

# Magnon Kerr effect in a strongly coupled cavity-magnon system

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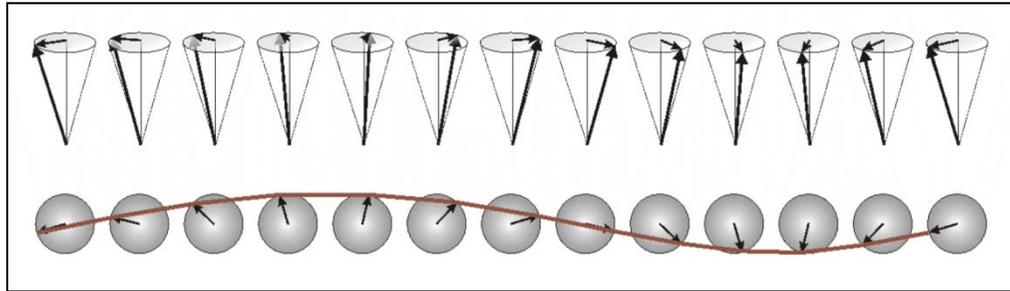
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# Outline

- **Introduction**
  - Spin waves (magnons)
- Cavity quantum electrodynamics (QED) with magnons  
D. Zhang *et al.*, npj Quantum Informtion 1, 15014 (2015).
- Magnon Kerr effect in a cavity QED system  
Y.P. Wang *et al.*, Phys. Re. B 94, 224410 (2016).
- Conclusions & Outlook

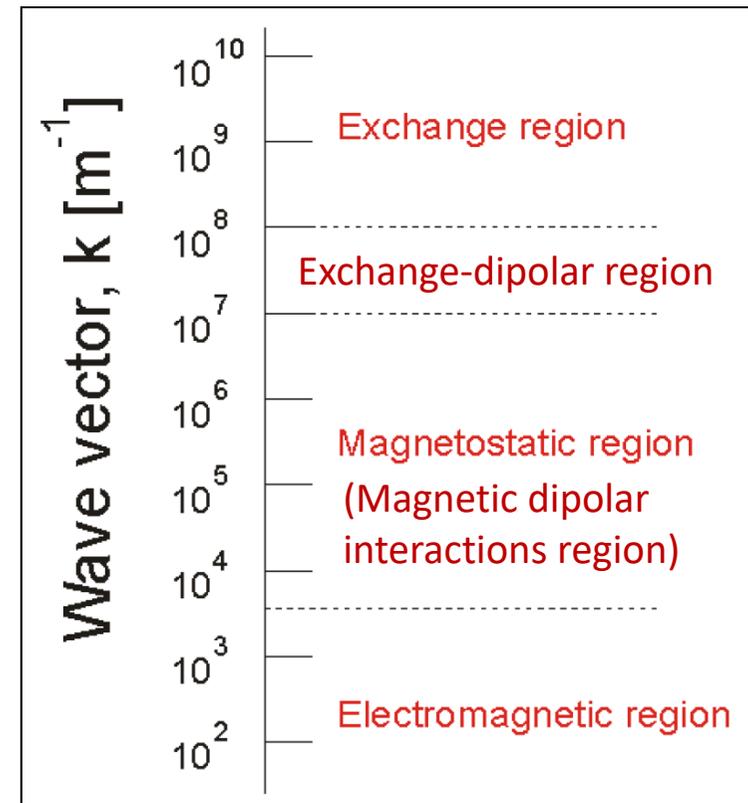
# Spin waves (magnons)



Landau-Lifshitz equation

$$\frac{d}{dt} \vec{M} = -\gamma \mu_0 \vec{M} \times \vec{H}_{\text{eff}} + T_D$$

+ Maxwell equations



M. Cottam and D. Tilley, *Introduction to surface and superlattice excitations* (IoP, 2004).

- Robust extended spatial mode
- High spin density  $10^{21} - 10^{22} \text{ cm}^{-3}$

Quantum theory  
for spin waves in the  
long-wavelength limit

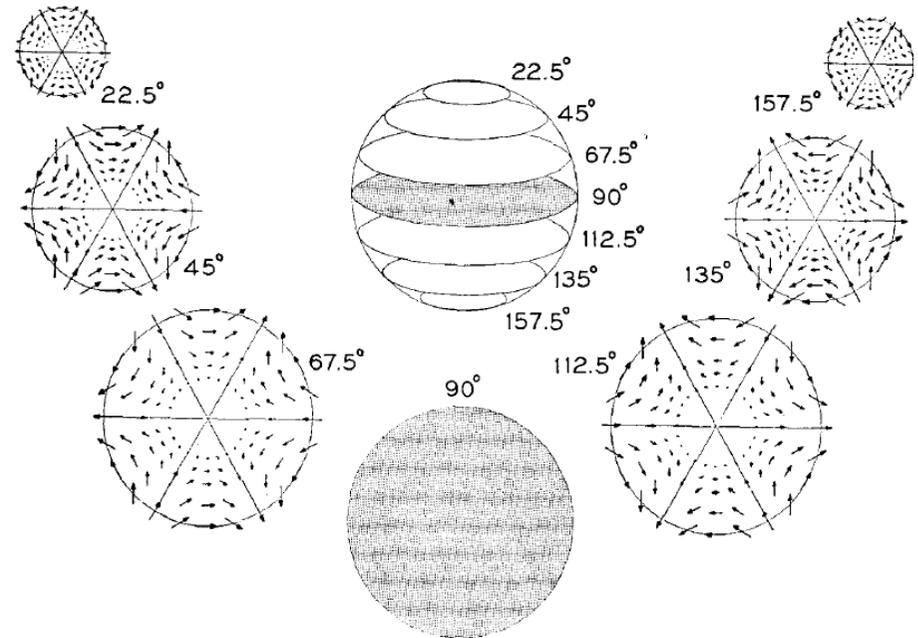
# Magnetostatic wave

Magnetic dipolar interactions (dominating)

Exchange interactions (ignored)

$$\Lambda k^2 \ll 1, \Lambda = \text{exchange constant}$$

- Ferromagnetic resonance (FMR) mode  
(Kittel mode; uniform precession mode)
- Magnetostatic (MS) modes  
(Non-uniform precession modes)



L.R. Walker, J. Appl. Phys. 29, 318 (1958).

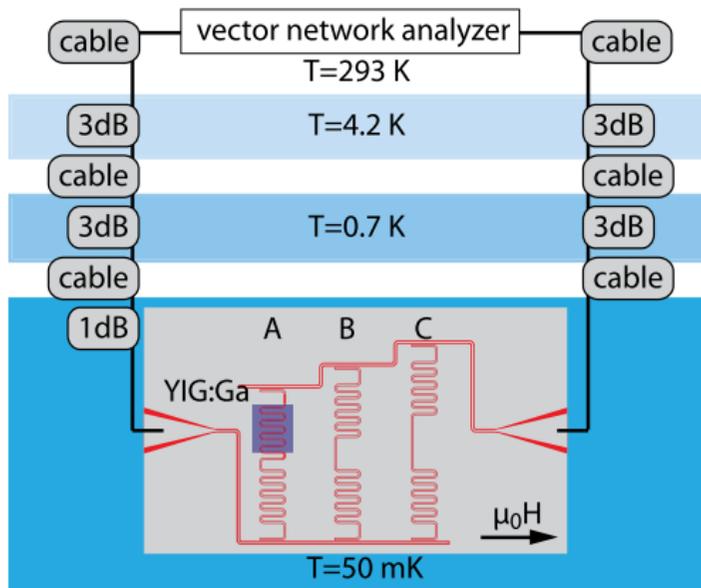
FMR mode in ferromagnetic sphere:

$$f_m = \frac{\gamma_e}{2\pi} |\bar{B}_0| + f_{m,0} \quad \begin{array}{l} \gamma_e/2\pi = 28 \text{ GHz/T, electron gyromagnetic ratio} \\ f_{m,0} < 10 \text{ MHz, determined by anisotropy field} \end{array}$$

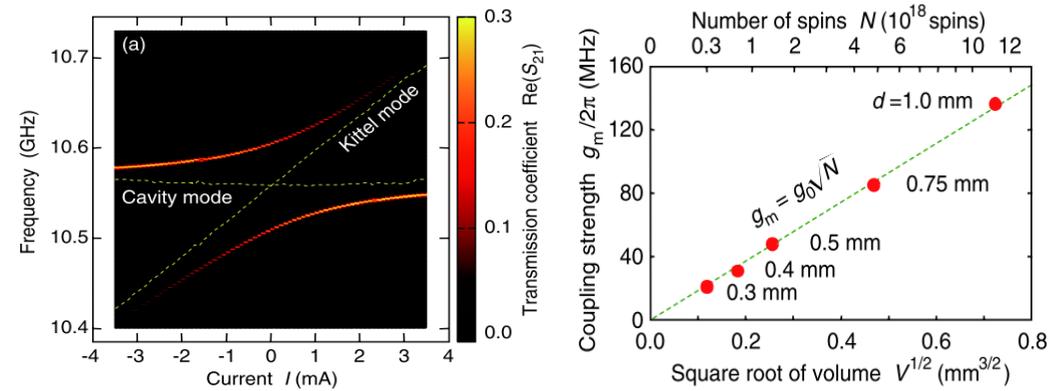
**YIG** = yttrium iron garnet (Y<sub>3</sub>Fe<sub>5</sub>O<sub>12</sub>)

Quantum  
magnonics

# Recently reported experimental results on YIG sphere in cavity

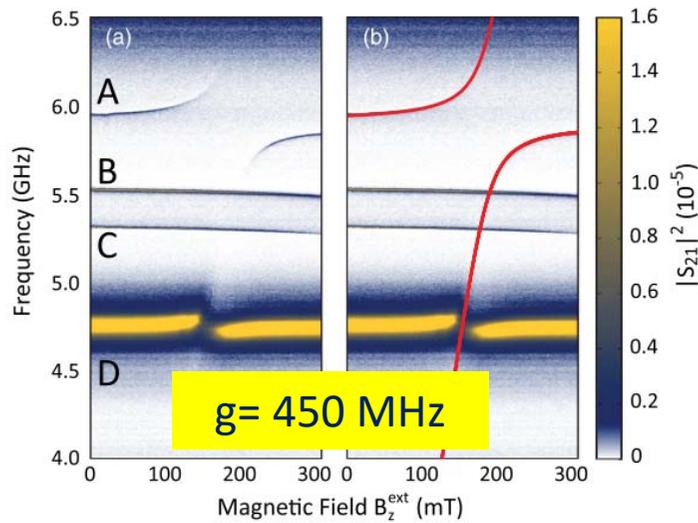


## Single photon level

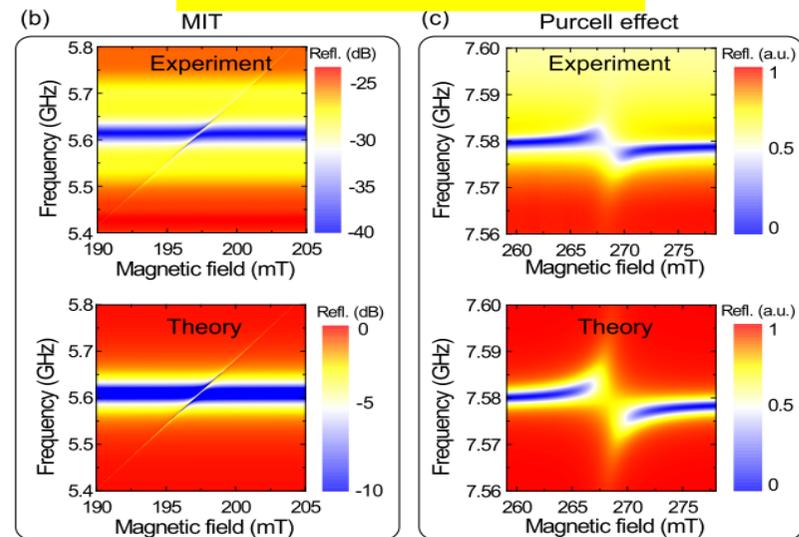


Y. Tabuchi *et al.*, *Phys. Rev. Lett.* **113**, 083603 (2014).

## Room temperature



H. Huebl *et al.*, *Phys. Rev. Lett.* **111**, 127003 (2013).



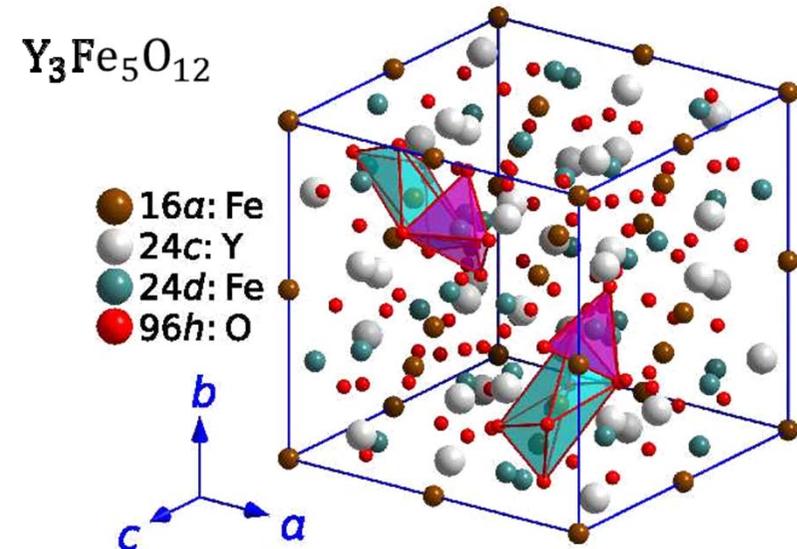
X. Zhang *et al.*, *Phys. Rev. Lett.* **113**, 156401 (2014).

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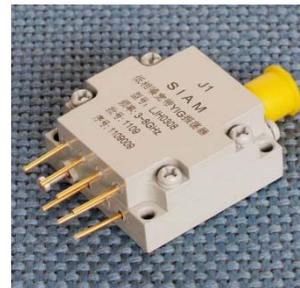
# Yttrium-iron-garnet (YIG, $Y_3Fe_5O_{12}$ )

- Ferrimagnetic insulator
- Narrow FMR linewidth
- Transparent at infrared
  
- High Curie temperature:  $\sim 550$  K
- Large spin density:  $4.2 \times 10^{21} \text{ cm}^{-3}$



Classic applications:

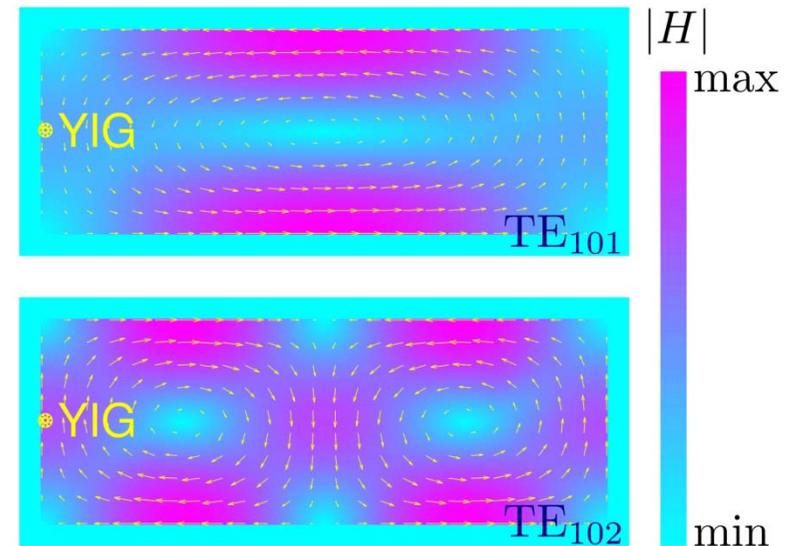
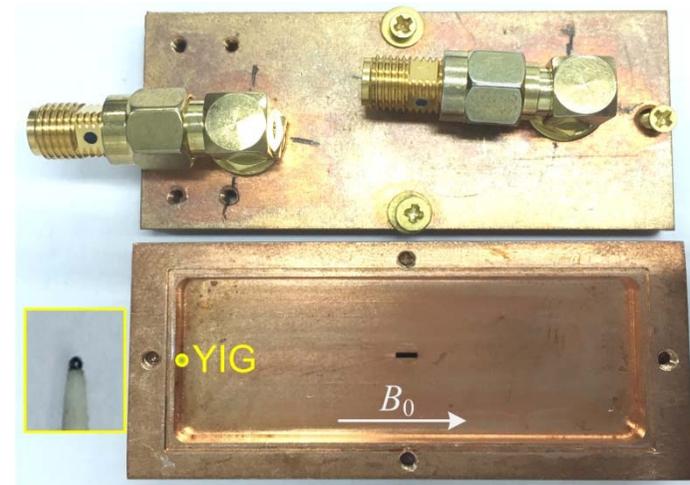
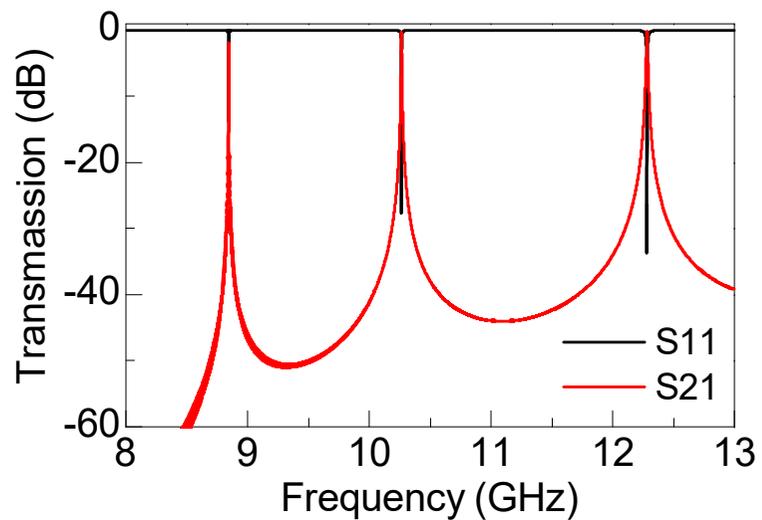
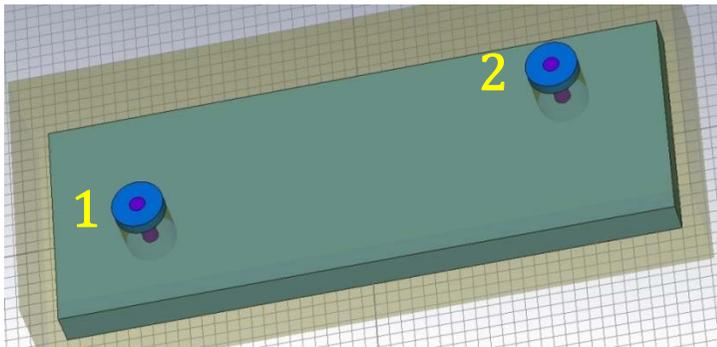
- Oscillator/ Filter
- Isolator/ Circulator



# 3D rectangular copper cavity

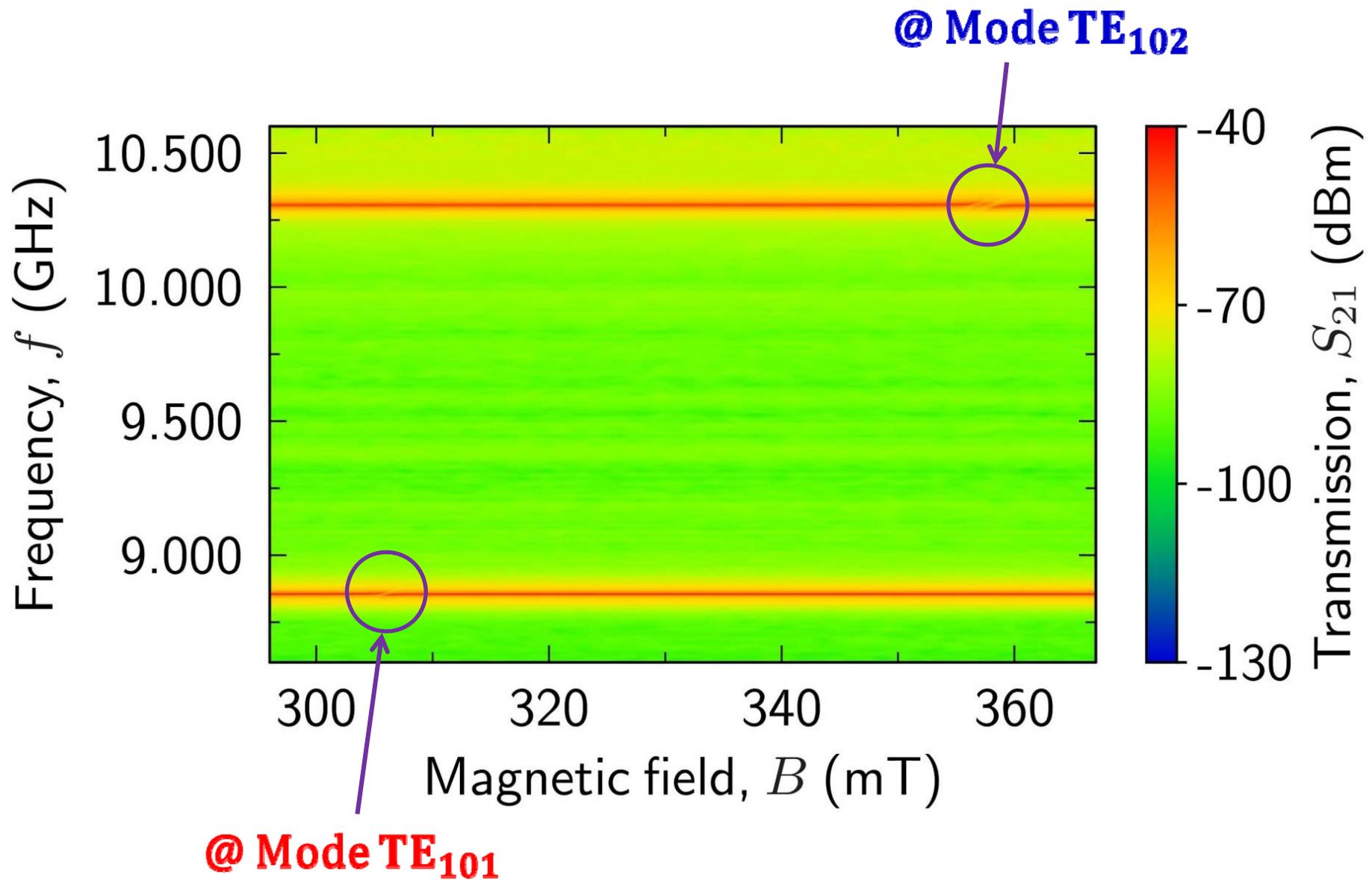
Two ports

Cavity size:  $50 \times 18 \times 3$  (mm<sup>3</sup>)



YIG diameter: 0.32 mm

# $S_{21}$ spectrum of cavity with a YIG sphere @ 22 mK



# $S_{21}$ @ cryogenic temperature (Cryo., 22 mK)

@ 22 mK average thermal photon number  $\sim 1 \times 10^{-2}$

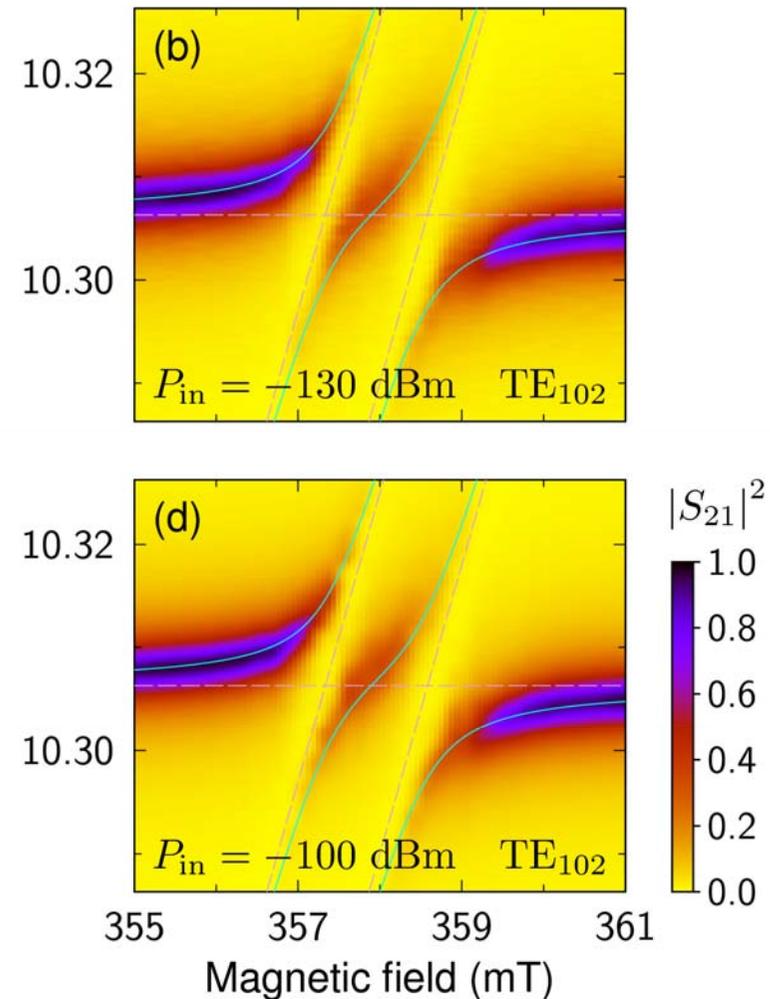
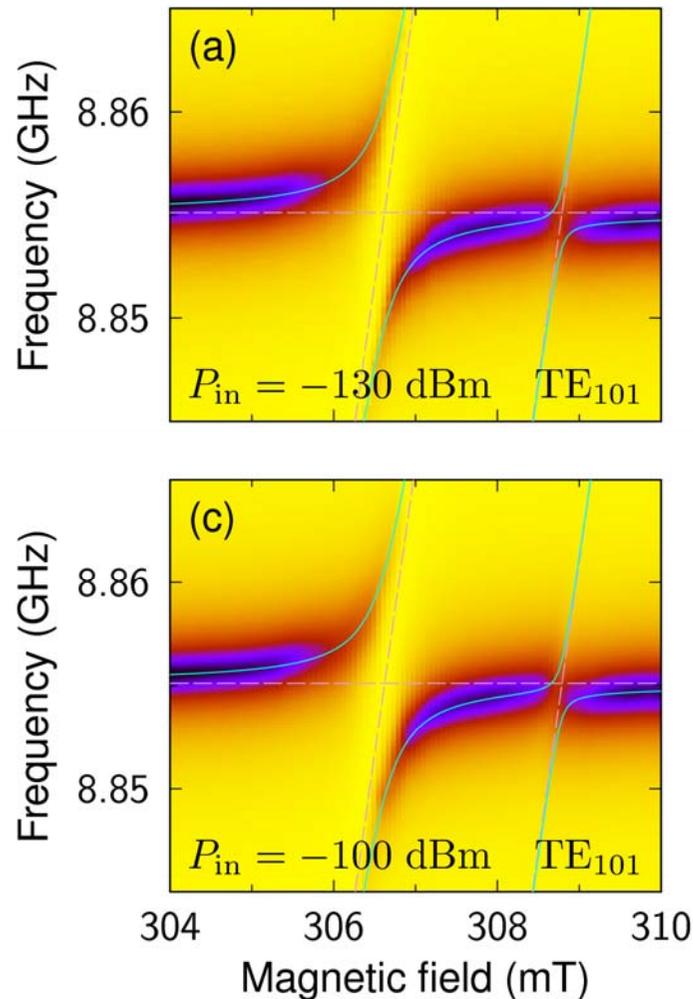
Average  
photon  
numbers

0.8 (<1)

800

@ Mode  $TE_{101}$

@ Mode  $TE_{102}$



# $S_{21}$ @ room temperature (R. T., 300 K)

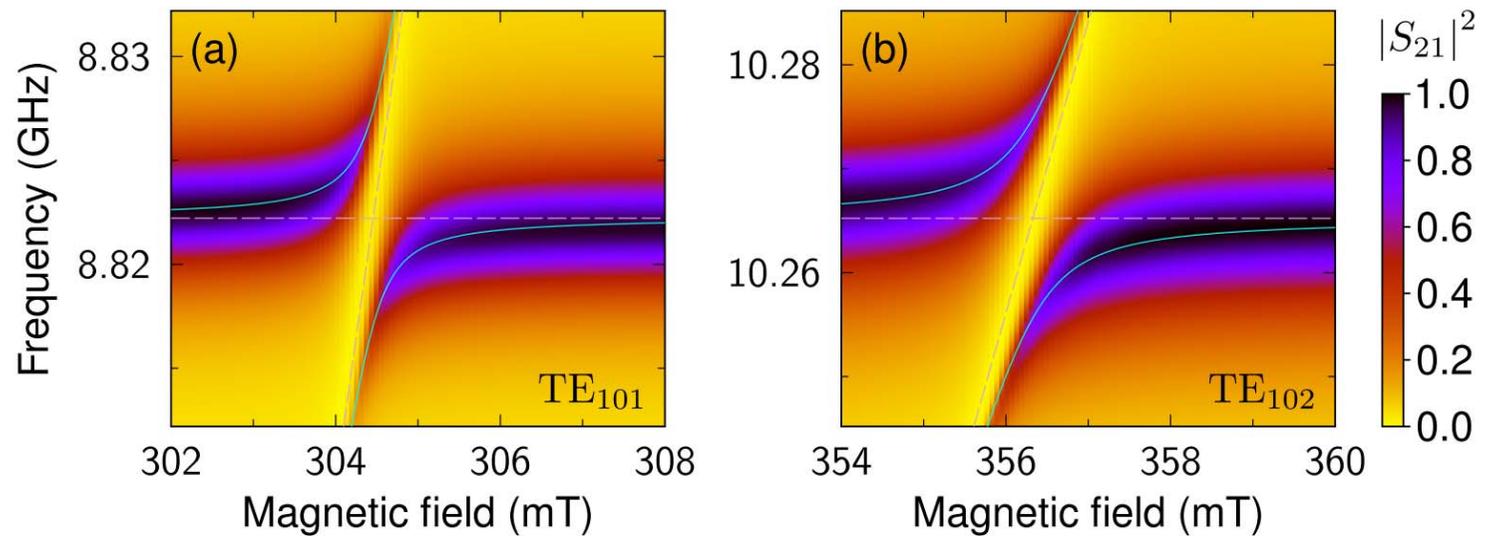
@ 300 K average thermal photon number  $\sim 700$

Average  
photon  
numbers

$1.8 \times 10^{10}$

@ Mode  $TE_{101}$

@ Mode  $TE_{102}$



The MS mode related anti-crossing has disappeared.

## $S_{21}$ spectrum – Input-output theory

$$S_{21}(\omega) = \frac{2\sqrt{\kappa_i \kappa_o}}{i(\omega - \omega_c) - \kappa_{\text{tot}} + \sum(\omega)},$$

$$\sum(\omega) = \frac{\tilde{g}_{\text{FMR}}^2}{i(\omega - \omega_{\text{FMR}}) - \gamma_{\text{FMR}}} + \frac{\tilde{g}_{\text{MS}}^2}{i(\omega - \omega_{\text{MS}}) - \gamma_{\text{MS}}}$$

Coupling strength  $\tilde{g}_m$  ( $m = \text{FMR}, \text{MS}$ )

Total cavity decay rate  $\kappa_{\text{tot}} = \kappa_i + \kappa_o + \kappa_{\text{int}}$

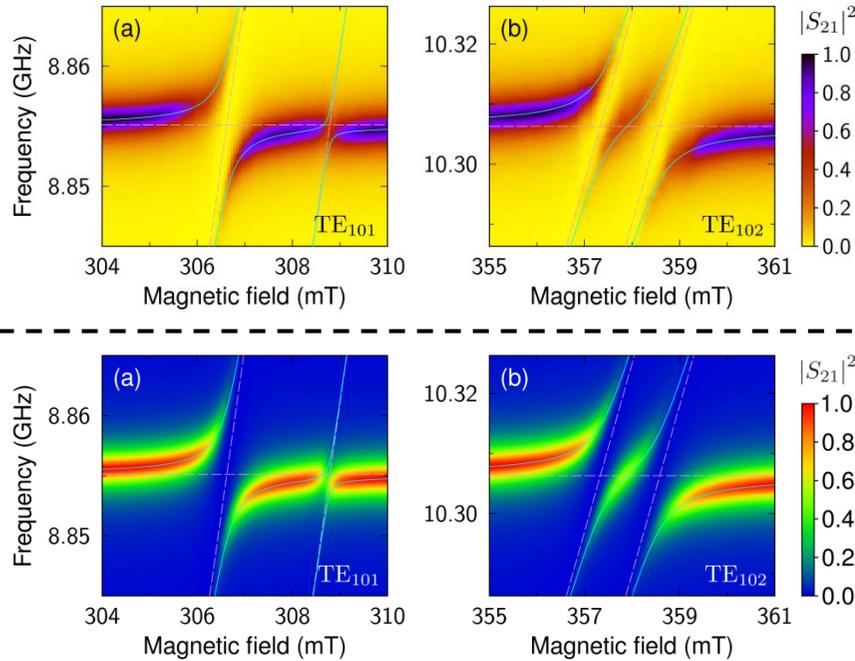
Magnon mode damping rate  $\gamma_m$  ( $m = \text{FMR}, \text{MS}$ )

**Strong coupling**  $\tilde{g}_m > \kappa_{\text{tot}}, \gamma_m$

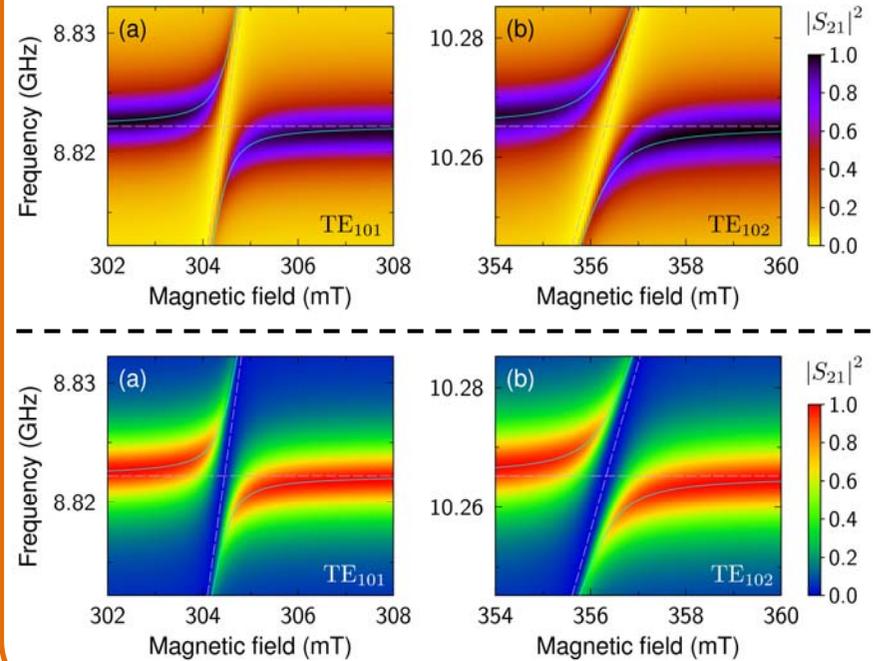
**Cooperativity**  $C \equiv \tilde{g}_m^2 / \kappa_{\text{tot}} \gamma_m > 1$

# Parameters extracted by fitting with the experiment

## Cryogenic temperature



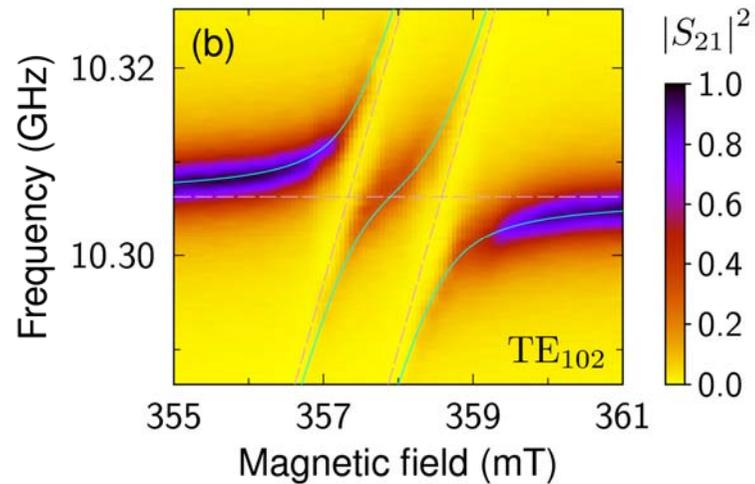
## Room temperature



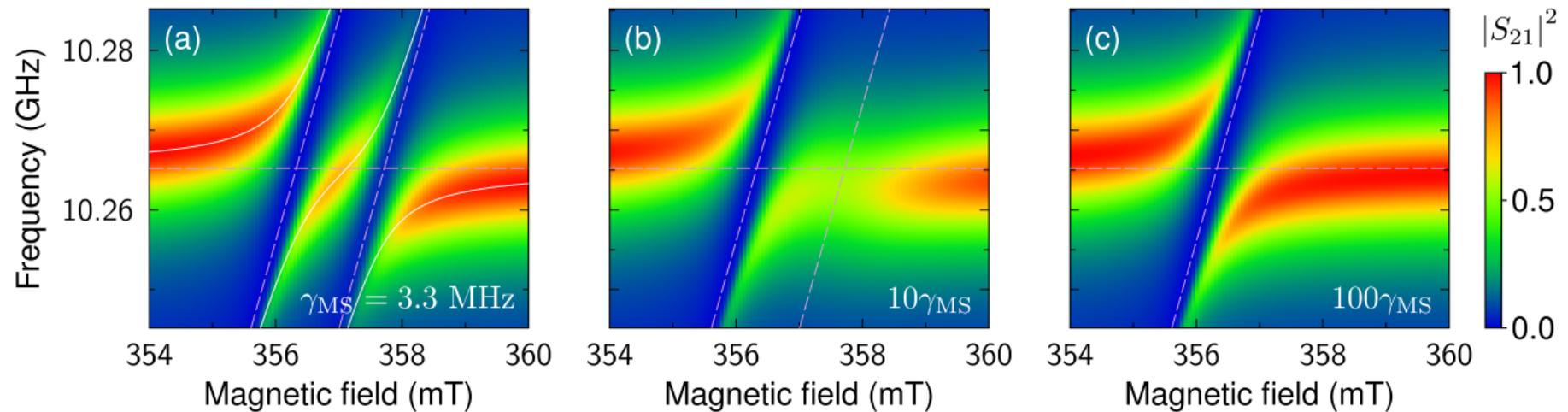
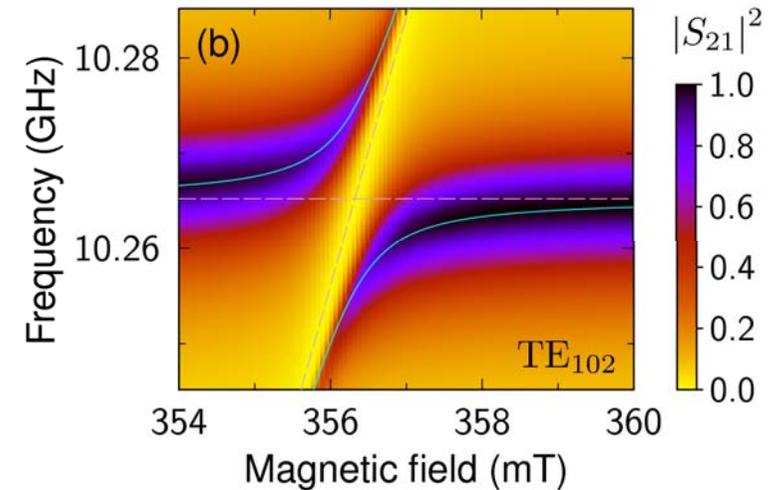
Temp.	Mode	$\tilde{g}_{\text{FMR}}, \tilde{g}_{\text{MS}}$ (/2 $\pi$ MHz)	$\kappa_{\text{tot}}$ (/2 $\pi$ MHz)	$\gamma_{\text{FMR}}, \gamma_{\text{MS}}$ (/2 $\pi$ MHz)	$C_{\text{FMR}}, C_{\text{MS}}$
Cryo. (22 mK)	TE <sub>101</sub>	5.1, 1.4	1.1	1.2, 2.7	22.1, 0.7
	TE <sub>102</sub>	7.5, 8.3	2.4	1.3, 3.3	18.0, 8.7
R. T. (300 K)	TE <sub>101</sub>	5.2	2.5	1.3	8.3
	TE <sub>102</sub>	9.6	5.9	1.5	10.4

# Comparison with simulations

## Cryogenic temperature



## Room temperature

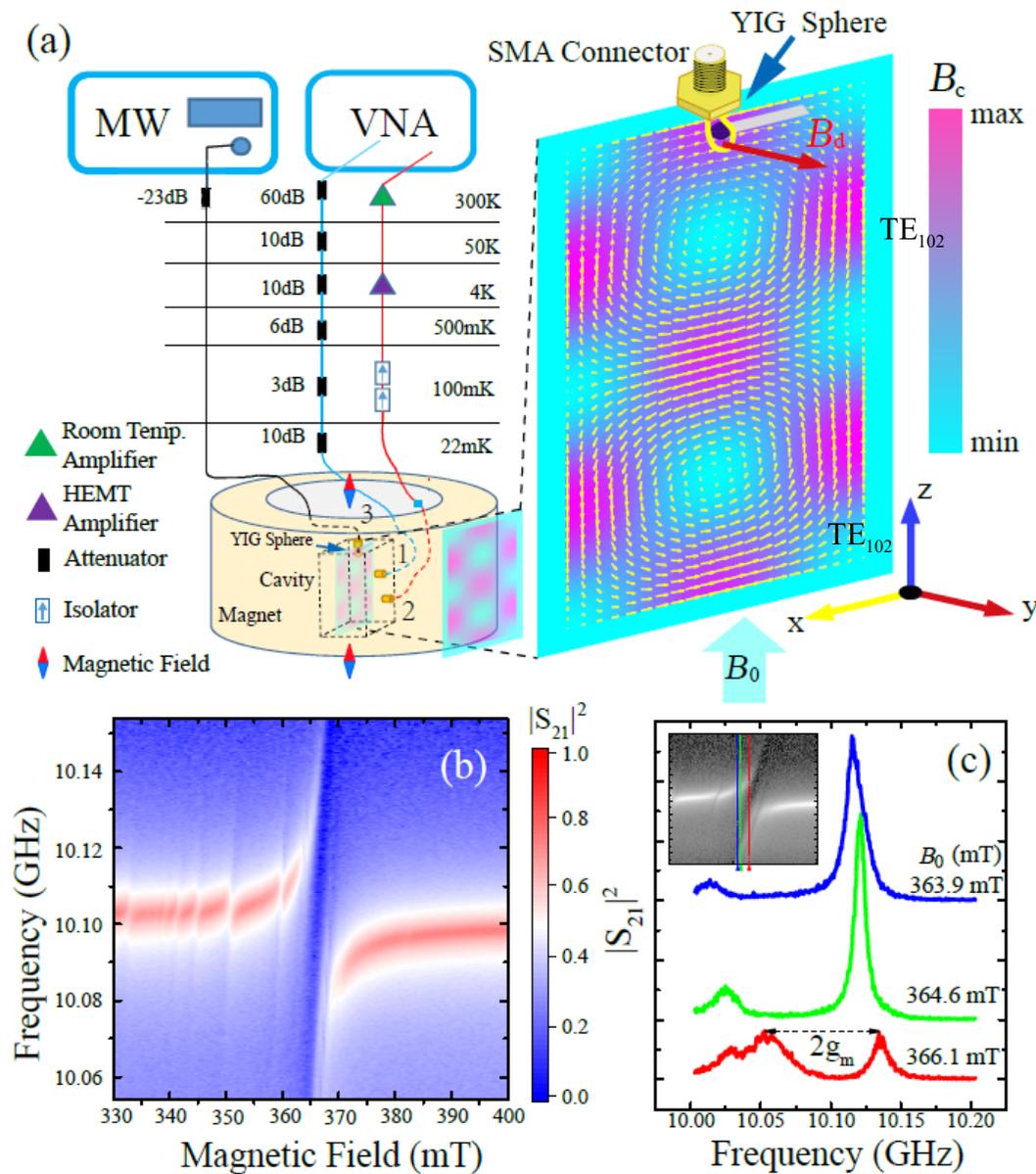


For the MS mode, the damping rate is increased two orders from Cryo. to R.T.

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# Experimental setup



Temperature: 22 mK

The YIG sphere is placed at the site where the magnetic field of the cavity mode TE<sub>102</sub> is maximal (as before).

Newly added:  
The YIG sphere can be directly pumped by a drive field.

# Cavity QED system with magnon Kerr effect

Total Hamiltonian:  $H = H_c + H_m + H_{\text{int}} + H_d.$

Here  $H_c = \omega_c a^\dagger a$      $H_m = -\gamma B_0 S_z - \frac{\mu_0 \gamma^2 K_{\text{an}}}{M^2 V_m} S_z^2$  ← Magnetocrystalline anisotropic energy

$$H_{\text{int}} = g_s (S^+ + S^-) (a^\dagger + a) \equiv 2g_s S_x (a^\dagger + a)$$

$$H_d = \Omega_s (S^+ + S^-) (e^{i\omega_d t} + e^{-i\omega_d t}) \equiv 4\Omega_s S_x \cos(\omega_d t)$$

Holstein-Primakoff transformation:

$$S^+ = \left( \sqrt{2S - b^\dagger b} \right) b, \quad S^- = b^\dagger \left( \sqrt{2S - b^\dagger b} \right), \quad S_z = S - b^\dagger b$$

For the low-lying excitations with  $\langle b^\dagger b \rangle / 2S \ll 1$ ,  $S^+ \approx b\sqrt{2S}$ ,  $S^- \approx b^\dagger\sqrt{2S}$ .

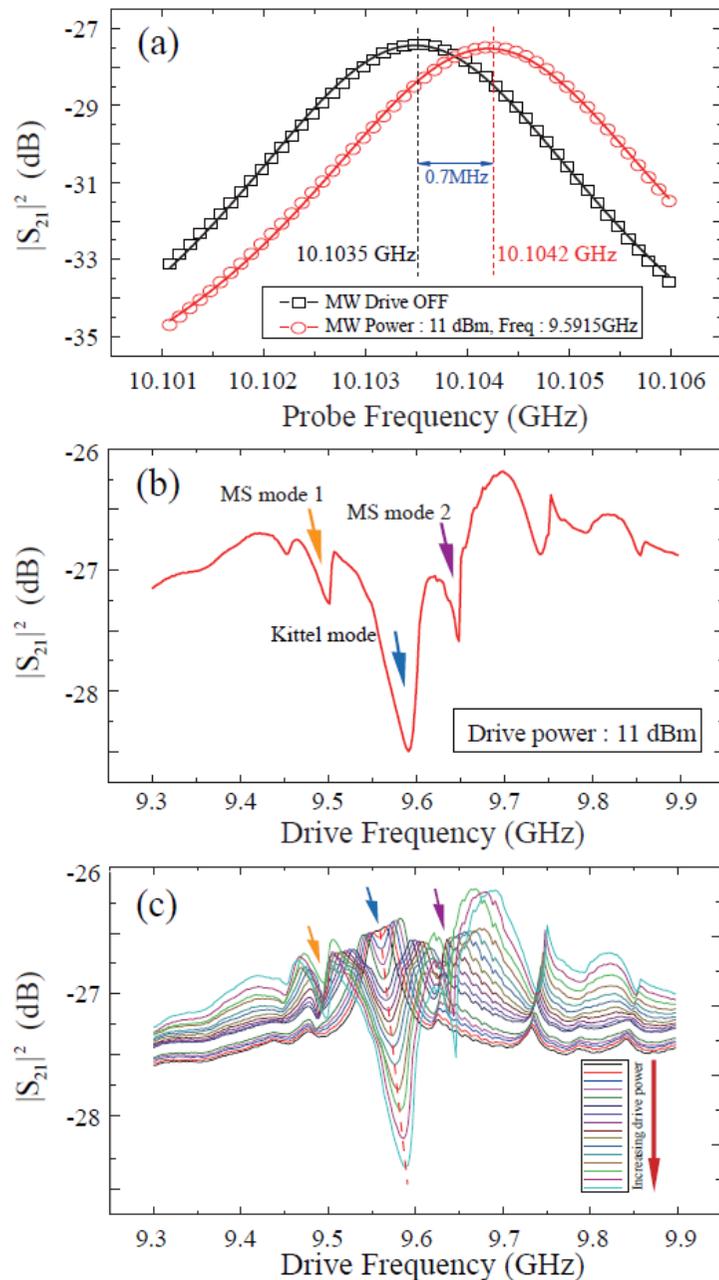
In the rotating-wave approximation, the total Hamiltonian of the coupled hybrid system becomes

$$H = \omega_c a^\dagger a + \omega_m b^\dagger b + K b^\dagger b b^\dagger b \\ + g_m (a^\dagger b + a b^\dagger) + \Omega_d (b^\dagger e^{-i\omega_d t} + b e^{i\omega_d t}),$$

$$K = \mu_0 K_{\text{an}} \gamma^2 / (M^2 V_m)$$

**Kerr effect of magnons** owing to the magnetocrystalline anisotropy.

# Kerr-effect-induced frequency shifts



In the **dispersive** regime, when considerable magnons are generated by a drive field, the effective Hamiltonian of the system is

$$H_{\text{eff}} = \left[ \omega_c + \frac{g_m^2}{\Delta} + \frac{2g_m^2}{\Delta^2} K \langle b^\dagger b \rangle \right] a^\dagger a + \left[ \omega_m - \frac{g_m^2}{\Delta} + \left( 1 - \frac{2g_m^2}{\Delta^2} \right) K \langle b^\dagger b \rangle \right] b^\dagger b + \Omega'_d (b^\dagger e^{-i\omega_d t} + b e^{i\omega_d t}),$$

$$\Omega'_d = \left[ 1 - \frac{1}{2(\omega_c - \omega_d)} \left( \frac{g_m^2}{\Delta} + \frac{2g_m^2}{\Delta^2} K \langle b^\dagger b \rangle \right) \right] \Omega_d,$$

where  $\Delta = \omega_c - \omega_m$ .

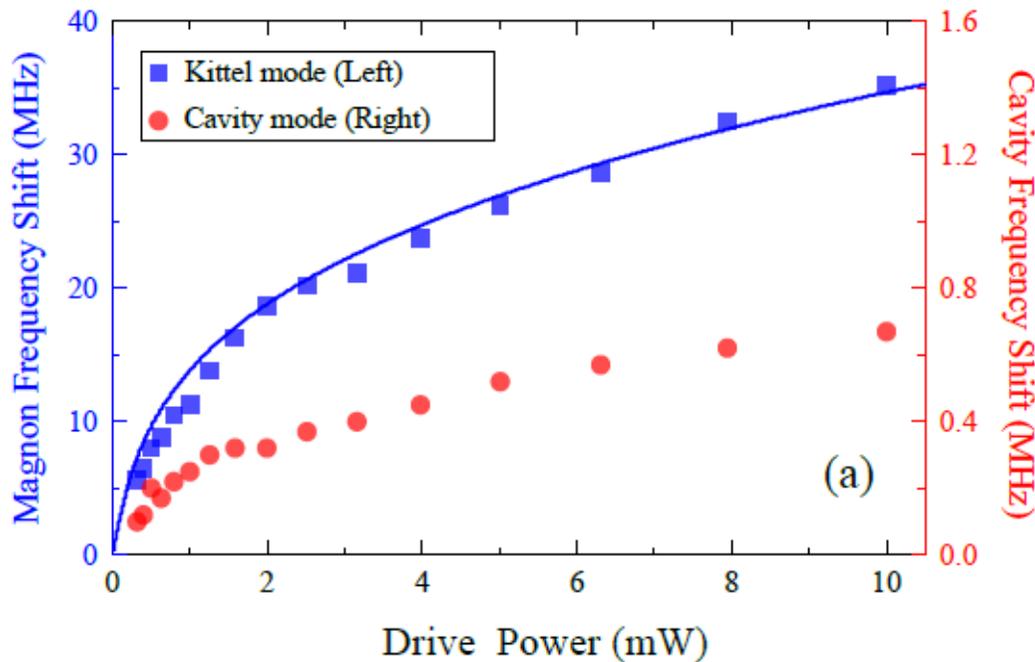
Kerr-effect-induced **shift of the central cavity frequency:**

$$\Delta_c = (2g_m^2/\Delta^2) K \langle b^\dagger b \rangle$$

Kerr-effect-induced **shift of the magnon frequency:**

$$\Delta_m = (1 - 2g_m^2/\Delta^2) K \langle b^\dagger b \rangle \approx K \langle b^\dagger b \rangle$$

# Kerr-effect-induced frequency shifts vs. Drive power



Kerr-effect-induced shift of the central cavity frequency:

$$\Delta_c = (2g_m^2/\Delta^2)K\langle b^\dagger b \rangle$$

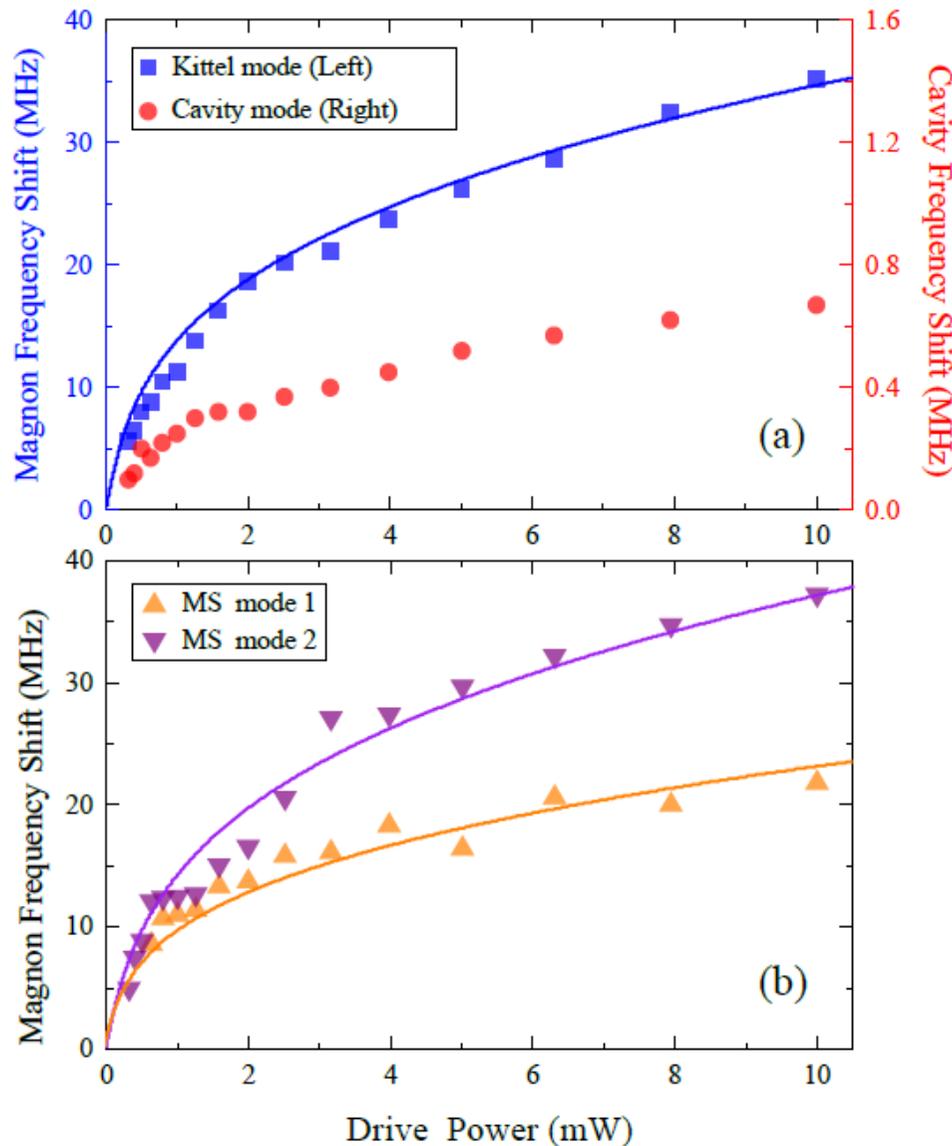
Kerr-effect-induced shift of the magnon frequency:

$$\Delta_m = (1 - 2g_m^2/\Delta^2)K\langle b^\dagger b \rangle \approx K\langle b^\dagger b \rangle$$

➡  $\Delta_c \propto \Delta_m$

Therefore, it is predicted that the **shift of the central cavity frequency** should have a *similar behavior* as the frequency shift of the **Kittel mode**.

# Magnon frequency shift vs. Drive power



Using a Langevin approach, we obtain *the relation between the magnon frequency shift and the drive power*.

$$\left[ \Delta_m^2 + \left( \frac{\gamma_m}{2} \right)^2 \right] \Delta_m - cP = 0$$

- The experimental results for the Kittel mode agree well with this relation (derived for an uniformly magnetized YIG sphere).
- The experimental results for the MS modes deviate from this relation, which confirms the deviations of the MS modes from homogeneous magnetization.

# Conclusions

Experimental study of a cavity QED with magnons at both cryogenic and room temperatures.

- Robustness of the FMR mode against temperature.
- A drastic increase of the damping rate of the MS mode from cryogenic to room temperature.

Experimental demonstration of the magnon Kerr effect in a cavity QED system

- Kerr-effect-induced shift of the central cavity mode & Kerr-effect-induced magnon frequency shift.
- The experimental results for the Kittel mode agree well with the analytical relation between the magnon frequency shift and the drive power.

# Outlook: Quantum magnonics & solid-state hybrids

- Ultrastrong coupling of magnons to photons
- Tim-domain properties of magnons
- Demonstrate non-classical states of magnon mode
- Coupling with qubits (*single transmon qubit, U. of Tokyo/RIKEN group, Science, 2015*).
- More nonlinear effects (bistability, chaos)
- BEC with magnons?

## Main Collaborators:

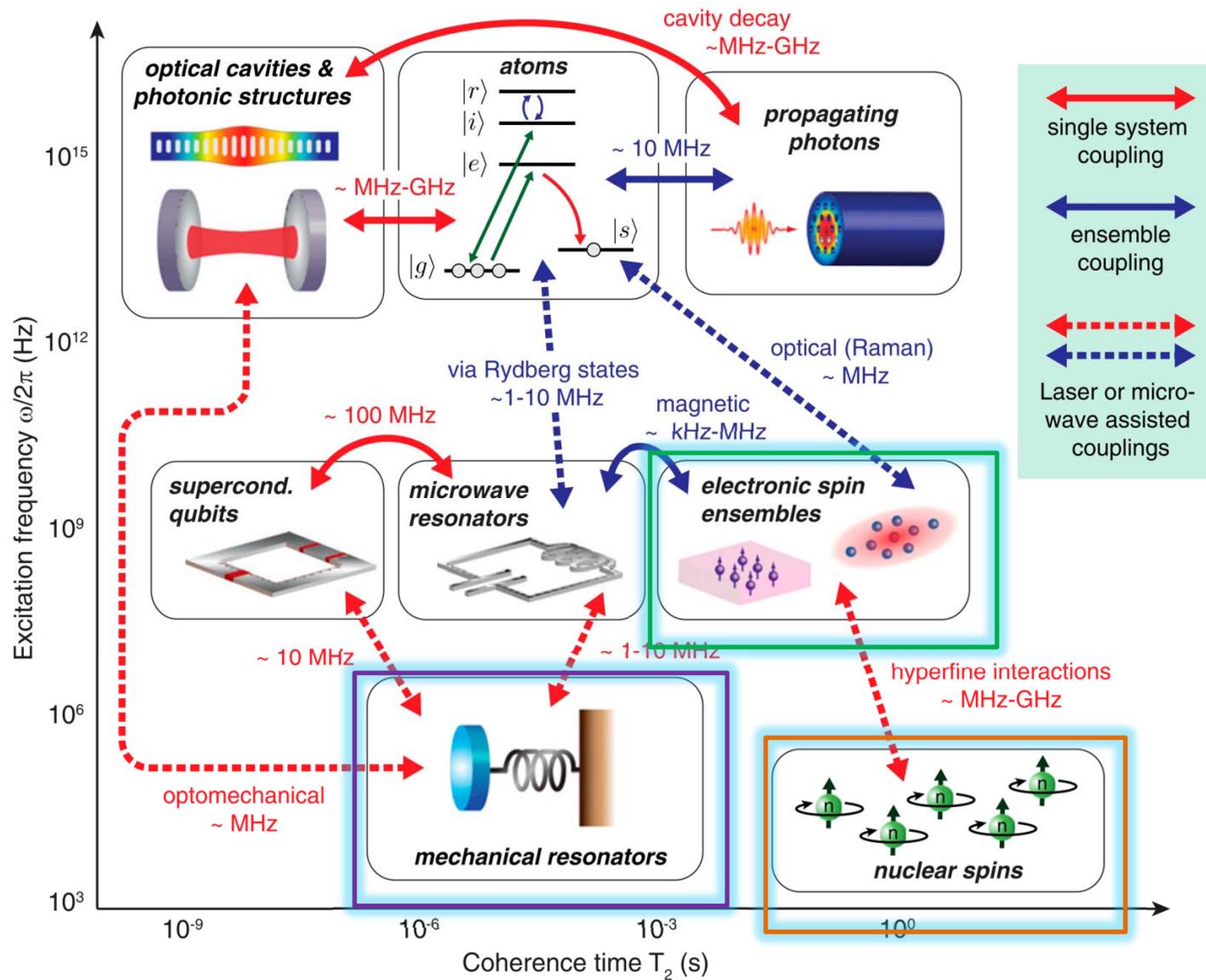
CSRC: Dengke Zhang, Yi-Pu Wang (experimental) & Guo-Qiang Zhang (theoretical)

Tie-Fu Li (Tsinghua & CSRC), Franco Nori (*RIKEN/Univ. of Michigan*)

& Can-Ming Hu (Univ. of Manitoba)

Thank you  
for your attention!

# Hybrid quantum systems



Kurizki *et al.*, "Quantum Technologies with Hybrid Systems" *PNAS*, 2015;  
 Xiang *et al.*, "Hybrid quantum circuits", *Rev. Mod. Phys.* 2013.

# Magnetic resonance

Resonant frequency

Nuclear magnetic resonance (NMR)

Radio-frequency

## Paramagnetic materials

Electron paramagnetic resonance (EPR)

Microwave

---Electron spin resonance (ESR)

## Ferromagnetic material

Ferromagnetic resonance (FMR)

Microwave

Ferrimagnetic resonance (FiMR)

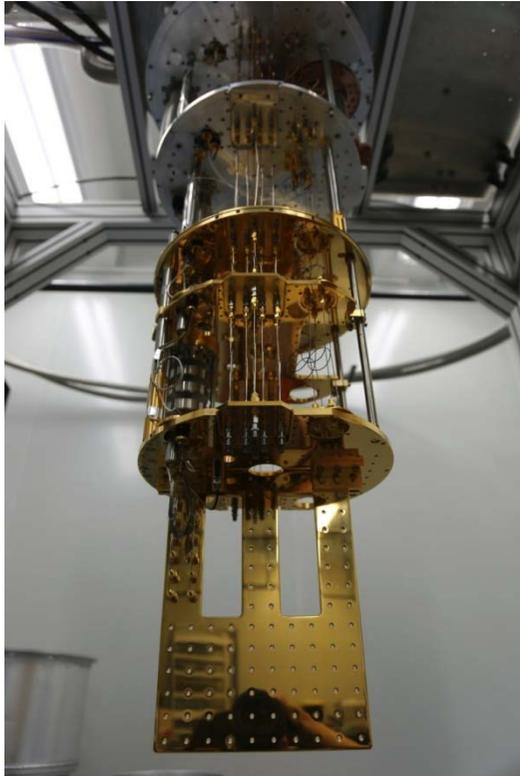
Microwave — Far infrared

Antiferromagnetic resonance (AFMR)

Microwave — Far infrared

# Measurement setups

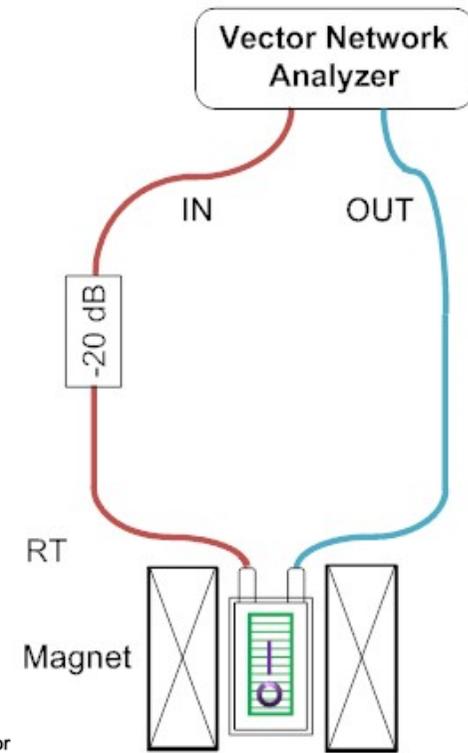
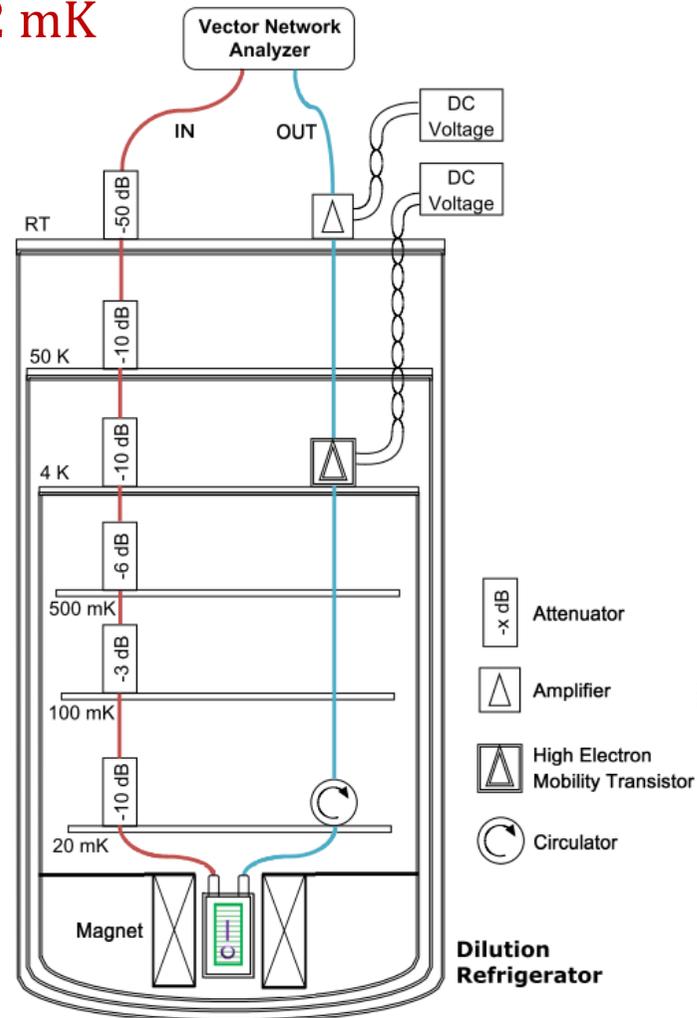
Cryogenic  
temperature  
(Cryo.)  
22 mK



Bluefors cryo-free dilution  
refrigerator:

Base temperature 7mK

Room temperature (R. T.)  
300 K

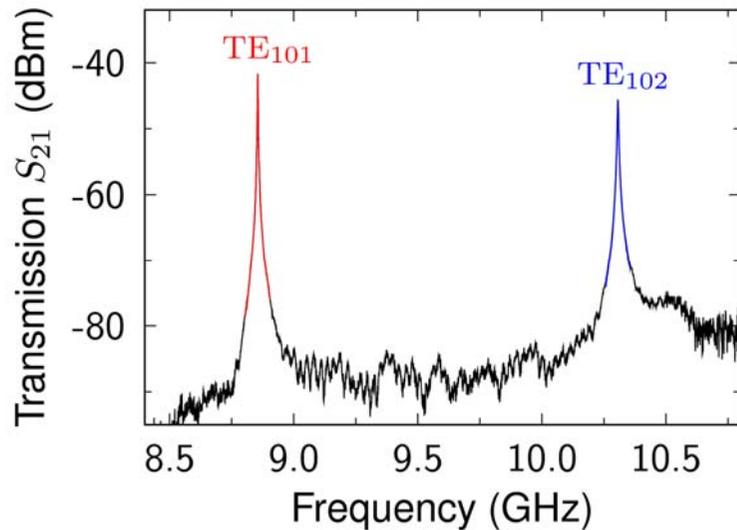


# $S_{21}$ spectra of cavity without a YIG sphere

$$\text{TE}_{101} : \kappa_{i,1}/2\pi = 0.19 \text{ MHz}; \quad \kappa_{o,1}/2\pi = 0.20 \text{ MHz}$$

$$\text{TE}_{102} : \kappa_{i,2}/2\pi = 0.85 \text{ MHz}; \quad \kappa_{o,2}/2\pi = 0.99 \text{ MHz}$$

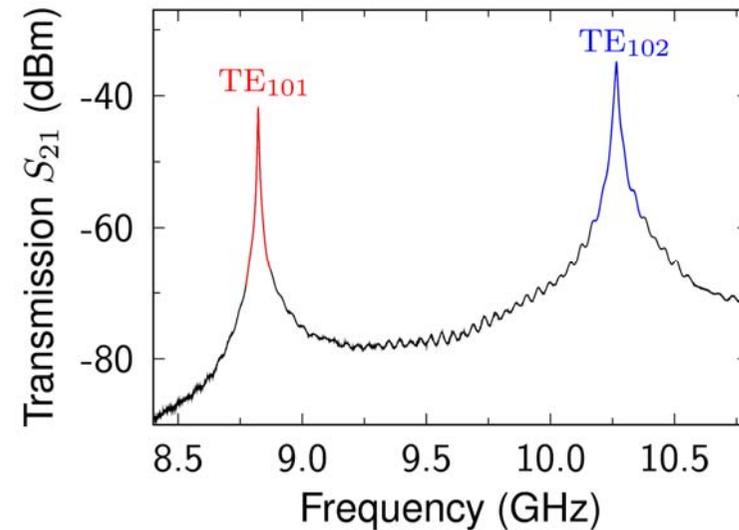
## Cryogenic temperature (22 mK)



$$\kappa_{\text{tot},1}/2\pi = 1.1 \text{ MHz}; \quad \kappa_{\text{int},1}/2\pi = 0.71 \text{ MHz}$$

$$\kappa_{\text{tot},2}/2\pi = 2.4 \text{ MHz}; \quad \kappa_{\text{int},2}/2\pi = 0.56 \text{ MHz}$$

## Room temperature



$$\kappa_{\text{tot},1}/2\pi = 2.5 \text{ MHz}; \quad \kappa_{\text{int},1}/2\pi = 2.11 \text{ MHz}$$

$$\kappa_{\text{tot},2}/2\pi = 5.9 \text{ MHz}; \quad \kappa_{\text{int},2}/2\pi = 4.06 \text{ MHz}$$

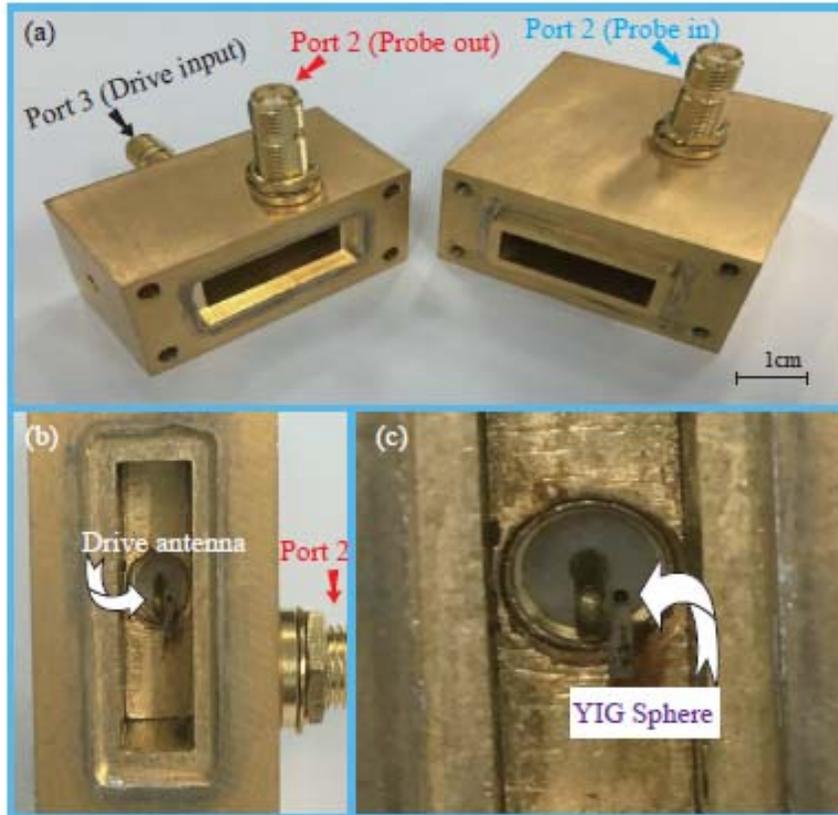
# Experimental parameters

Temp.	Mode	Input power (dBm)	Input photon number	Thermal photon number
Cryo. (22 mK)	TE <sub>101</sub>	-130	0.8	$\sim 1 \times 10^{-2}$
	TE <sub>102</sub>	-100	800	$\sim 1 \times 10^{-2}$
R. T. (300 K)	TE <sub>101</sub>	-130	0.7	$\sim 1 \times 10^{-2}$
	TE <sub>102</sub>	-100	656	$\sim 1 \times 10^{-2}$
	TE <sub>101</sub>	-20	$1.8 \times 10^{10}$	705
	TE <sub>102</sub>	-20	$1.1 \times 10^{10}$	606

Quantum limit is reached at ~22 mK.

Question: Strong magnon-photon coupling was reached @ R. T. for the MS mode?

# The 3D cavity containing a YIG sample



The 3D cavity has inner dimensions:  
 $44.0 \times 20.0 \times 6.0 \text{ mm}^3$

The diameter of the YIG sphere: 1mm