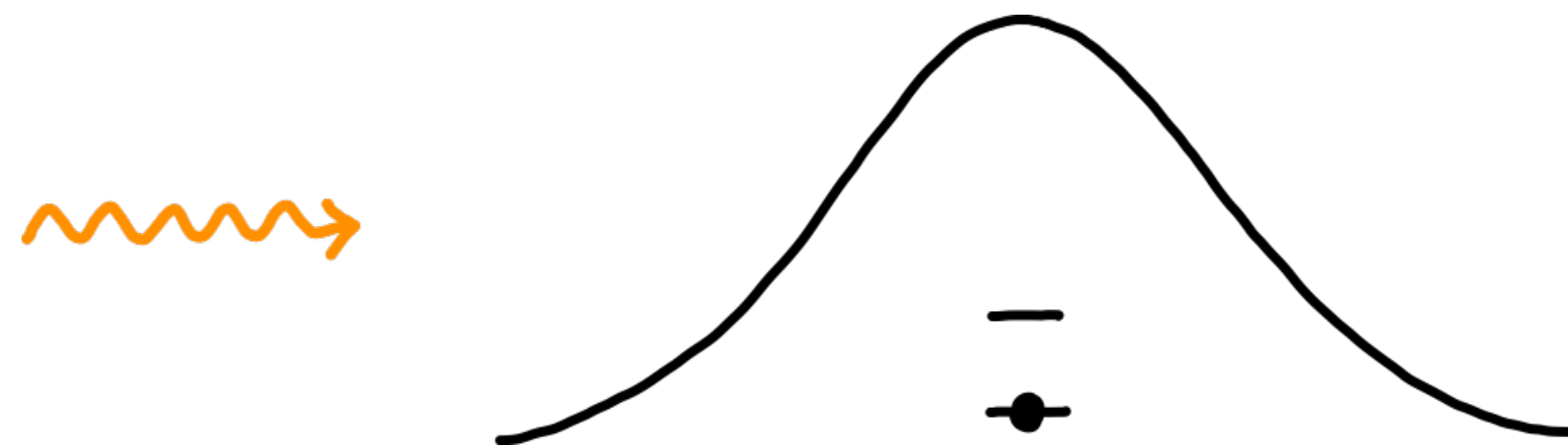
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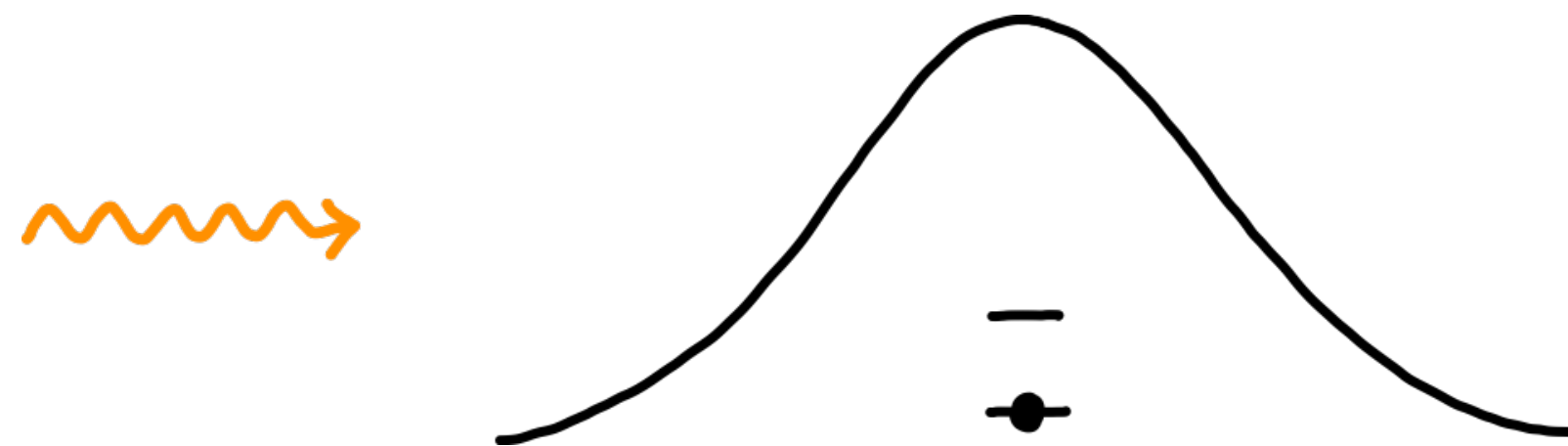
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E.g. for ion trap based quantum computing

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 - Atoms decay the faster, the more sharply they are initially localized
 - Coherent delocalization different from incoherent delocalization
 - predicts new Cherenkov-like effect
- Prepared to study quantum channel capacities of the light matter interaction



Thank you! Questions?

