

Hidden Teleportation Power for Entangled Quantum States

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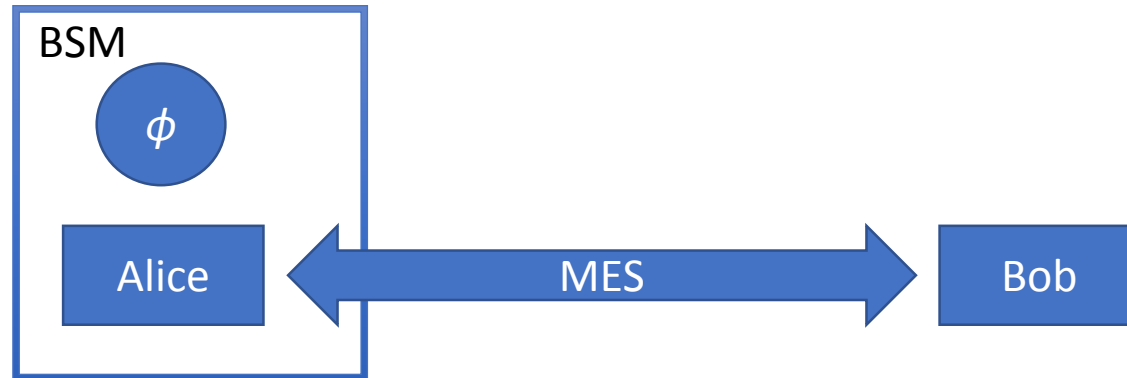
Introduction

- Quantum Teleportation [1]:



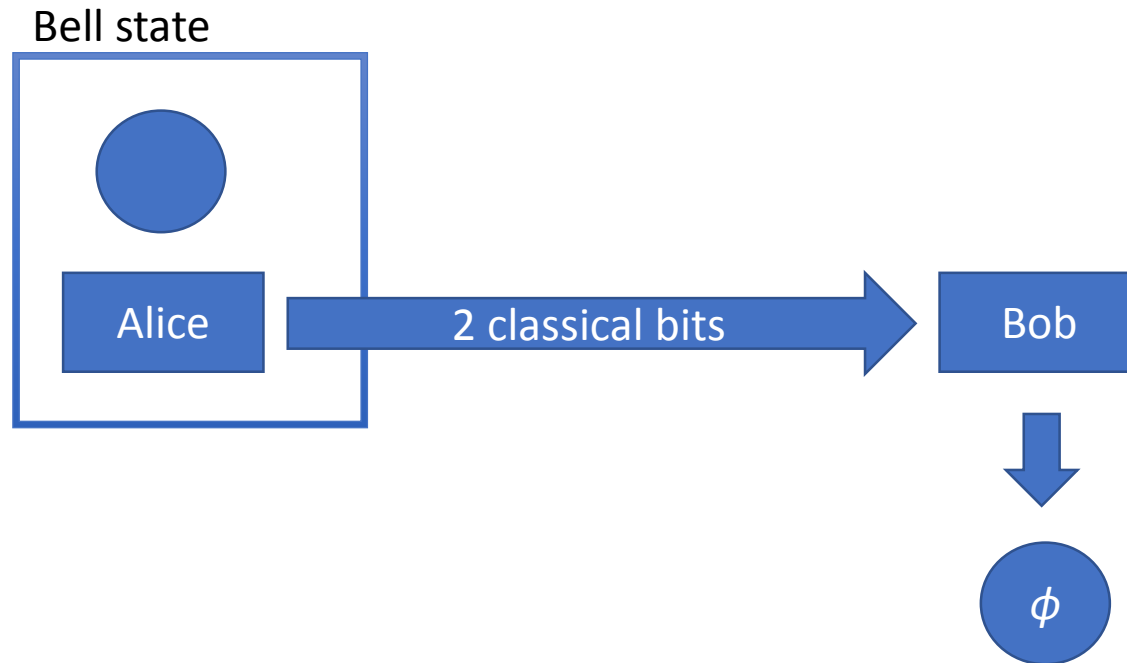
Introduction

- Quantum Teleportation [1]:



Introduction

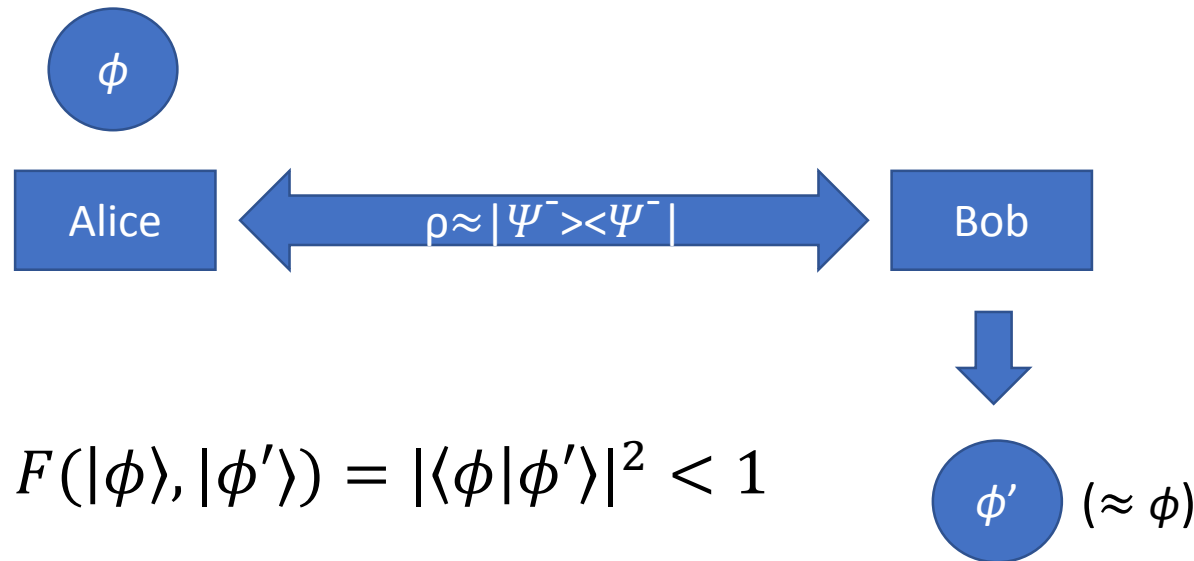
- Quantum Teleportation [1]:



[1] Bennett *et. al.*, Phys. Rev. Lett, **70** (1993)

Introduction

- What if they do not share a maximally entangled state?



- The fidelity $F(|\phi\rangle, |\phi'\rangle) = |\langle\phi|\phi'\rangle|^2 < 1$

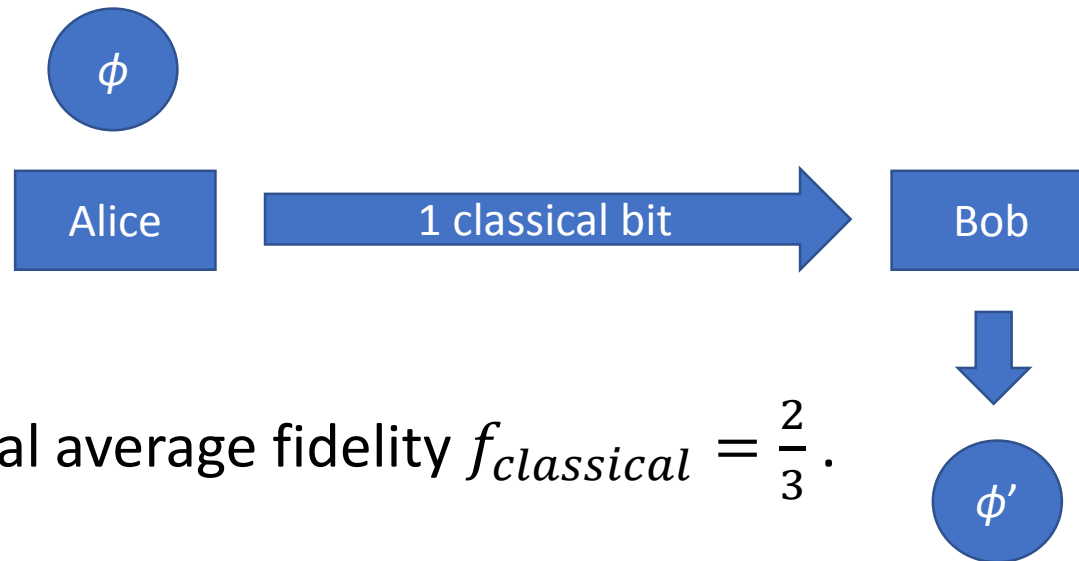
Introduction

- For qubit case:
If Alice and Bob can only have **classical** resource:



Introduction

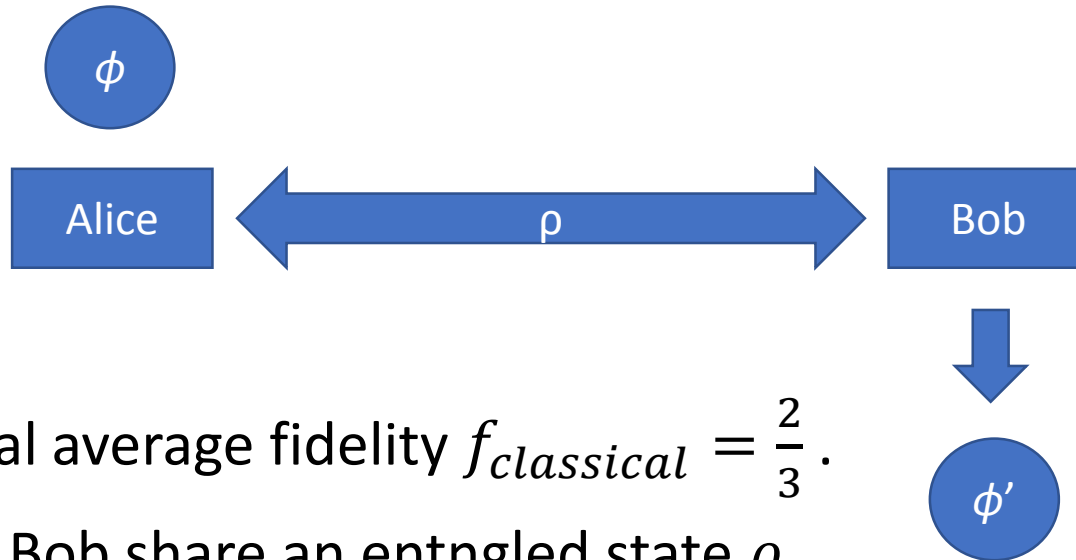
- For qubit case:
If Alice and Bob can only have **classical** resource:



- The maximal average fidelity $f_{classical} = \frac{2}{3}$.

Introduction

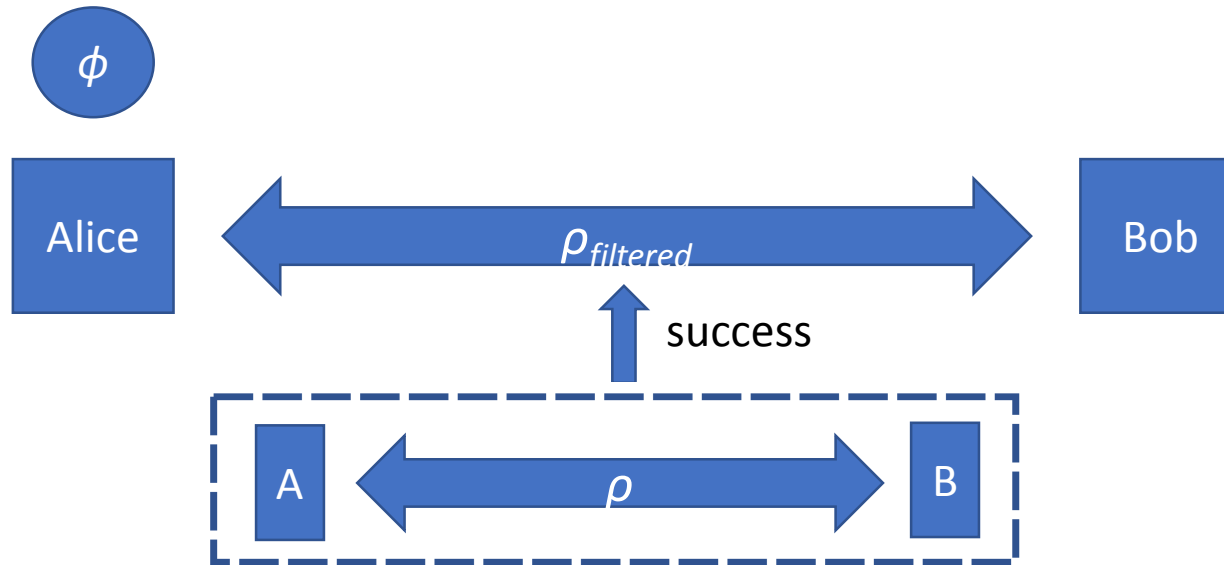
- For qubit case:
If Alice and Bob can only have **classical** resource:



- The maximal average fidelity $f_{classical} = \frac{2}{3}$.
- If Alice and Bob share an entangled state ρ and $f_{ave}^{max}(\rho) > \frac{2}{3}$, we say this entangled state ρ is **useful** for teleportation.

Hidden teleportation power

- Local filtering operation:



- For qubit case, if an entangled state is useless $f_{ave}^{max}(\rho) < \frac{2}{3}$ initially, but become useful $f_{ave}^{max}(\rho_{filtered}) > \frac{2}{3}$ after **successful local filtering**, then we say this entangled state ρ has **hidden teleportation power**.

Hidden teleportation power for qudit states

- The hidden teleportation power for qubit case has been studied. [2], [3]

- Relating **fidelity** and **singlet fraction** [4]:

Suppose that Alice and Bob share an entangled state ρ acting on $H_A \otimes H_B = \mathbb{C}^d \otimes \mathbb{C}^d$, then

$$f_{ave}^{max}(\rho) = \frac{SF(\rho)d+1}{d+1}$$

where $SF(\rho) = \max_{\Psi} \langle \Psi | \rho | \Psi \rangle$ is the maximal singlet fraction of the state ρ .

- $f_{classical} = \frac{2}{d+1}$

An entangled state is **useful** if $f_{ave}^{max}(\rho) > \frac{2}{d+1}$. (or $SF(\rho) > \frac{1}{d}$)

[4] Horodecki *et. al.*, Phys. Rev. A, **60** (1999)

[2] Horodecki *et. al.*, Phys. Rev. Lett, **78** (1997) [3] F. Verstraete, H. Verschelde, Phys. Rev. Lett, **90** (2003)

Hidden teleportation power for qudit states

- Consider following entangled state share between Alice and Bob:

$$\rho(q) = q|\Phi^+\rangle\langle\Phi^+| + (1 - q)|01\rangle\langle 01|, 0 < q \leq 1$$

$$|\Phi^+\rangle = (|00\rangle + |11\rangle + |22\rangle + \dots + |(d-1)(d-1)\rangle)/\sqrt{d}$$

- Goal: find an optimal filter A s.t. the **overall** teleportation fidelity is maximal:

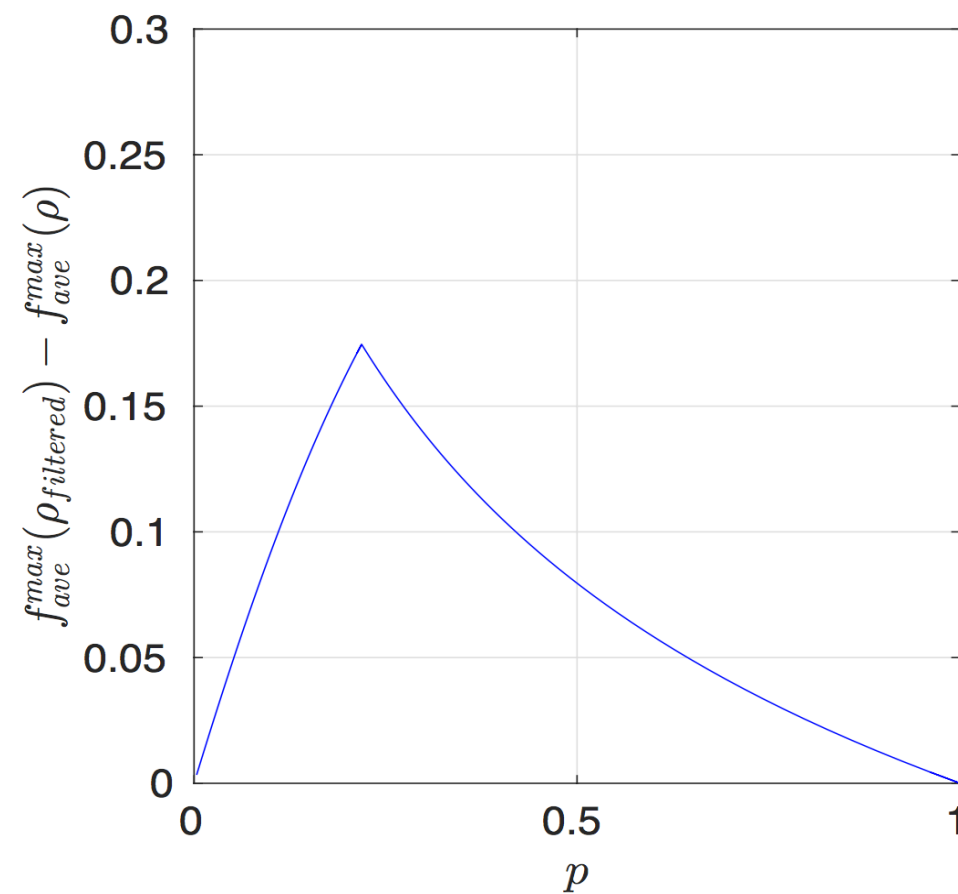
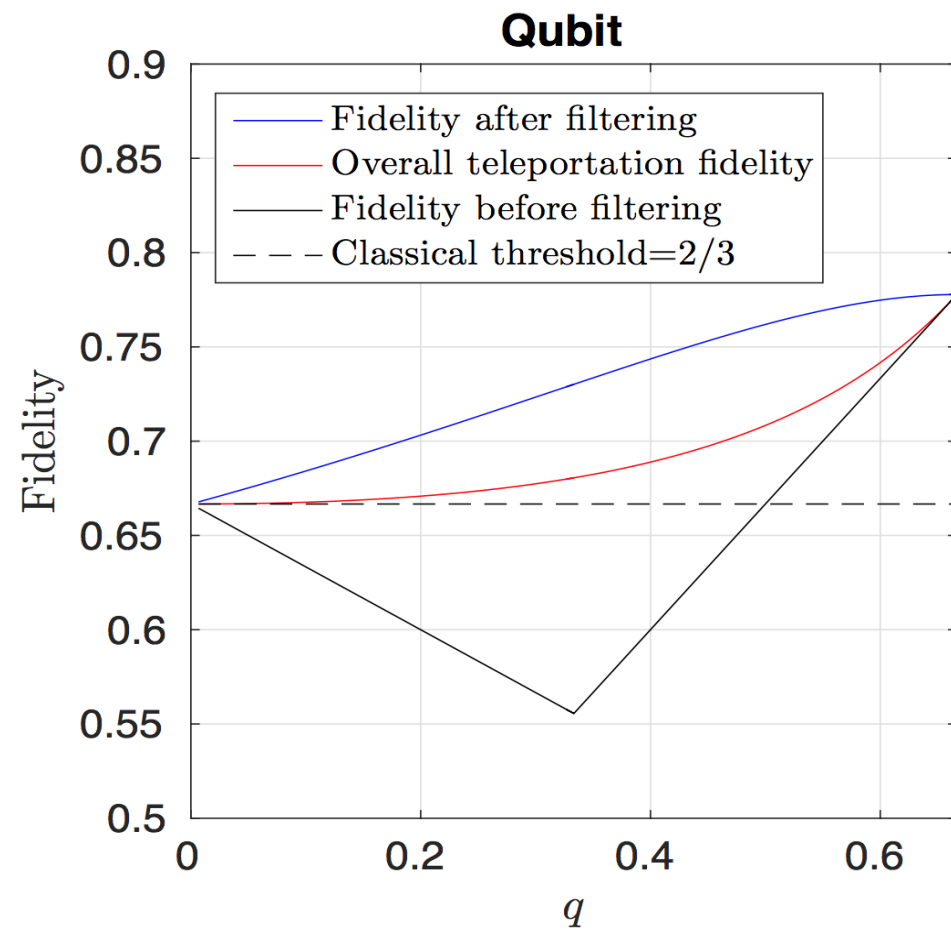
$$f_{\text{overall}} = pf_{\text{ave}}^{\text{max}}(\rho_{\text{filtered}}) + (1 - p)\frac{2}{d+1}$$

$p = \text{tr}[(A \otimes I)\rho(A \otimes I)^\dagger]$ is the **success probability** of local filtering,

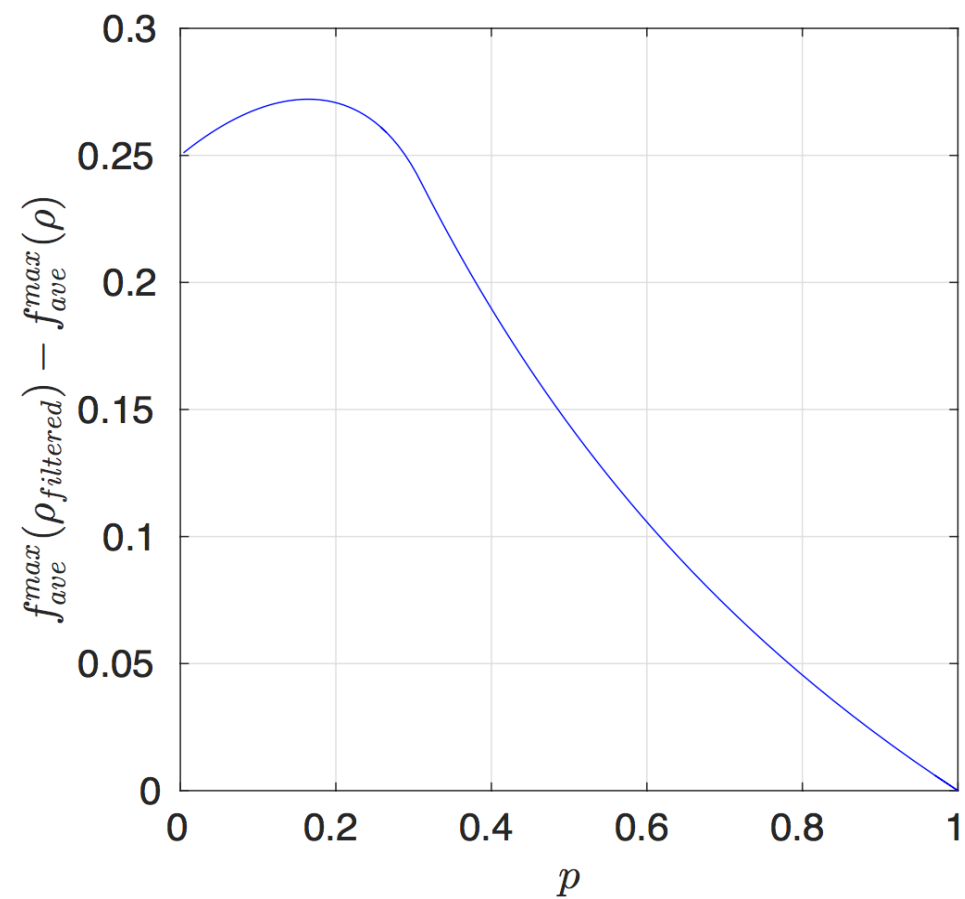
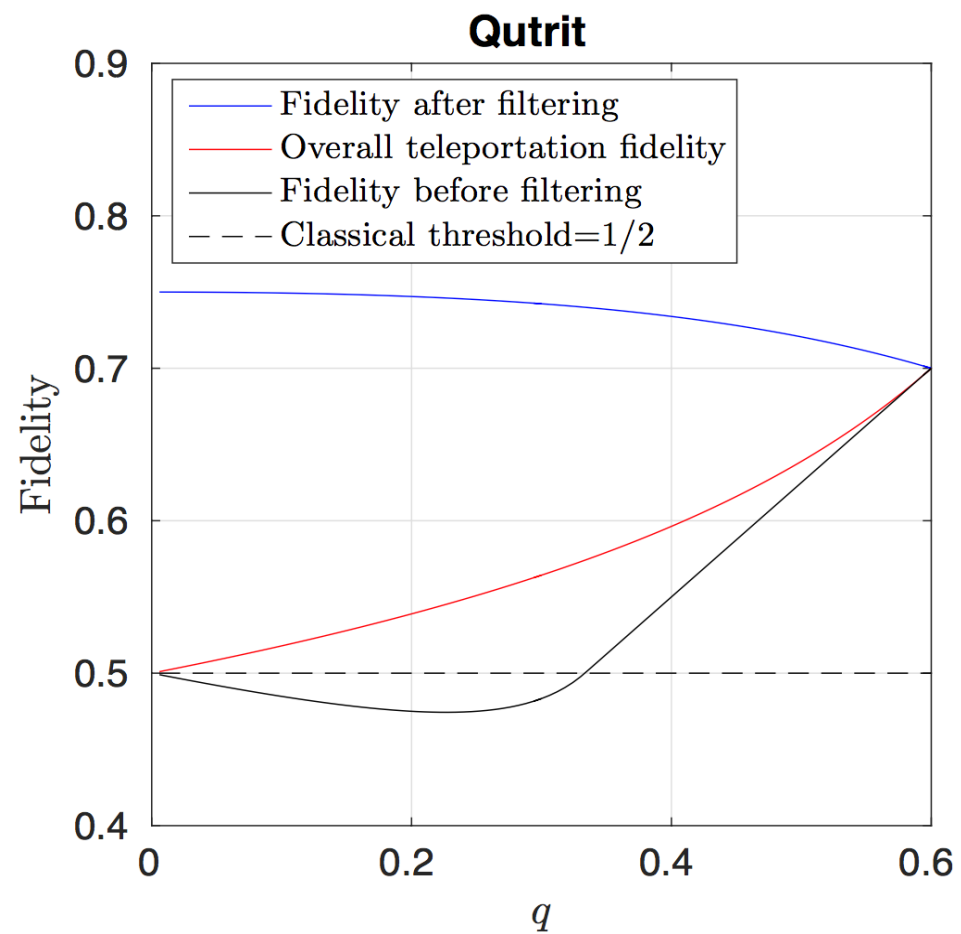
$\rho_{\text{filtered}} = \frac{(A \otimes I)\rho(A \otimes I)^\dagger}{p}$ is the state after **successful** local filtering.

- The optimal filter $A = \begin{bmatrix} \frac{(d-1)q}{d(q-1)} & 0 & \dots & 0 \\ 0 & 1 & & \vdots \\ \vdots & & \ddots & 0 \\ 0 & \dots & 0 & 1 \end{bmatrix}, 0 < q \leq \frac{d}{2d-1}$

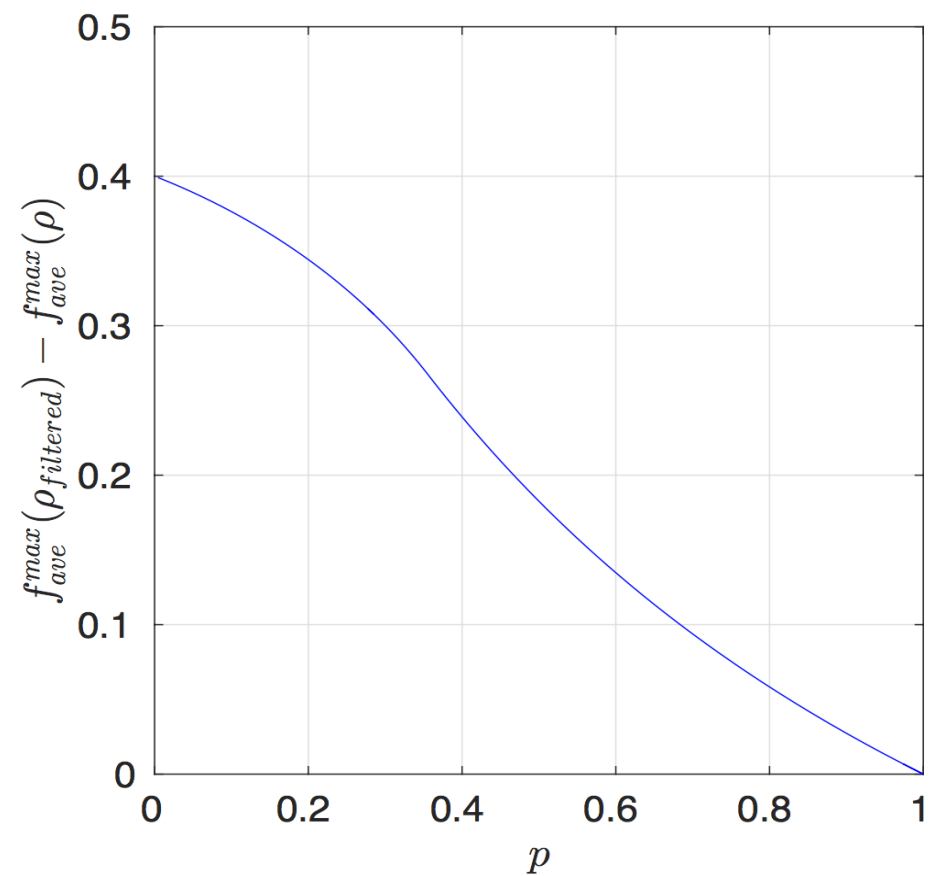
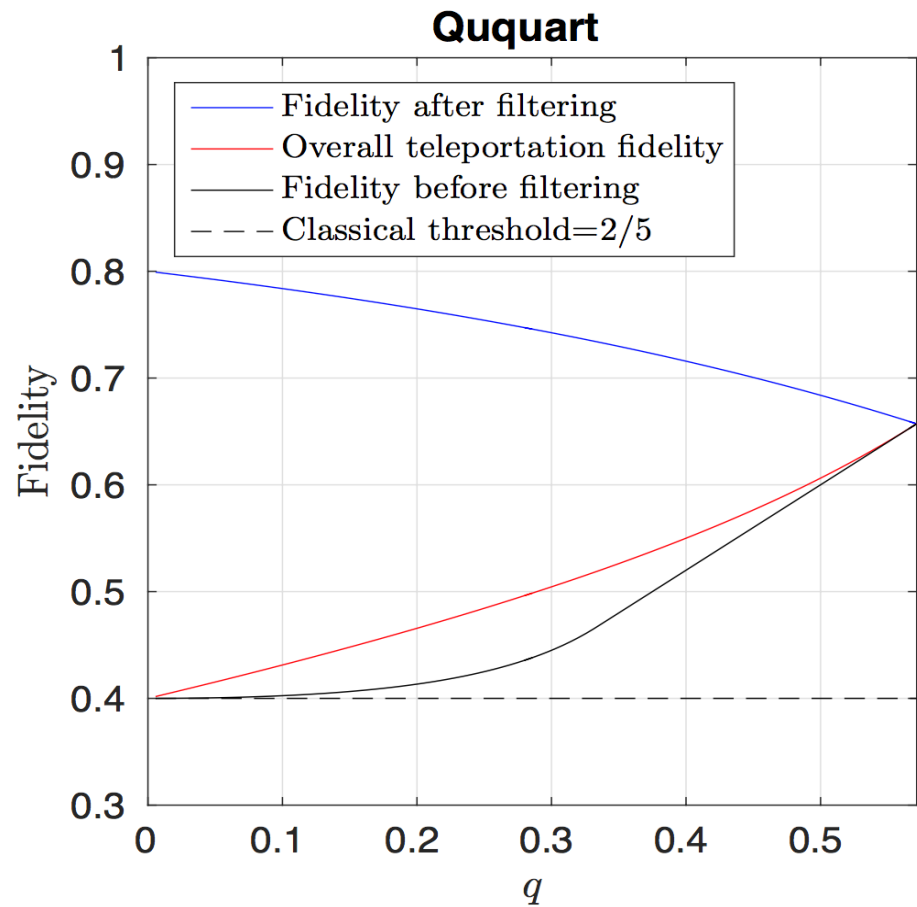
Results



Results



Results



Hidden teleportation power for Werner state

- Werner State:

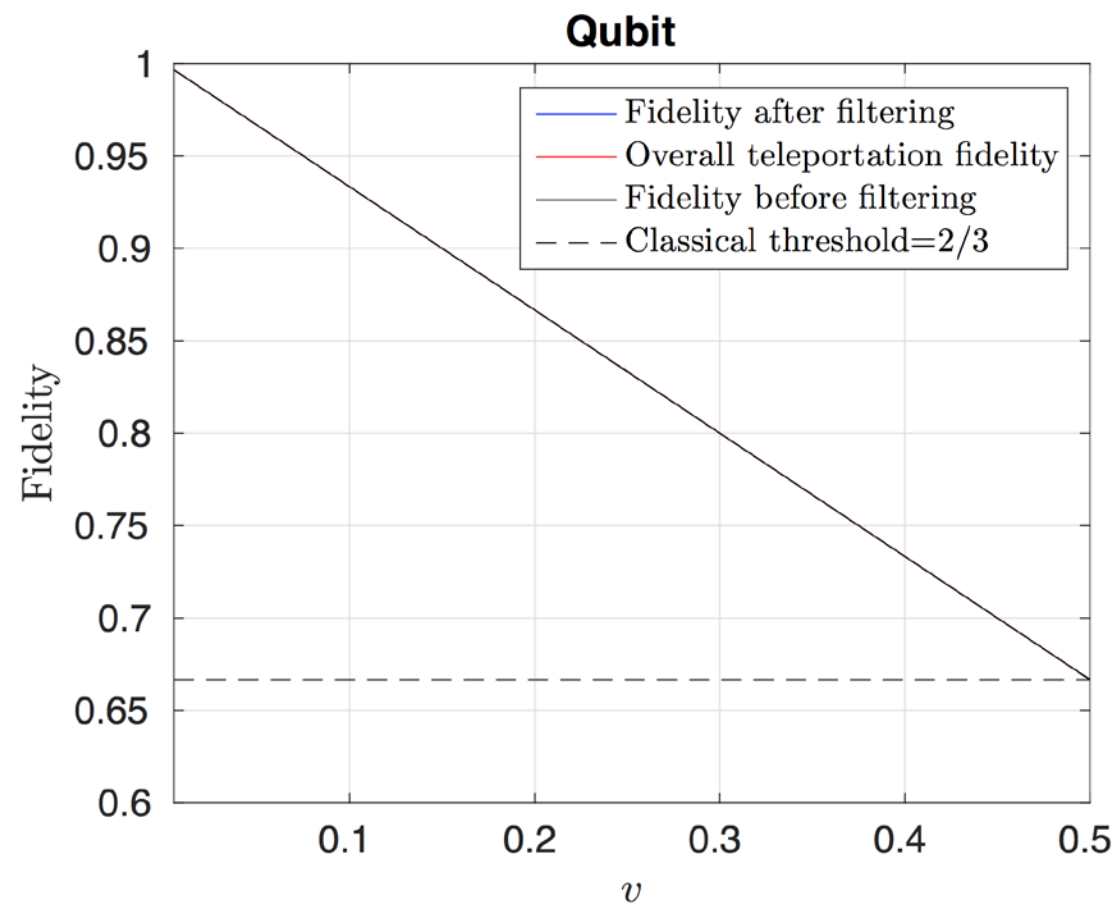
$$\rho(v) = v \frac{2}{d(d+1)} P_{sy} + (1-v) \frac{2}{d(d-1)} P_{asy}, \quad 0 \leq v \leq 1$$

$$\text{where } P = \sum_{i,j} |i\rangle\langle j| \otimes |j\rangle\langle i|, \quad P_{sy} = \frac{I+P}{2}, \quad P_{asy} = \frac{I-P}{2}$$

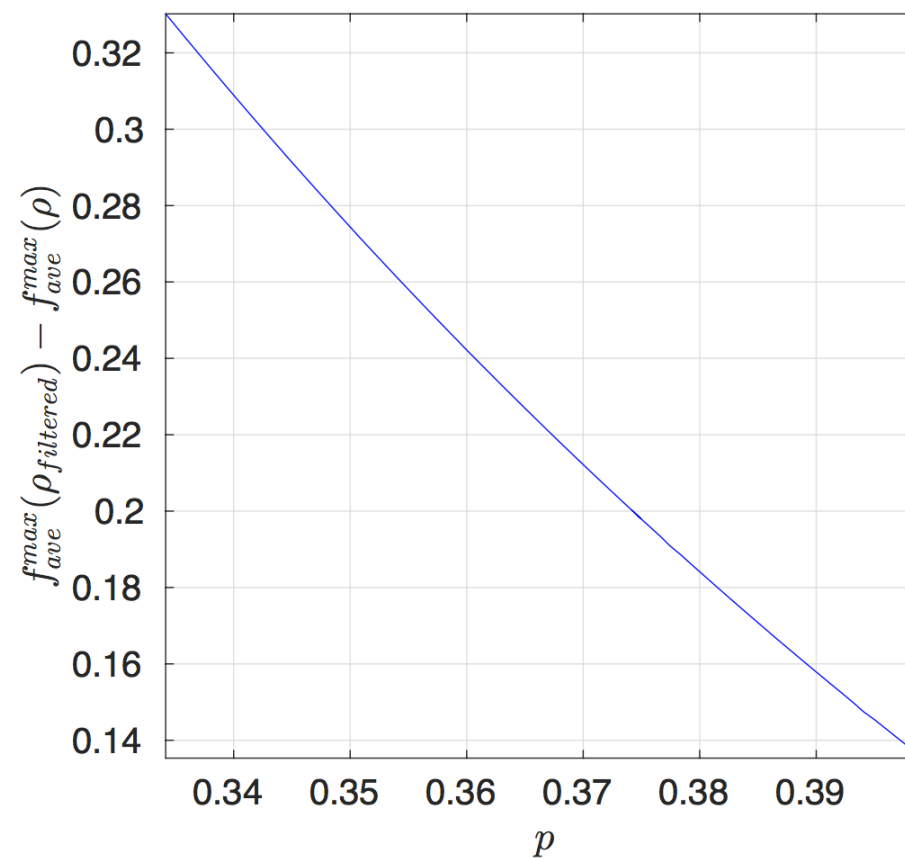
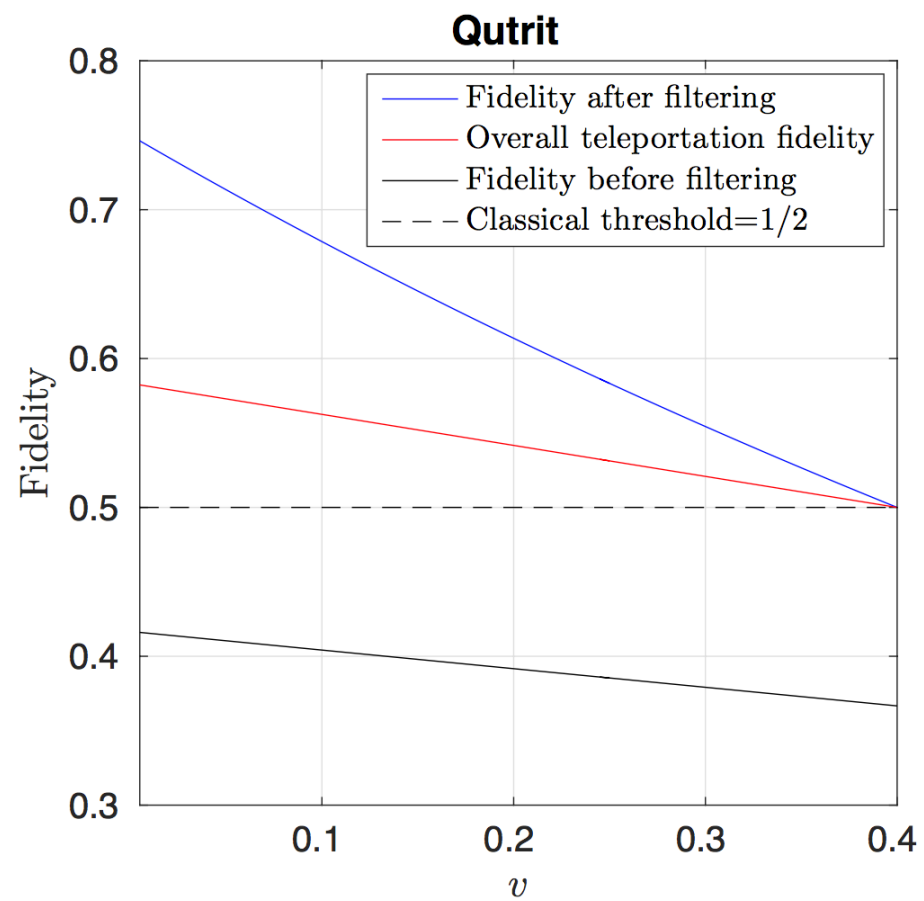
- The optimal filters A and B :

$$A = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix}, \quad B = \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix}, \quad 0 \leq v < \frac{d+1}{4d-2}$$

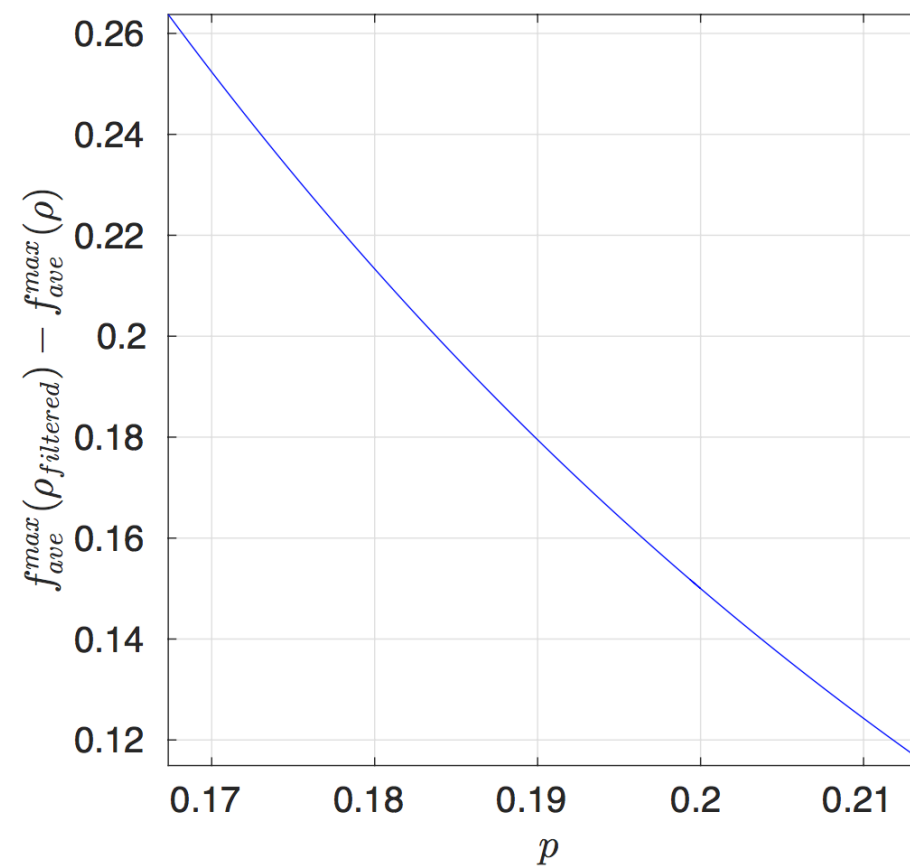
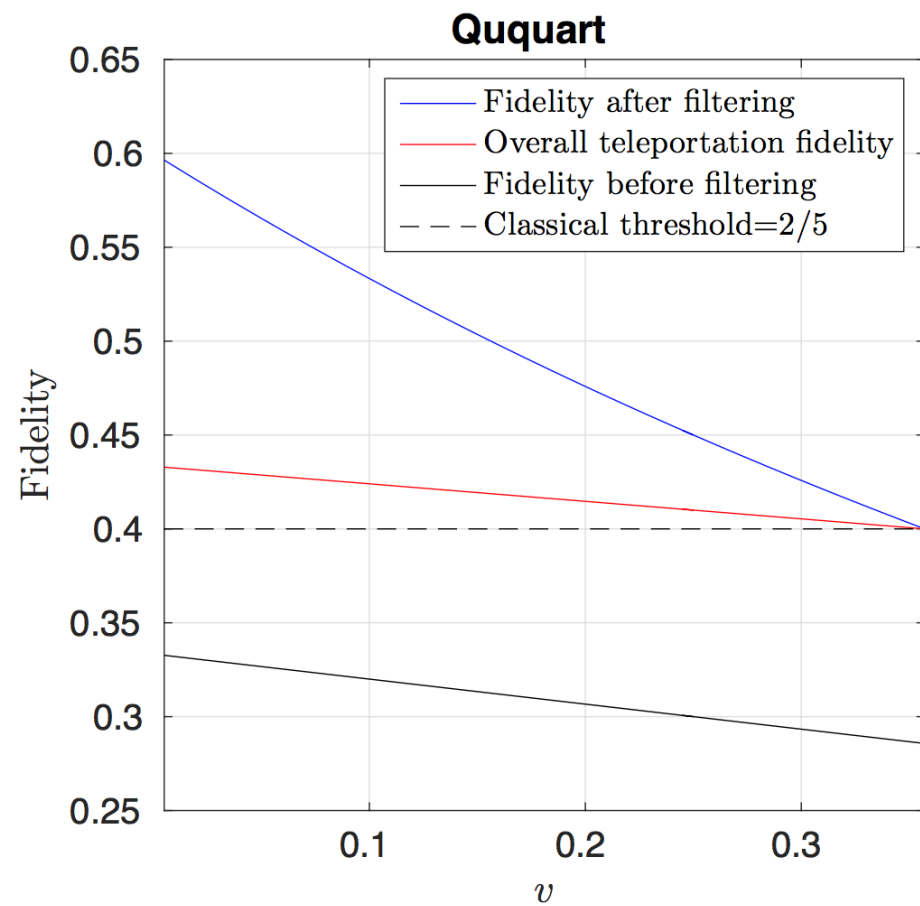
Results



Results



Results



Summary

- We summarize our results in the following table:

d	Teleportation power	Teleportation power after local filtering	Hidden teleportation power
2	useless ($0 < q \leq \frac{1}{2}$)	always useful	yes
3	useless ($0 < q \leq \frac{1}{3}$)	always useful	yes
4	always useful	always useful	no

d	Teleportation power	Teleportation power after local filtering	Hidden teleportation power
2	always useful	always useful	no
3	always useless	always useful	yes
4	always useless	always useful	yes

- There is a **tradeoff** between the **success probability** of local filtering and **how much fidelity can be increased** by local filtering.
- What about the **general** entangled state?

Back-up

- For qubit case, the maximal SF can be obtained by only consider 4 Bell states; but for higher dimensional case is not:

$$\rho(F) = q|\psi\rangle\langle\psi| + (1-q)|01\rangle\langle 01|$$

$$|\psi\rangle = (|00\rangle + |11\rangle + |22\rangle) / \sqrt{3}$$

$$SF = \max_{|\psi\rangle} \langle \psi | \rho | \psi \rangle = \max_{U_B} \langle \psi | (I \otimes U)^\dagger \rho (I \otimes U) | \psi \rangle$$

$$\bullet A = \begin{bmatrix} A_{11} & A_{12} & A_{13} \\ A_{21} & A_{22} & A_{23} \\ A_{31} & A_{32} & A_{33} \end{bmatrix} = e^{i\theta} \begin{bmatrix} a_{11} & a_{12} + ib_{12} & a_{13} + ib_{13} \\ a_{21} + ib_{21} & a_{22} + ib_{22} & a_{23} + ib_{23} \\ a_{31} + ib_{31} & a_{32} + ib_{32} & a_{33} + ib_{33} \end{bmatrix}$$

Back-up

- $$f_{overall} = \frac{1}{3} [1 + \frac{q}{3} |A_{11} + A_{22} + A_{33}|^2 + (1-q) |A_{21}|^2 - \frac{q}{3} (|A_{11}|^2 + |A_{22}|^2 + |A_{33}|^2) - \frac{q}{3} (\text{all the off-diagonal terms of A}) - (1-q) (|A_{11}|^2 + |A_{21}|^2 + |A_{31}|^2)]$$

$$= \frac{1}{3} [1 + (q-1)a_{11}^2 + \frac{2q}{3} a_{11}(a_{22} + a_{33}) + \frac{2q}{3} (a_{22}a_{33} + b_{22}b_{33})]$$

- Generalize to d-dimensional quantum state:

$$f_{overall} = \frac{1}{d} [1 + (q-1) a_{11}^2 + \frac{2q}{d} a_{11} \sum_{i=2}^d a_{ii} + \frac{2q}{d} \sum_{i=2, j>i}^{d-1} a_{ii} a_{jj}]$$

$$= \frac{1}{d} \{ 1 + (q-1) a_{11}^2 + \frac{q(d-1)}{d} [2a_{22} + (d-2)] \}$$

$$* \frac{\partial K_d}{\partial a_{11}} = 0, \text{ we can find } a_{11} = \frac{(d-1)q}{d(q-1)}$$

$$f_{overall} = \frac{1}{d} [1 + \frac{q^2(d-1)^2}{d^2(1-q)} + \frac{q(d-1)(d-2)}{d}].$$

$$A = \begin{bmatrix} \frac{(d-1)q}{d(q-1)} & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & 1 \end{bmatrix}$$

Back-up

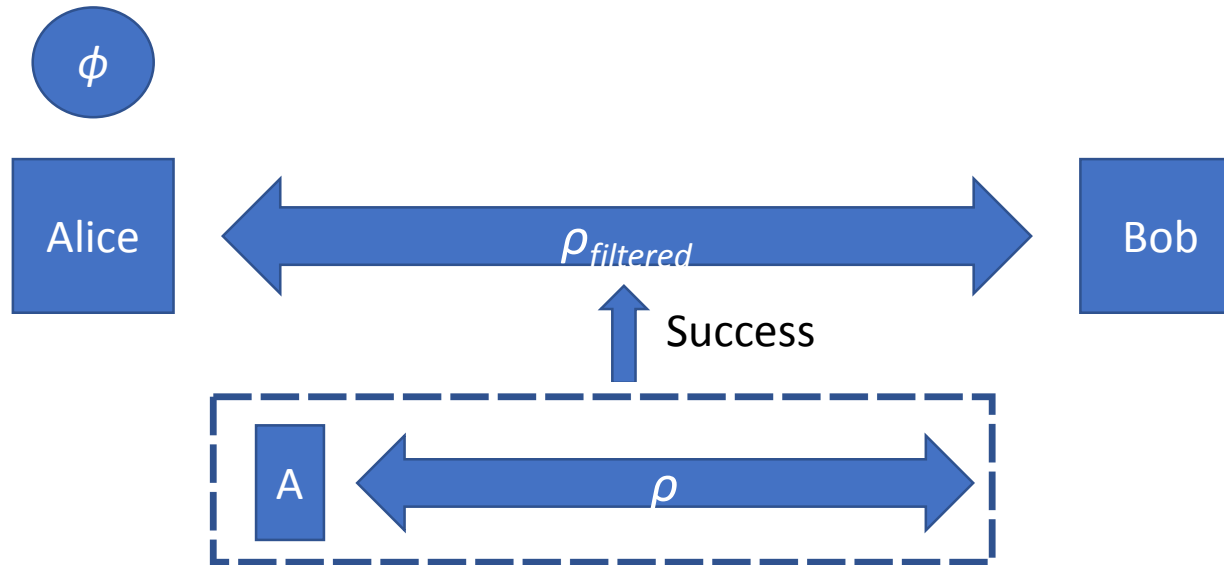
- The problem of hidden teleportation for qubit case has been studied: Optimal teleportation with a mixed state of two qubit (F. Verstraete, H. Verschelde, 2003)

Theorem:

The bipartite local filtering operations probabilistically bring an entangled mixed state of two qubits to a state with highest possible fidelity are given by the filtering operations bringing the state into its unique Bell-diagonal normal form, yielding a fidelity larger than $1/2$.

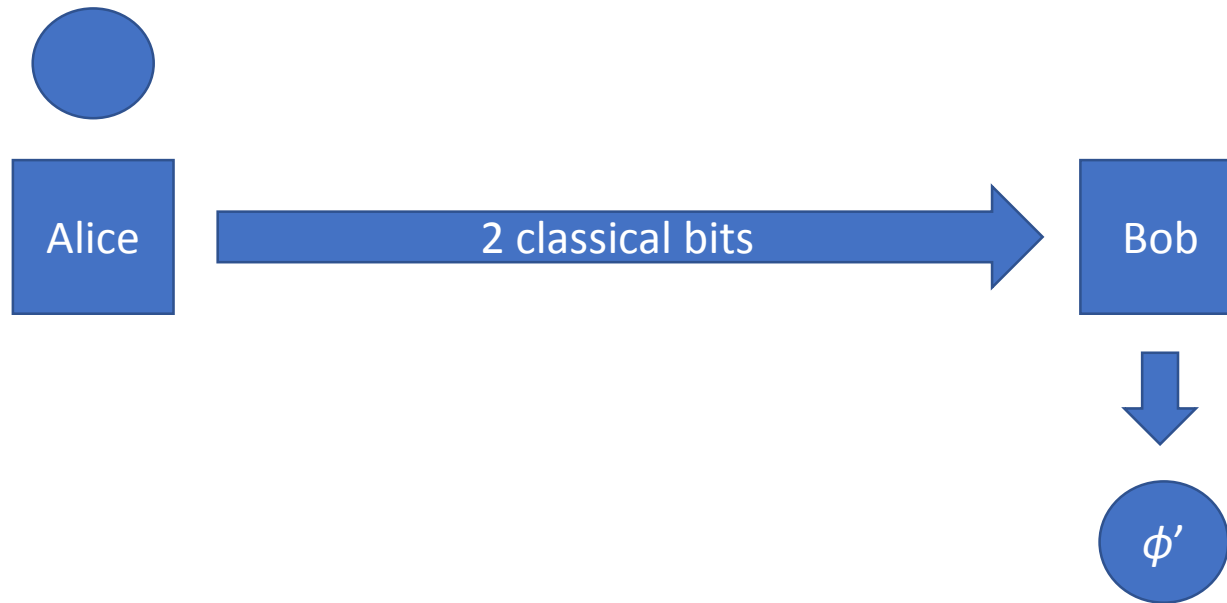
Protocol

- Success



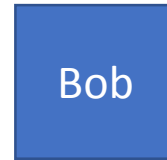
Protocol

- Success



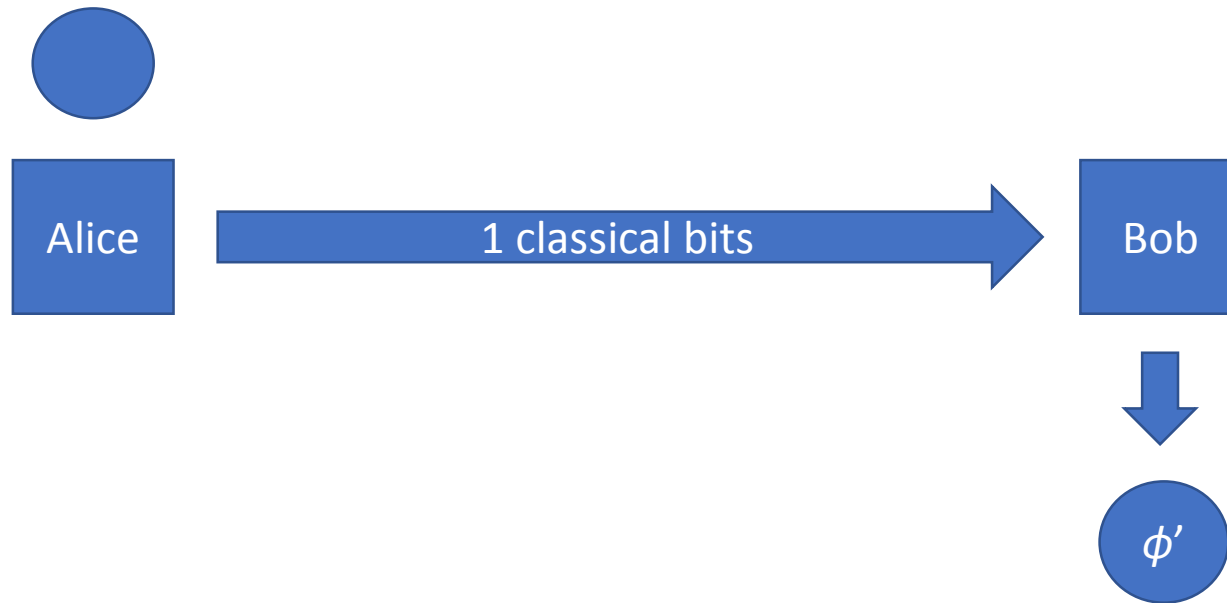
Protocol

- Fail



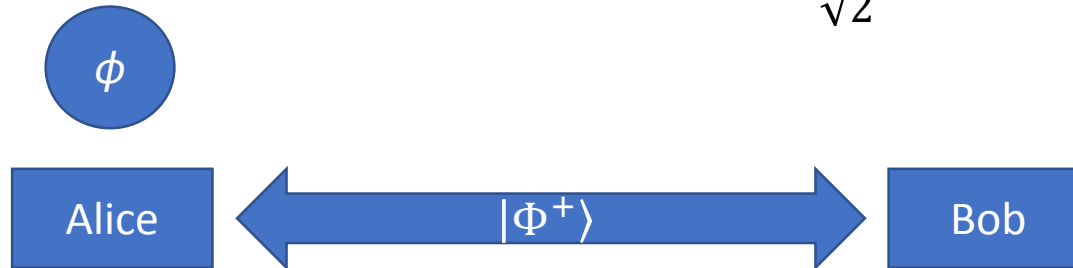
Protocol

- Fail



Back-up

- Quantum Teleportation [1]
 $|\phi\rangle = \alpha|0\rangle + \beta|1\rangle$



Bell basis :

$$|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

$$|\Phi^-\rangle = \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle)$$

$$|\Psi^+\rangle = \frac{1}{\sqrt{2}}(|01\rangle + |10\rangle)$$

$$|\Psi^-\rangle = \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle)$$

$\Rightarrow$

$$|00\rangle = \frac{1}{\sqrt{2}}(|\Phi^+\rangle + |\Phi^-\rangle)$$

$$|11\rangle = \frac{1}{\sqrt{2}}(|\Phi^+\rangle - |\Phi^-\rangle)$$

$$|01\rangle = \frac{1}{\sqrt{2}}(|\Psi^+\rangle + |\Psi^-\rangle)$$

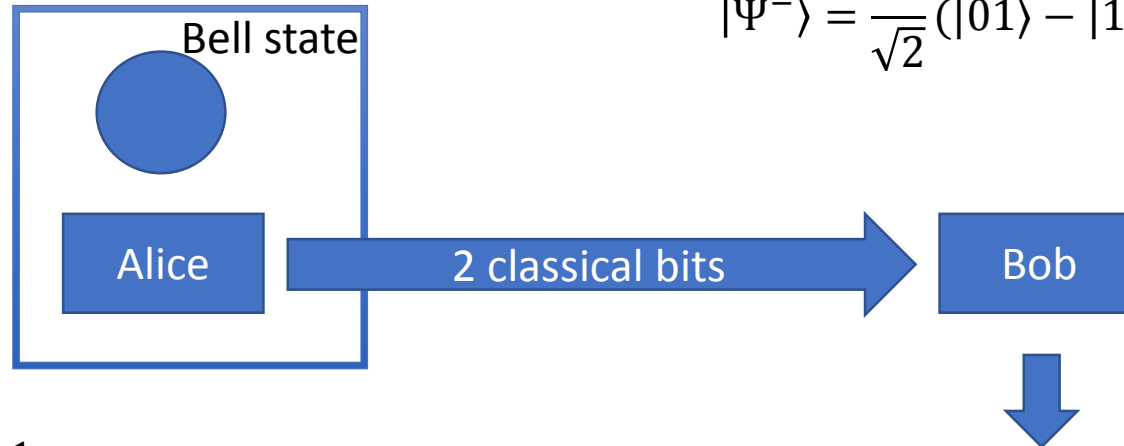
$$|10\rangle = \frac{1}{\sqrt{2}}(|\Psi^+\rangle - |\Psi^-\rangle)$$

- $$\begin{aligned}
 |\phi\rangle|\Phi^+\rangle &= \frac{1}{2}(\alpha|000\rangle + \alpha|011\rangle + \beta|100\rangle + \beta|111\rangle) \\
 &= \frac{1}{2}[|\Phi^+\rangle(\alpha|0\rangle + \beta|1\rangle) + |\Phi^-\rangle(\alpha|0\rangle - \beta|1\rangle) \\
 &\quad + |\Psi^+\rangle(\alpha|1\rangle + \beta|0\rangle) + |\Psi^-\rangle(\alpha|1\rangle - \beta|0\rangle)]
 \end{aligned}$$

Back-up

- Quantum Teleportation [1]

$$|\phi\rangle = \alpha|0\rangle + \beta|1\rangle$$



- $$|\phi\rangle|\Phi^+\rangle = \frac{1}{2}(\alpha|000\rangle + \alpha|011\rangle + \alpha|100\rangle + \alpha|111\rangle)$$

$$= \frac{1}{2}[|\Phi^+\rangle(\alpha|0\rangle + \beta|1\rangle) + |\Phi^-\rangle(\alpha|0\rangle - \beta|1\rangle)$$

$$+ |\Psi^+\rangle(\alpha|1\rangle + \beta|0\rangle) + |\Psi^-\rangle(\alpha|1\rangle - \beta|0\rangle)]$$

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$\Rightarrow$

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$$|10\rangle = \frac{1}{\sqrt{2}}(|\Psi^+\rangle - |\Psi^-\rangle)$$