

# Sonic Compton scattering in an analogue gravity model

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Analogue gravity models provide an indirect way to probe the physics of relativistic systems for which the actual (i.e., non-analogue) experiments are currently inaccessible, provided that the analogue model can be faithfully mapped to the actual physical system of interest. Here, we investigate a simple toy model of phonon scattering from a Newtonian particle as an analogue to Lorentz-violating Compton scattering. In particular, we wish to know what *in-universe observers* – observers who use sound in their analogue gravity medium to operationally measure distance and duration [1] – can determine from carrying out scattering experiments of this kind. Unsurprisingly, the Newtonian particle appears to be a Lorentz-violating particle to *in-universe observers*; scattering experiments from these particles can be used by *in-universe observers* to determine their absolute state of motion with respect to the medium.

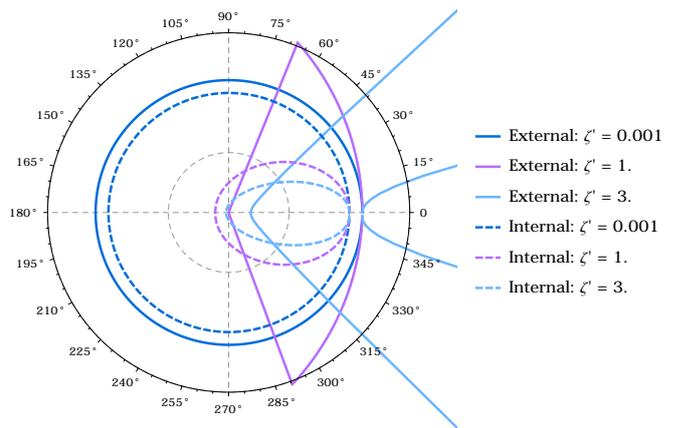
Hawking’s famous result that black holes should act as black-body emitters was built on assumptions whose validity is questionable. Unruh used a hydrodynamical model as a sonic (or acoustic) analogue to a black hole to study, by analogy, how a breakdown of these assumptions might manifest (provided, of course, that the analogy is appropriate) [2]. This started the modern day analogue gravity research effort, and follow-up work by Unruh and others suggests that Hawking’s result should be robust to a breakdown of his original assumptions.

In an effort to understand how far analogue gravity models can be pushed, we investigate a toy model of Lorentz-violating Compton scattering and compare it with a toy model of Lorentz-obeying Compton scattering. We start with an analogue-gravity medium whose excitations are phonons that we treat as a scalar field; the phonon field obeys the Klein-Gordon equation and thus admits a (sonic-)Lorentz symmetry with characteristic speed  $c_s$ , the speed of sound in the analogue gravity medium. The frame in which the analogue gravity medium is stationary defines the *laboratory frame*, and it is in this frame that we define all of our dynamical physical quantities and perform all calculations. In the laboratory frame, the phonon dispersion relationship is  $E = pc_s$  with  $\mathbf{p} = \hbar\mathbf{k}$ .

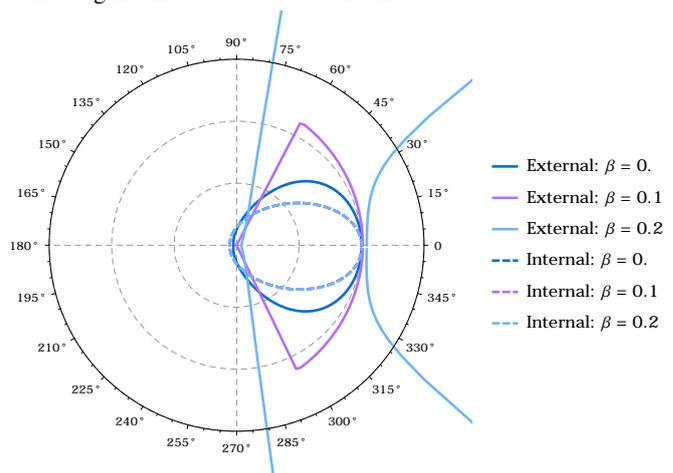
We then separately consider phonon scattering from two types of particles:

- i. *Internal particles* possess a sonically-relativistic dispersion relationship  $\gamma_s^2 E^2 = \gamma_s^2 p^2 c_s^2 + m^2 c_s^4$  (where here  $E$  and  $\mathbf{p}$  have their usual Newtonian definition and are measured in the laboratory frame, and  $\gamma_s$  is the sonic Lorentz-factor which is defined as per the usual Lorentz factor but with characteristic speed  $c_s$  the speed of sound);
- ii. *External particles* possess a Newtonian dispersion relationship  $E = p^2/2m$  (where, again,  $E$  and  $\mathbf{p}$  have their usual Newtonian definition).

By performing the sonic analogue to Compton scattering from external particles, *in-universe observers* can determine their state of motion with respect to their medium. By comparison, phonon scattering from internal particles is a process that is entirely Lorentz obeying, and reveals nothing to *in-universe observers* about their state of motion with respect to their medium.



(a)  $\beta := v/c_s \approx 0.3106$ . For scattering from external particles,  $\zeta' = 0.001$  results in scattering at all angles;  $\zeta' = 1$  is a critical value for which scattering is restricted to a closed window of angles;  $\zeta' = 3$  is past the critical value of  $\zeta'$ : here we see two disjoint curves, though only the “inner” (or rightmost) such curve is physical. Note that in this plot we have  $\zeta' := \hbar\omega'/mc_s^2$ , where  $\zeta$  and  $\omega$  are primed, i.e. these are co-moving *in-universe observer* frame measured values.



(b)  $\zeta := \hbar\omega/mc_s^2 = 1.8$ . For scattering from external particles,  $\beta = 0$  results in scattering at all angles;  $\beta = 0.1$  is a critical value for which scattering is restricted to a closed window of angles;  $\beta = 0.2$  is past the critical value of  $\beta$ : here we see two disjoint curves, though only the “inner” (or rightmost) such curve is physical. Note that all of the dashed curves sit on-top of one another. Finally, note that  $\zeta$  and  $\omega$  are unprimed, i.e. these are laboratory frame measured values.

Figure 1: Differential scattering cross-sections of phonons scattering from external (solid) and internal (dashed) particles in the co-moving *in-universe observer* frame where (1a) phonons have different initial energies and particles have fixed initial velocity, and (1b) phonons have fixed initial energy and particles have different initial velocities.

## References

- [1] Todd, S.L., Menicucci, N.C., “Sound Clocks and Sonic Relativity,” *Foundations of Physics* **47**, 1267–1293 (2017)
- [2] Unruh, W. G., “Experimental Black-Hole Evaporation?,” *Physical Review Letters* **46**, 1351–1353 (1981).