# LHV models for quantum states and measurements

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# Joint work with



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#### Quantum Nonlocality



Data:  $p(ab|xy) = \operatorname{Tr}(\varrho M_{a|x} \otimes M_{b|y})$ 







# Q1: LHV model for an entangled quantum state

All 
$$\leftarrow \rho$$
 All Measurements All





## **Question 1**

# General method for constructing LHV models for entangled quantum states

Hirsch et al. arxiv 2015 also: Cavalcanti et al. arxiv 2015

# **Question 1**

# General method for constructing LHV models for entangled quantum states

- Applicable to any entangled state
- Can be implemented on a standard computer
- Converging sequence of tests

Hirsch et al. arxiv 2015 also: Cavalcanti et al. arxiv 2015

$$\rho_W^{\mu} = \mu \left| \phi_+ \right\rangle \left\langle \phi_+ \right| + (1 - \mu) \frac{\mathbf{1}}{4}$$



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# What about a generic state?

Map the problem to a simpler one

LHV model for ALL measurements on  $\rho$ 



LHV model for a **FINITE** set of measurements on  $\rho'$  (close to  $\rho$ )

Initial state  $\rho'$ 

Take finite sets of meas  $\{M_{a|x}\}$  and  $\{M_{b|y}\}$ 





#### Initial state $\rho'$

Take finite sets of meas  $\{M_{a|x}\}$  and  $\{M_{b|y}\}$ 

Check  $p(a,b|xy) = tr(\rho' M_{a|x} M_{b|y})$  is local



Initial state  $\rho'$ 

Take finite sets of meas  $\{M_{a|x}\}$  and  $\{M_{b|y}\}$ 

Check  $p(a,b|xy) = tr(\rho' M_{a|x} M_{b|y})$  is local



 $\boldsymbol{\rho}'$  is local for all **noisy measurements** 

$$M^{\eta}_{\pm|\vec{x}} = \frac{1}{2} (\mathbf{1} \pm \eta \, \hat{x} \cdot \vec{\sigma})$$

Initial state  $\rho'$ 

Take finite sets of meas  $\{M_{a|x}\}$  and  $\{M_{b|y}\}$ Check  $p(a,b|xy) = tr(\rho' M_{a|x} M_{b|y})$  is local

η

 $\rho'$  is local for all noisy measurements  $M^{\eta}_{\pm|\vec{x}} = \frac{1}{2}(\mathbf{1} \pm \eta \, \hat{x} \cdot \vec{\sigma})$ 

Noisy version of  $\rho'$  is local for all pure meas <u>Target state</u>  $\rho = \eta^2 \rho + (1-\eta^2) \sigma_{sep}$ 

### **Application: Werner states**



$$\rho_W^{\mu} = \mu \left| \phi_+ \right\rangle \left\langle \phi_+ \right| + (1 - \mu) \frac{\mathbf{1}}{4}$$





## Iterative procedure



## To be explored

- POVMs
- Higher dimensions
- Multipartite systems

## Question 2

# Can all incompatible sets of measurements lead to Bell violation?



#### Incompatibility $\rightarrow$ Bell violation ??



Incompatibility  $\rightarrow$  Bell violation ??

**Projective measurements** 

Incompatibility (commutativity) → CHSH violation Khalfin Tsirelson'85



Incompatibility  $\rightarrow$  Bell violation ??

#### **Projective measurements**

Incompatibility (commutativity) → CHSH violation Khalfin Tsirelson'85

What about POVMs?

Joint measurability

POVM {M<sub>a</sub>} & {M<sub>b</sub>} are JM if there is joint POVM {C<sub>ab</sub>} s.t.  $M_a = \Sigma_b C_{ab} \& M_b = \Sigma_a C_{ab}$  Joint measurability

POVM {M<sub>a</sub>} & {M<sub>b</sub>} are JM if there is joint POVM {C<sub>ab</sub>} s.t.  $M_a = \Sigma_b C_{ab} \& M_b = \Sigma_a C_{ab}$ 



Partial JM does not imply fully JM



Heinossari'08, Liang et al'11

#### JM vs nonlocality



2 binary POVMs Wolf et al. 2007
 Incompatibility → CHSH violation

#### JM vs nonlocality



2 binary POVMs Wolf et al. 2007
 Incompatibility → CHSH violation

Incompatibility steering
 Quintino et al & Uola et al '14



#### Our candidate

Continuous set of noisy qubit measurements

 $\mathcal{M} = \{M_{\pm|\hat{x}}^{\eta}\}$  for all Bloch vectors  $\vec{x}$ 

where 
$$M^{\eta}_{\pm | \vec{x}} = \frac{1}{2} (\mathbf{1} \pm \eta \, \hat{x} \cdot \vec{\sigma})$$

#### Our candidate

Continuous set of noisy qubit measurements





### Our candidate

Continuous set of noisy qubit measurements











To show:  $p(ab|xy) = \operatorname{tr}(|\phi_{\theta}\rangle \langle \phi_{\theta}| M^{\eta}_{a|\hat{x}} \otimes \Pi_{b|\hat{y}})$  is local

Step 2  

$$\mathcal{M} \longleftarrow |\phi_{\theta}\rangle \longrightarrow \text{All projective meas }\Pi_{b|y}$$
To show:  $p(ab|xy) = \operatorname{tr}(|\phi_{\theta}\rangle \langle \phi_{\theta}| M_{a|\hat{x}}^{\eta} \otimes \Pi_{b|\hat{y}})$  is local  
 $\operatorname{tr}(|\phi_{\theta}\rangle \langle \phi_{\theta}| M_{a|\hat{x}}^{\eta} \otimes \Pi_{b|\hat{y}}) = \operatorname{tr}(\rho_{\theta}^{\eta} \Pi_{a|\hat{x}} \otimes \Pi_{b|\hat{y}})$ 

$$\operatorname{tr}(|\phi_{\theta}\rangle \langle \phi_{\theta}| M_{a|\hat{x}}^{\eta} \otimes \Pi_{b|\hat{y}}) = \operatorname{tr}(\rho_{\theta}^{\eta} \Pi_{a|\hat{x}} \otimes \Pi_{b|\hat{y}})$$
noisy meas
pure state
 $\rho_{\theta}^{\eta} = \eta |\phi_{\theta}\rangle \langle \phi_{\theta}| + (1 - \eta) \frac{\pi}{2} \otimes \rho_{B}$ 





### Step 3

Show that 
$$\rho_{\theta}^{\eta} = \eta \ket{\phi_{\theta}} \langle \phi_{\theta} | + (1 - \eta) \frac{\mathbf{1}}{2} \otimes \rho_{B}$$

is local for  $\eta > \frac{1}{2}$  and for all  $\theta \in [0, \pi/4]$ .



### Step 3

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Show that 
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#### Summary

LHV model for a set of incompatible POVMs

$$\mathcal{M}^{\eta}_{A} = \{M^{\eta}_{\pm|\hat{x}}\}$$
 for all Bloch vectors  $\vec{\mathbf{x}}$ 



#### Summary

#### LHV model for a set of incompatible POVMs





#### Open questions

• LHV model for set of few measurements



- Notion of incompatibility corresponding to Bell nonlocality
- Activation