

Postquantum steering

Ana Belén Sainz, Nicolas Brunner, Daniel Cavalcanti
Paul Skrzypczyk and Tamás Vértesi

Phys. Rev. Lett. 115, 190403 (2015)

Steering

Alice

Bob

ρ_A



Steering

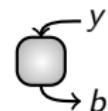
Alice

Bob

$$\rho_A \quad \bullet$$

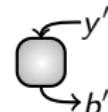


$$\sigma_{b|y}^A \quad \bullet$$



$$\sigma_{b'|y'}^A \quad \bullet$$

:



Steering

Fix y \longrightarrow ensemble: $\{\sigma_{b|y}^A\}_b$, \longrightarrow $\rho_A = \sum_b \sigma_{b|y}^A$

Assemblage: $\{\sigma_{b|y}^A\}_{b,y}$. $p(b|y) = \text{tr}(\sigma_{b|y}^A)$.

Steering

Fix y \longrightarrow ensemble: $\{\sigma_{b|y}^A\}_b$, \longrightarrow $\rho_A = \sum_b \sigma_{b|y}^A$

Assemblage: $\{\sigma_{b|y}^A\}_{b,y}$. $p(b|y) = \text{tr}(\sigma_{b|y}^A)$.

Quantum: $\sigma_{b|y}^A = \text{tr}_B (\mathbb{1}_A \otimes M_{b|y} \rho_{AB})$

Given an assemblage, **could it have a classical explanation?**

Steering

Fix $y \rightarrow$ ensemble: $\{\sigma_{b|y}^A\}_b, \rightarrow \rho_A = \sum_b \sigma_{b|y}^A$

Assemblage: $\{\sigma_{b|y}^A\}_{b,y}, p(b|y) = \text{tr}(\sigma_{b|y}^A)$.

Quantum: $\sigma_{b|y}^A = \text{tr}_B (\mathbb{1}_A \otimes M_{b|y} \rho_{AB})$

Given an assemblage, **could it have a classical explanation?**

Set of “classical” assemblages

Steering Inequality:
 $\text{tr} \sum_{by} F_{by} \sigma_{b|y}^A \leq \beta_{US}$

Steering

Fix $y \rightarrow$ ensemble: $\{\sigma_{b|y}^A\}_b, \rightarrow \rho_A = \sum_b \sigma_{b|y}^A$

Assemblage: $\{\sigma_{b|y}^A\}_{b,y}, p(b|y) = \text{tr}(\sigma_{b|y}^A)$.

Quantum: $\sigma_{b|y}^A = \text{tr}_B (\mathbb{1}_A \otimes M_{b|y} \rho_{AB})$

Given an assemblage, **could it have a classical explanation?**

Set of “classical” assemblages

Steering Inequality:
 $\text{tr} \sum_{by} F_{by} \sigma_{b|y}^A \leq \beta_{US}$

Given an assemblage, **could it have a quantum explanation?**

Bipartite steering

Given $\{\sigma_{b|y}^A\}_{b,y}$, $\rho_A = \sum_b \sigma_{b|y}^A$, $\text{tr}(\rho_A) = 1$

$\exists \rho_{AB}$, $\{M_{b|y}\}_{b,y}$ st $\sigma_{b|y}^A = \text{tr}_B (\mathbb{1}_A \otimes M_{b|y} \rho_{AB})$

¹N. Gisin, Helvetica Physica Acta 62, 363 (1989).

L. P. Hughston, R. Jozsa and W. K. Wootters, Phys. Lett. A 183, 14 (1993).

Bipartite steering

Given $\{\sigma_{b|y}^A\}_{b,y}$, $\rho_A = \sum_b \sigma_{b|y}^A$, $\text{tr}(\rho_A) = 1$

$\exists \rho_{AB}$, $\{M_{b|y}\}_{b,y}$ st $\sigma_{b|y}^A = \text{tr}_B (\mathbb{1}_A \otimes M_{b|y} \rho_{AB})$

- Alice and Bob: Yes ! GHJW theorem¹
- Multipartite scenarios?

¹N. Gisin, Helvetica Physica Acta 62, 363 (1989).

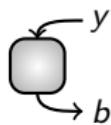
L. P. Hughston, R. Jozsa and W. K. Wootters, Phys. Lett. A 183, 14 (1993).

Steering: multipartite scenarios

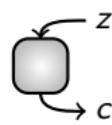
Alice

$$\bullet$$
$$\sigma_{bc|yz}^A$$

Bob



Charlie

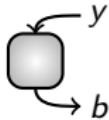


Steering: multipartite scenarios

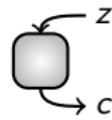
Alice

$$\bullet \\ \sigma_{bc|yz}^A$$

Bob



Charlie



Fix $y, z \rightarrow$ ensemble: $\{\sigma_{bc|yz}^A\}_{b,c}, \rightarrow \rho_A = \sum_{b,c} \sigma_{bc|yz}^A$

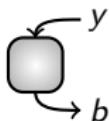
Assemblage: $\{\sigma_{bc|yz}^A\}_{b,y,c,z} \cdot p(bc|yz) = \text{tr}(\sigma_{bc|yz}^A)$.

Steering: multipartite scenarios

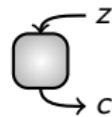
Alice

$$\bullet \\ \sigma_{bc|yz}^A$$

Bob



Charlie



Fix $y, z \rightarrow$ ensemble: $\{\sigma_{bc|yz}^A\}_{b,c}, \rightarrow \rho_A = \sum_{b,c} \sigma_{bc|yz}^A$

Assemblage: $\{\sigma_{bc|yz}^A\}_{b,y,c,z} \cdot p(bc|yz) = \text{tr}(\sigma_{bc|yz}^A)$.

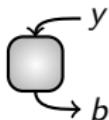
No Signalling: $\sum_b \sigma_{bc|yz}^A = \sum_b \sigma_{bc|y'z}^A, \quad \sum_c \sigma_{bc|yz}^A = \sum_c \sigma_{bc|yz'}^A$

Steering: multipartite scenarios

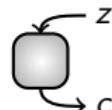
Alice

$$\bullet \\ \sigma_{bc|yz}^A$$

Bob



Charlie



Fix $y, z \rightarrow$ ensemble: $\{\sigma_{bc|yz}^A\}_{b,c}, \rightarrow \rho_A = \sum_{b,c} \sigma_{bc|yz}^A$

Assemblage: $\{\sigma_{bc|yz}^A\}_{b,y,c,z} \quad p(bc|yz) = \text{tr}(\sigma_{bc|yz}^A)$.

No Signalling: $\sum_b \sigma_{bc|yz}^A = \sum_b \sigma_{bc|y'z}^A, \quad \sum_c \sigma_{bc|yz}^A = \sum_c \sigma_{bc|yz'}^A$

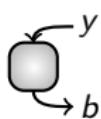
$\exists \rho_{ABC}, \{M_{b|y}\}_{b,y}, \{M_{c|z}\}_{c,z}$ st $\sigma_{bc|yz}^A = \text{tr}_B (\mathbb{1}_A \otimes M_{b|y} \otimes M_{c|z} \rho_{ABC})$

Postquantum steering: example

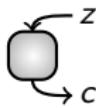
Alice

$$\bullet \\ \sigma_{bc|yz}^A$$

Bob



Charlie



$$b, c, y, z \in \{0, 1\}$$

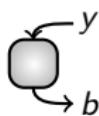
$$\rho_A = \frac{1}{2}.$$

Postquantum steering: example

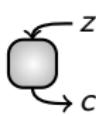
Alice

$$\sigma_{bc|yz}^A \bullet$$

Bob



Charlie



$$b, c, y, z \in \{0, 1\}$$

$$\rho_A = \frac{1}{2}.$$

- $(y, z) = (0, 0), (0, 1), (1, 0)$:

$$\sigma_{bc|yz}^A = \begin{cases} \frac{1}{4}, & \text{if } b = c, \\ 0, & \text{if } b \neq c, \end{cases}$$

- $(y, z) = (1, 1)$:

$$\sigma_{bc|yz}^A = \begin{cases} 0, & \text{if } b = c, \\ \frac{1}{4}, & \text{if } b \neq c, \end{cases}$$

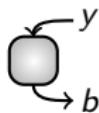
$$\sum_b \sigma_{bc|yz}^A = \frac{1}{4}, \quad \sum_c \sigma_{bc|yz}^A = \frac{1}{4}$$

Postquantum steering: example

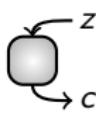
Alice

$$\bullet \\ \sigma_{bc|yz}^A$$

Bob



Charlie



$$b, c, y, z \in \{0, 1\}$$

$$\rho_A = \frac{1}{2}.$$

- $(y, z) = (0, 0), (0, 1), (1, 0)$:

$$\sigma_{bc|yz}^A = \begin{cases} \frac{1}{4}, & \text{if } b = c, \\ 0, & \text{if } b \neq c, \end{cases}$$

$$\begin{aligned} p(bc|yz) &= \text{tr} \left(\sigma_{bc|yz}^A \right) \\ &\stackrel{\text{quantum}}{=} \text{tr}_{BC} \left(M_{b|y} \otimes M_{c|z} \rho_{BC} \right) \end{aligned}$$

- $(y, z) = (1, 1)$:

$$\sigma_{bc|yz}^A = \begin{cases} 0, & \text{if } b = c, \\ \frac{1}{4}, & \text{if } b \neq c, \end{cases}$$

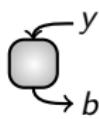
$$\sum_b \sigma_{bc|yz}^A = \frac{1}{4}, \quad \sum_c \sigma_{bc|yz}^A = \frac{1}{4}$$

Postquantum steering: example

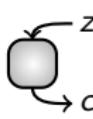
Alice

$$\bullet \\ \sigma_{bc|yz}^A$$

Bob



Charlie



$$b, c, y, z \in \{0, 1\}$$

$$\rho_A = \frac{1}{2}.$$

- $(y, z) = (0, 0), (0, 1), (1, 0)$:

$$\sigma_{bc|yz}^A = \begin{cases} \frac{1}{4}, & \text{if } b = c, \\ 0, & \text{if } b \neq c, \end{cases}$$

$$\begin{aligned} p(bc|yz) &= \text{tr} \left(\sigma_{bc|yz}^A \right) \\ &\stackrel{\text{quantum}}{=} \text{tr}_{BC} \left(M_{b|y} \otimes M_{c|z} \rho_{BC} \right) \end{aligned}$$

- $(y, z) = (1, 1)$:

$$\sigma_{bc|yz}^A = \begin{cases} 0, & \text{if } b = c, \\ \frac{1}{4}, & \text{if } b \neq c, \end{cases}$$

$$p(bc|yz) = \begin{cases} \frac{1}{2}, & \text{if } b \oplus c = yz, \\ 0, & \text{otherwise.} \end{cases}$$

$$\sum_b \sigma_{bc|yz}^A = \frac{1}{4}, \quad \sum_c \sigma_{bc|yz}^A = \frac{1}{4}$$

No quantum realisation for
the assemblage

Postquantum steering: example without postquantum nonlocality

- (1) Postquantum assemblage $\{\sigma_{bc|yz}^A\}_{b,y,c,z}$
- (2) Quantum correlations for every measurement by Alice:

$$p(abc|xyz) = \text{tr} (M_{a|x} \otimes \sigma_{bc|yz}^A)$$

(1) Postquantum assemblage $\sigma_{bc|yz}^A$

Steering inequality: F_{bcyz}

$$\text{tr} \sum_{bcyz} F_{bcyz} \sigma_{bc|yz}^A$$

(1) Postquantum assemblage $\sigma_{bc|yz}^A$

Steering inequality: F_{bcyz}

$$\text{tr} \sum_{bcyz} F_{bcyz} \sigma_{bc|yz}^A \leq \beta_Q$$

(1) Postquantum assemblage $\sigma_{bc|yz}^A$

Steering inequality: F_{bcyz}

$$\text{tr} \sum_{bcyz} F_{bcyz} \sigma_{bc|yz}^A \leq \beta_Q$$

How to compute β_Q ?

(1) Postquantum assemblage $\sigma_{bc|yz}^A$

Steering inequality: F_{bcyz}

$$\text{tr} \sum_{bcyz} F_{bcyz} \sigma_{bc|yz}^A \leq \beta_Q$$

How to compute β_Q ? \rightarrow upper bound

(1) Postquantum assemblage $\sigma_{bc|yz}^A$

Steering inequality: F_{bcyz}

$$\text{tr} \sum_{bcyz} F_{bcyz} \sigma_{bc|yz}^A \leq \beta_Q$$

How to compute β_Q ? \rightarrow upper bound

Almost quantum assemblages: $\tilde{Q} \supset Q$

$$\beta_{\tilde{Q}} \geq \beta_Q$$

(1) Postquantum assemblage $\sigma_{bc|yz}^A$

Steering inequality: F_{bcyz}

$$\mathrm{tr} \sum_{bcyz} F_{bcyz} \sigma_{bc|yz}^A \leq \beta_Q$$

How to compute β_Q ? \rightarrow upper bound

Almost quantum assemblages: $\tilde{Q} \supset Q$

$$\beta_{\tilde{Q}} \geq \beta_Q$$

$$\mathrm{tr} \sum_{bcyz} F_{bcyz} \sigma_{bc|yz}^A > \beta_{\tilde{Q}} \Rightarrow \sigma_{bc|yz}^A \text{ is postquantum}$$

Example without postquantum nonlocality

- (1) Postquantum assemblage $\{\sigma_{bc|yz}^A\}_{b,y,c,z}$
- (2) Quantum correlations for every measurement by Alice:

$$p(abc|xyz) = \text{tr} (M_{a|x} \otimes \sigma_{bc|yz}^A)$$

(2) Quantum correlations $p(abc|xyz)$

- (i) $p(abc|xyz)$ is local
- (ii) Real qubit assemblage, local for all projective measurements
- (iii) Qutrit assemblage, local for all POVMs².

²F. Hirsch, M. T. Quintino, J. Bowles and N. Brunner, Phys. Rev. Lett, 111, 160402 (2013).

(ii) Qubit assemblage, local for all PVM

$$\Pi_{a|x}(\mu) = \mu \Pi_{a|x} + (1 - \mu) \mathbb{1}/2, \quad \sigma_{bc|yz}^A(\mu) = \mu \sigma_{bc|yz}^A + (1 - \mu) \text{tr}(\sigma_{bc|yz}^A) \mathbb{1}/2$$

(ii) Qubit assemblage, local for all PVM

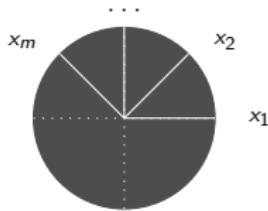
$$\Pi_{a|x}(\mu) = \mu \Pi_{a|x} + (1 - \mu) \mathbb{1}/2, \quad \sigma_{bc|yz}^A(\mu) = \mu \sigma_{bc|yz}^A + (1 - \mu) \text{tr}(\sigma_{bc|yz}^A) \mathbb{1}/2$$

- $p(abc|xyz) = \text{tr}_A (\Pi_{a|x}(\mu) \sigma_{bc|yz}^A) = \text{tr}_A (\Pi_{a|x} \sigma_{bc|yz}^A(\mu))$
- Noisy measurements are linear combinations of (finite number) PVMs.

(ii) Qubit assemblage, local for all PVM

$$\Pi_{a|x}(\mu) = \mu \Pi_{a|x} + (1 - \mu) \mathbb{1}/2, \quad \sigma_{bc|yz}^A(\mu) = \mu \sigma_{bc|yz}^A + (1 - \mu) \text{tr}(\sigma_{bc|yz}^A) \mathbb{1}/2$$

- $p(abc|xyz) = \text{tr}_A (\Pi_{a|x}(\mu) \sigma_{bc|yz}^A) = \text{tr}_A (\Pi_{a|x} \sigma_{bc|yz}^A(\mu))$
- Noisy measurements are linear combinations of (finite number) PVMs.

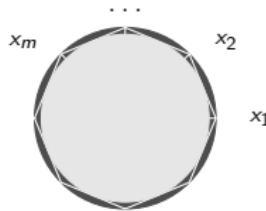


$\sigma_{bc|yz}^A$ **local for** $\{x_1, \dots, x_m\}$

(ii) Qubit assemblage, local for all PVM

$$\Pi_{a|x}(\mu) = \mu \Pi_{a|x} + (1 - \mu) \mathbb{1}/2, \quad \sigma_{bc|yz}^A(\mu) = \mu \sigma_{bc|yz}^A + (1 - \mu) \text{tr}(\sigma_{bc|yz}^A) \mathbb{1}/2$$

- $p(abc|xyz) = \text{tr}_A (\Pi_{a|x}(\mu) \sigma_{bc|yz}^A) = \text{tr}_A (\Pi_{a|x} \sigma_{bc|yz}^A(\mu))$
- Noisy measurements are linear combinations of (finite number) PVMs.

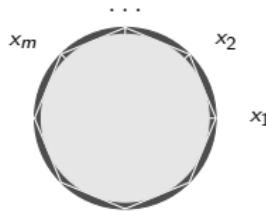


$\sigma_{bc|yz}^A$ **local for** $\{x_1, \dots, x_m\}$ \Leftrightarrow $\sigma_{bc|yz}^A$ **local for** $\Pi_{a|x}(\mu)$

(ii) Qubit assemblage, local for all PVM

$$\Pi_{a|x}(\mu) = \mu \Pi_{a|x} + (1 - \mu) \mathbb{1}/2, \quad \sigma_{bc|yz}^A(\mu) = \mu \sigma_{bc|yz}^A + (1 - \mu) \text{tr}(\sigma_{bc|yz}^A) \mathbb{1}/2$$

- $p(abc|xyz) = \text{tr}_A (\Pi_{a|x}(\mu) \sigma_{bc|yz}^A) = \text{tr}_A (\Pi_{a|x} \sigma_{bc|yz}^A(\mu))$
- Noisy measurements are linear combinations of (finite number) PVMs.



$\sigma_{bc|yz}^A$ **local for** $\{x_1, \dots, x_m\}$ \Leftrightarrow $\sigma_{bc|yz}^A$ **local for** $\Pi_{a|x}(\mu)$

\Leftrightarrow $\sigma_{bc|yz}^A(\mu)$ **local** $\forall \Pi_{a|x}$

Postquantum steering: example without postquantum nonlocality

- Four dichotomic measurements (X, Z)
- Search:

Fix $F_{bc|yz}$:

- Compute $\beta_{\tilde{Q}}$ (SDP).
- Find max violation of the inequality by the 'local' assemblages (SDP).
→ $\sigma_{bc|yz}^A$
- Compute $\beta^* = \text{tr} \sum_{bcyz} F_{bcyz} \sigma_{bc|yz}^*$, $\sigma_{bc|yz}^* := \sigma_{bc|yz}^A (\mu = \cos(\frac{\pi}{8}))$
If $\beta^* > \beta_{\tilde{Q}}$: done! , otherwise, change $F_{bc|yz}$, start over.

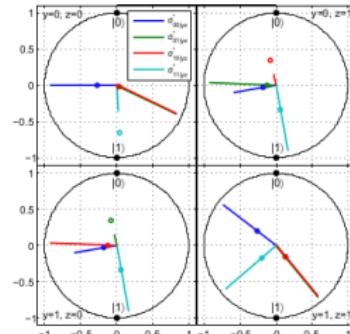
Postquantum steering: example without postquantum nonlocality

- Four dichotomic measurements (X, Z)
- Search:

Fix $F_{bc|yz}$:

- Compute $\beta_{\tilde{Q}}$ (SDP).
- Find max violation of the inequality by the 'local' assemblages (SDP).
→ $\sigma_{bc|yz}^A$
- Compute $\beta^* = \text{tr} \sum_{bcyz} F_{bcyz} \sigma_{bc|yz}^*$, $\sigma_{bc|yz}^* := \sigma_{bc|yz}^A (\mu = \cos(\frac{\pi}{8}))$
If $\beta^* > \beta_{\tilde{Q}}$: done! , otherwise, change $F_{bc|yz}$, start over.

$\sigma_{bc|yz}^*$ is a postquantum qubit assemblage and always gives quantum correlations for PVMs



Postquantum steering: example without postquantum nonlocality

- Four dichotomic measurements (X, Z)
- Search:

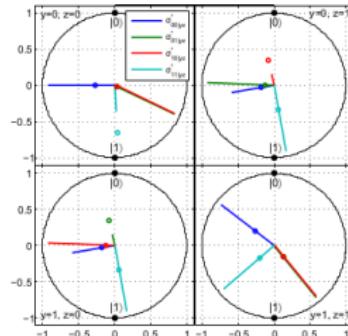
Fix $F_{bc|yz}$:

- Compute $\beta_{\tilde{Q}}$ (SDP).
- Find max violation of the inequality by the 'local' assemblages (SDP).
→ $\sigma_{bc|yz}^A$
- Compute $\beta^* = \text{tr} \sum_{bcyz} F_{bcyz} \sigma_{bc|yz}^*$, $\sigma_{bc|yz}^* := \sigma_{bc|yz}^A (\mu = \cos(\frac{\pi}{8}))$
If $\beta^* > \beta_{\tilde{Q}}$: done! , otherwise, change $F_{bc|yz}$, start over.

$\sigma_{bc|yz}^*$ is a postquantum qubit assemblage and always gives quantum correlations for PVMs

$$\tilde{\sigma}_{bc|yz}^* = \frac{1}{3} \sigma_{bc|yz}^* + \frac{2}{3} \text{tr} (\sigma_{bc|yz}^*) |2\rangle\langle 2|$$

$\tilde{\sigma}_{bc|yz}^*$ is a postquantum qutrit assemblage and always gives quantum correlations for POVMs



Summary and open questions

- Steering beyond quantum theory → multipartite scenarios
- Genuinely new effect
→ postquantum steering $\not\Rightarrow$ postquantum nonlocality
- Fundamental difference between bipartite and multipartite scenarios

Summary and open questions

- Steering beyond quantum theory → multipartite scenarios
- Genuinely new effect
→ postquantum steering $\not\Rightarrow$ postquantum nonlocality
- Fundamental difference between bipartite and multipartite scenarios
- Insight on the characterisation of quantum phenomena
- General framework for non-signalling assemblages
→ quantify postquantumness
- Information-theoretic applications of postquantum steering

Thanks !!!

Ana Belén Sainz, Nicolas Brunner, Daniel Cavalcanti, Paul Skrzypczyk, Tamás Vértesi

Phys. Rev. Lett. 115, 190403 (2015)