

# Comb photon and intriguing features in Comb laser-atom interaction

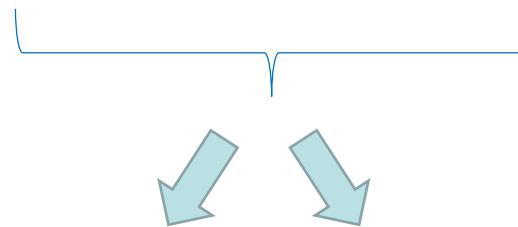
*Wang-Yau Cheng (鄭玉曜)*

*Department of Physics, National Central University, Taiwan*

2021 workshop on quantum technology

*I am a comb laser or “clock” expert, not  
“quantum optics”*

However, it might be a good chance to sell  
my unique comb laser to people in the field  
of “quantum optics”

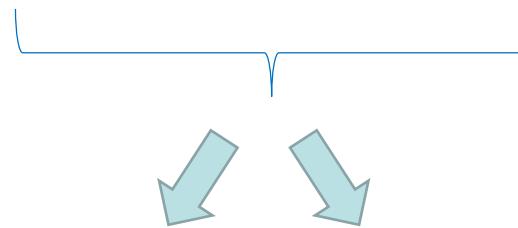


Quantum state  
of light

Photon-atom  
interaction

*I am a comb laser or “clock” expert, not  
“quantum optics”*

However, it might be a good chance to sell  
my unique comb laser to people in the field  
of “quantum optics”



Quantum state  
of light

Photon-atom  
interaction

# Outline

## 1. What is comb laser?

- From the perspective of spectroscopist
- From the perspective of “quantum optics” people

## 2. 40-femtosecond pulse train simultaneously resolves Rb and Cs spectra with 3-kHz resolution

1. Spectral line narrowing
2. Multi-pathway AC stark shift
3. Quantum interfered spectra

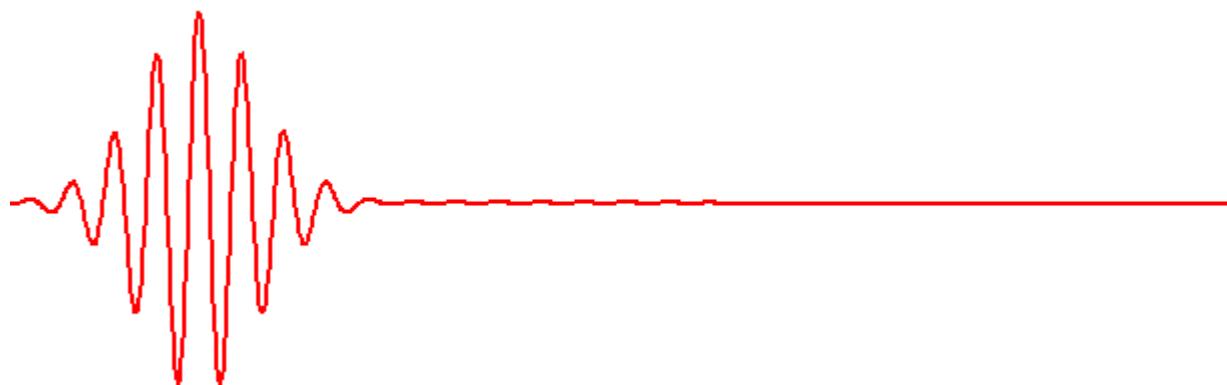
## 3. Application: a novel comb laser without need of expensive cesium clock

What is comb laser (I):

from the perspective of “laser spectroscopy”

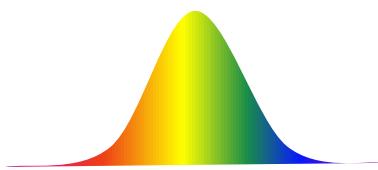
# Viewpoint of time domain:

- High peak power (compared to CW laser)
  - Femtosecond time scale
  - Fixed carrier-envelope phase
- } Good for controlling  
extremely nonlinear optics
- Good for selecting  
atomic quantum states



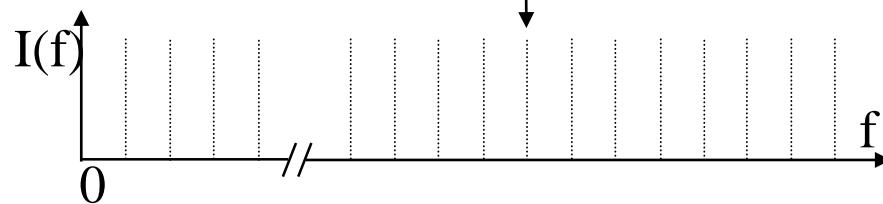
# Viewpoint of frequency domain:

- Wide-band

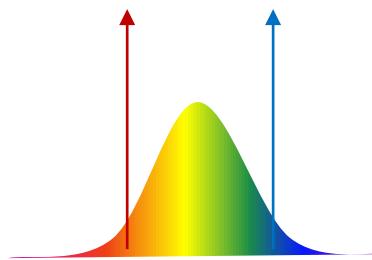


$F_n$  fixed to < 1 kHz instability

- High-resolution



- High coherence



Coherent-state photons  
with different colors

In this talk, we demonstrate:

“wide-band (40 nm) & high resolution (5 kHz) laser spectroscopy”

What is comb laser (II):

from the perspective of “quantum optics”

Cohere state is a superposition of Folk states photons

$$|\alpha\rangle_t = e^{-|\alpha|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha e^{-i\omega t})^n}{\sqrt{n!}} e^{-i\omega t/2} |n\rangle$$

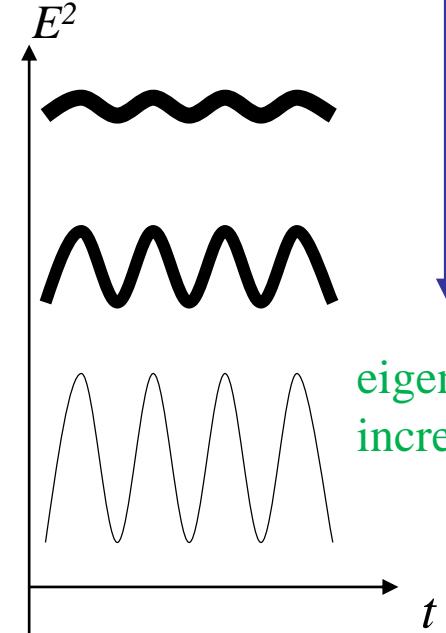
Coherent state photon reach the minimum criteria of uncertainty principle

$$\Delta q \times \Delta p = \frac{1}{2} \hbar$$

$$p \sim H \sim \frac{\partial E}{\partial t} \sim \text{phase of } E$$

$q \sim \text{amplitude of } E$

Photon number state



eigenvalue  $\alpha$  increased

over-complete

$$g^{(2)}(\tau) = 1 \quad \langle \alpha | \alpha' \rangle \neq 0$$

# Frequency-stabilized laser provide “Coherent state” photon

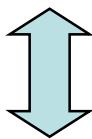
*frequency instability 40 Hz  
Over  $3.2 \times 10^{14}$  Hz duty cycle*



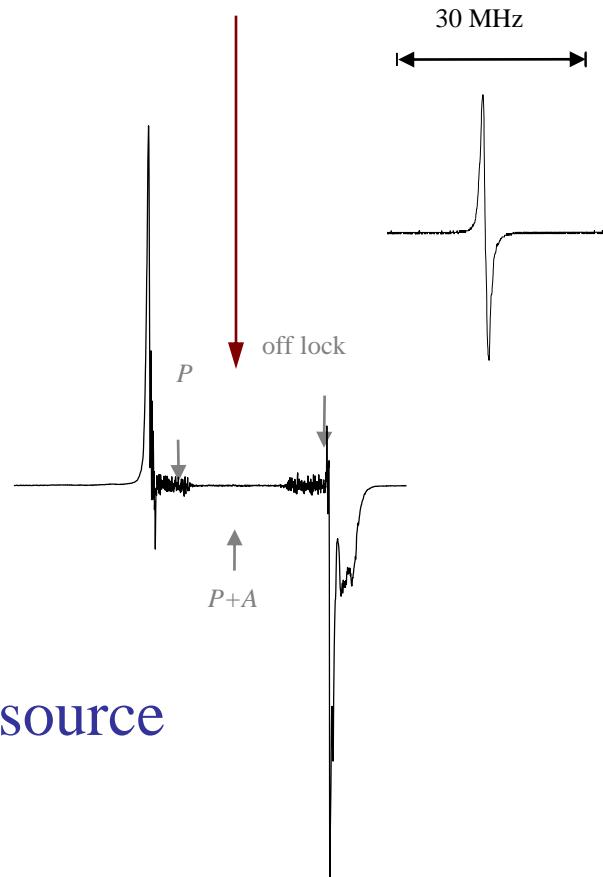
cesium 6S-8S  
dipole not-allowed transition

Opt. lett. **32**, 563 (2007)

electronics for  
feedback control



Coherent-state light source



Comb laser is a superposition of  $N$  ( $N=10^5\text{-}10^6$ ) frequency-stabilized lasers

$$|\alpha_{\omega_1}\rangle_t (= e^{-|\alpha|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha e^{-i\omega_1 t})^n}{\sqrt{n!}} e^{-i\omega_1 t/2} |n\rangle)$$

$$+ |\alpha_{\omega_2}\rangle_t (= e^{-|\alpha|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha e^{-i\omega_2 t})^n}{\sqrt{n!}} e^{-i\omega_2 t/2} |n\rangle)$$

⋮

$$+ |\alpha_{\omega_N}\rangle_t (= e^{-|\alpha|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha e^{-i\omega_N t})^n}{\sqrt{n!}})$$

They all have  
the same phase  
(mode locked)

$$\Delta q \times \Delta p = \frac{1}{2} \hbar$$

$$p \sim H \sim \frac{\partial E}{\partial t} \sim \text{phase of } E$$

$q \sim \text{amplitude of } E$

Superposition of all coherence-state light with wide-band mode locked carrier frequencies (femtosecond pulse train)



Via Spontaneous down conversion



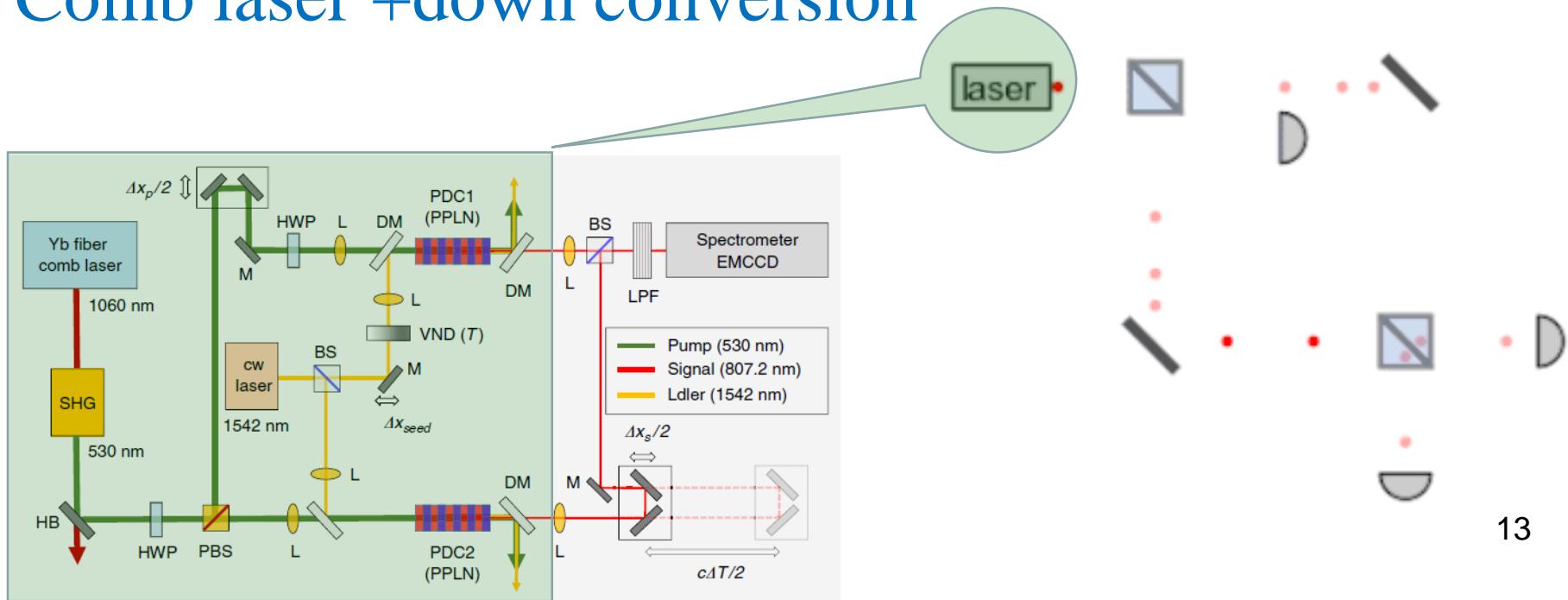
Still keep the anti-bunching property

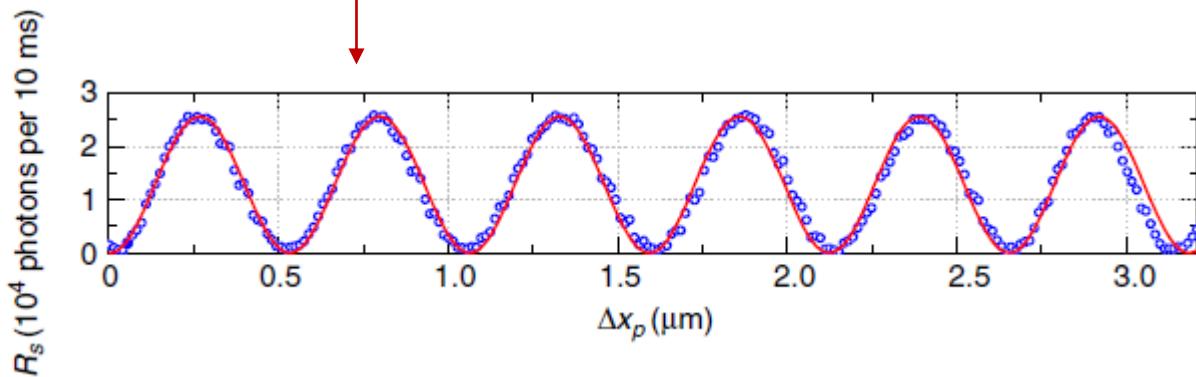
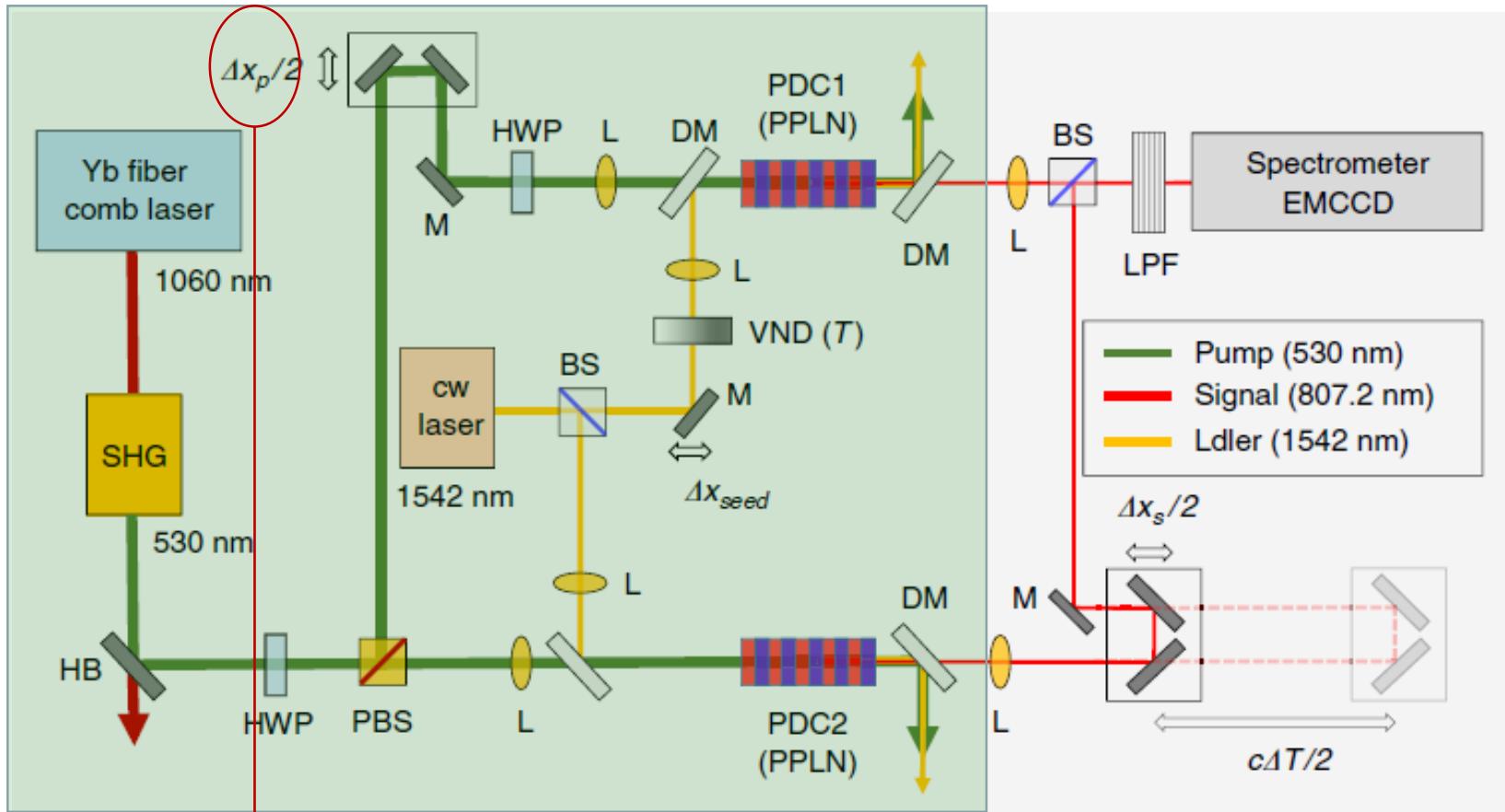
# Frequency comb single-photon interferometry

Sun Kyung Lee<sup>1</sup>, Noh Soo Han<sup>1</sup>, Tai Hyun Yoon<sup>1,2</sup> & Minhaeng Cho<sup>1,3</sup>

Use comb laser as a light source for “which-way interferometer” experiment

## Comb laser +down conversion





Application:  
Quantum metrology  
(Quantum tick-tack clock)

Their very recent report by using comb laser:

SCIENCE ADVANCES | RESEARCH ARTICLE

PHYSICS

## Quantitative complementarity of wave-particle duality

Tai Hyun Yoon<sup>1,2\*</sup> and Minhaeng Cho<sup>1,3\*</sup>

Yoon and Cho, *Sci. Adv.* 2021; 7 : eabi9268

18 August 2021

# Quantum interference in comb laser-atom interaction

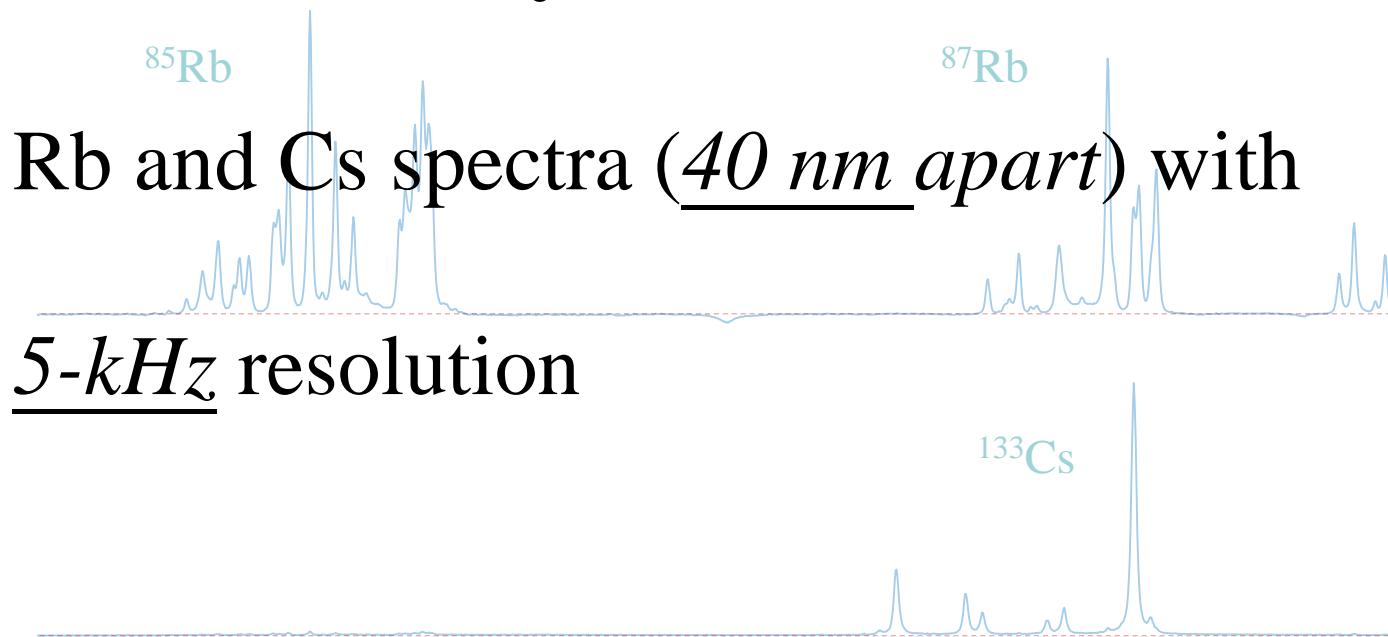
Tz-Wei Liu



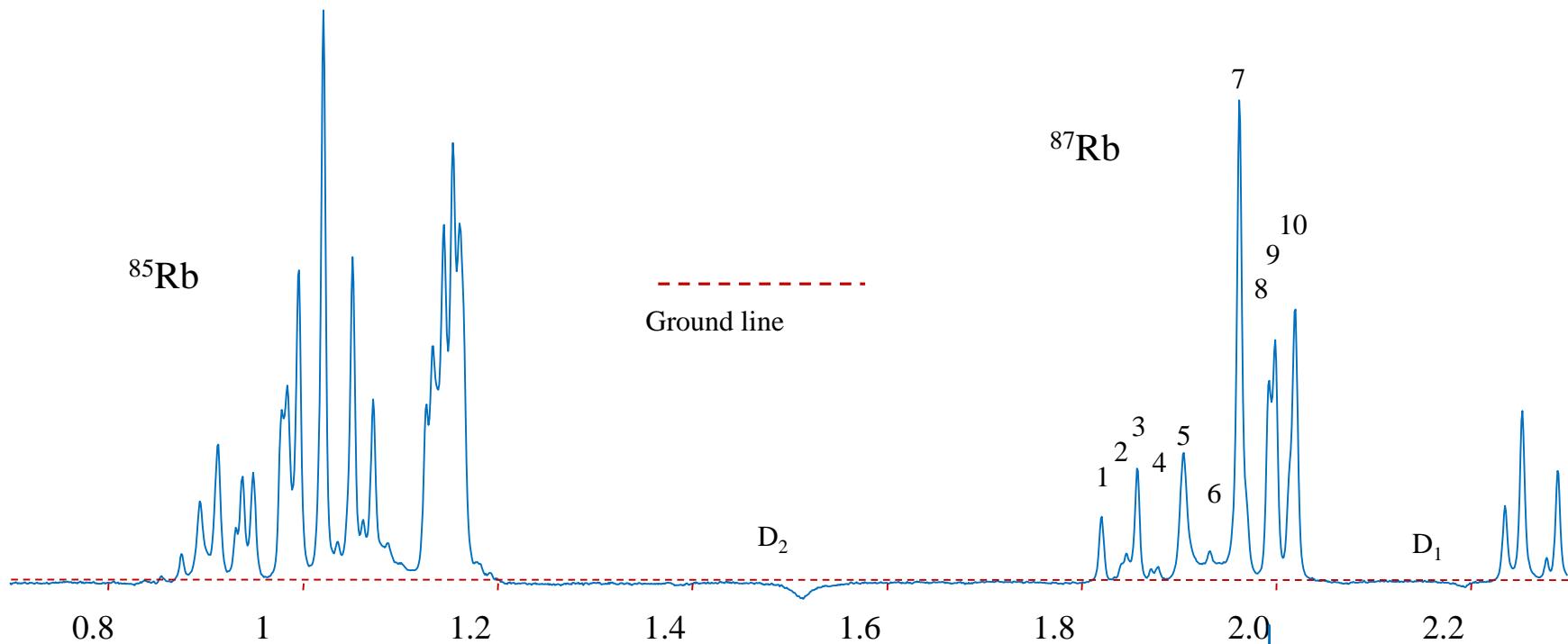
He is currently looking for a teaching job 16

$40 \times 10^{-15}$  second pulse laser

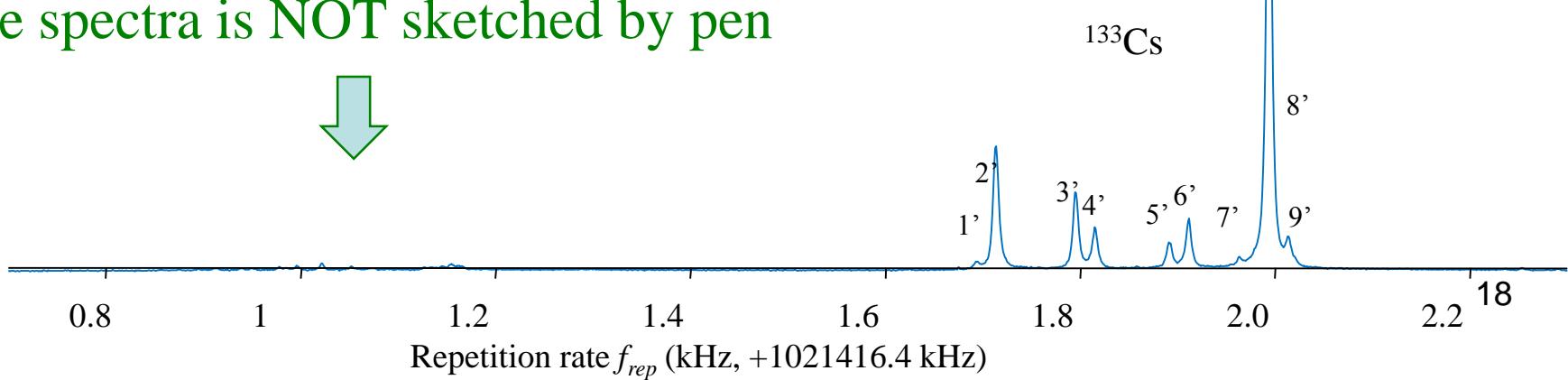
**Simultaneously** resolves



# Direct frequency comb spectroscopy in record resolution

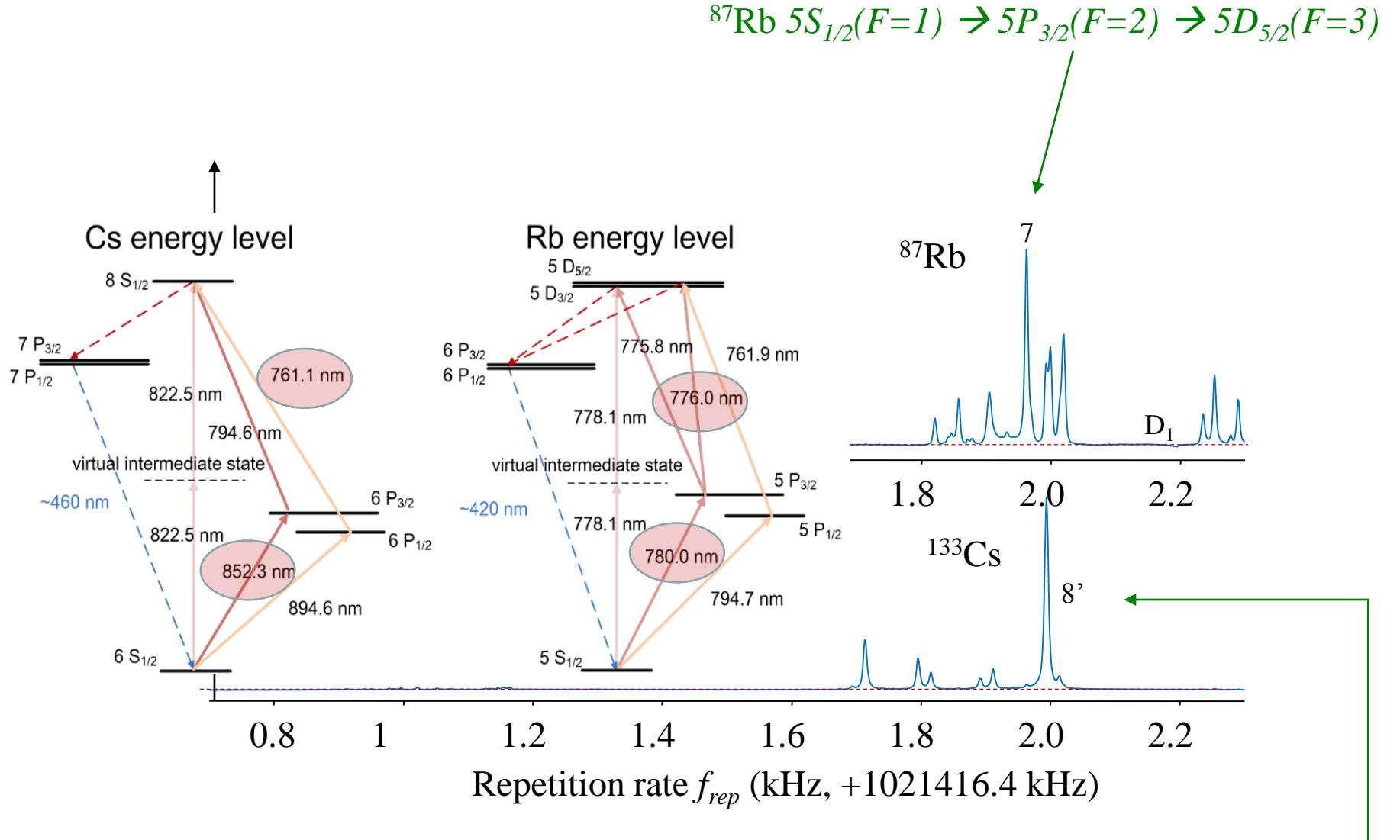


The spectra is NOT sketched by pen



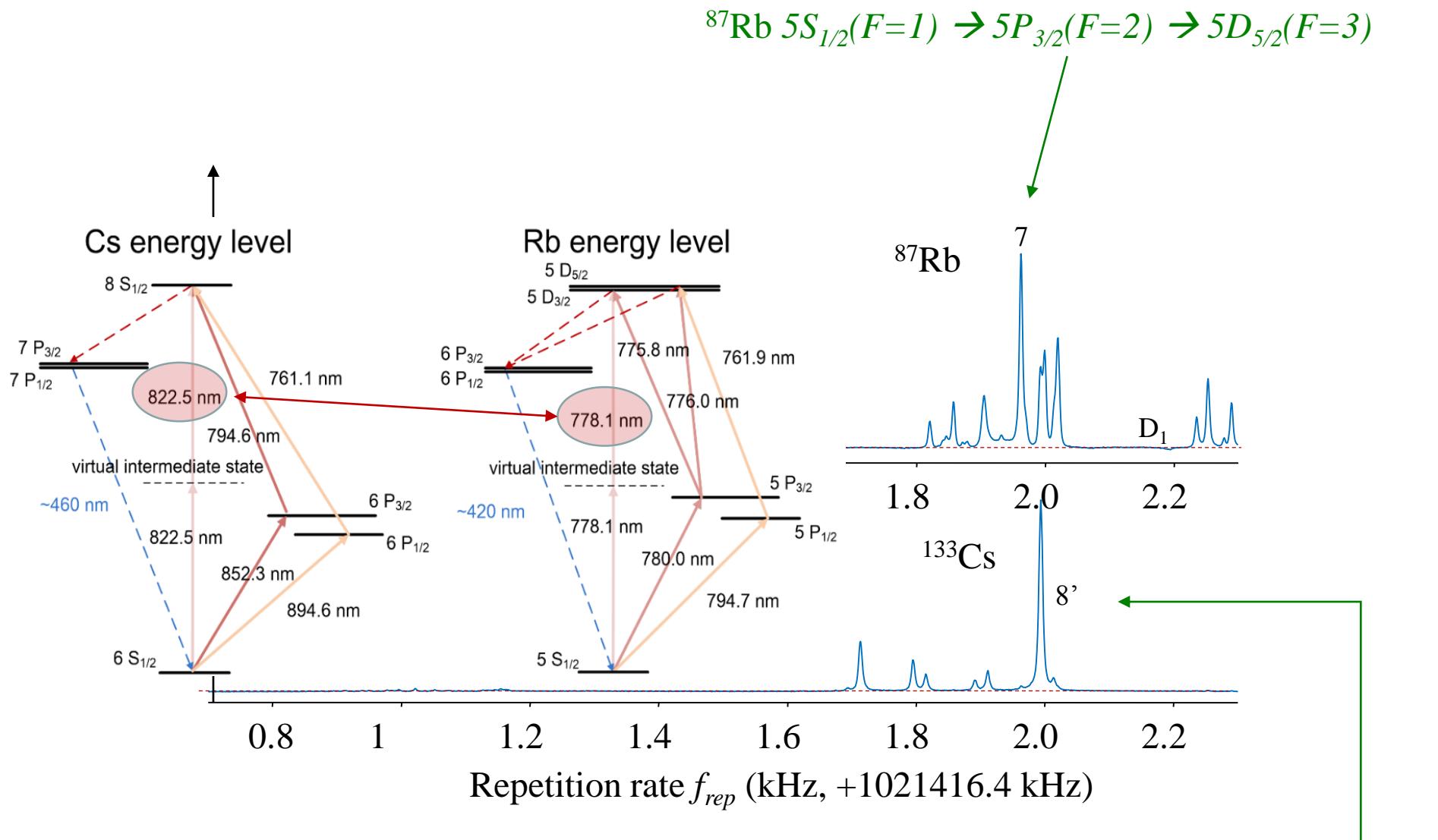
Repetition rate  $f_{rep}$  (kHz, +1021416.4 kHz)

# Two clock transitions



$^{133}\text{Cs } 6S_{1/2}(F=1) \rightarrow 5P_{3/2}(F=2) \rightarrow 5D_{5/2}(F=3)$

The two clocks separates for 40 nm, resolved simultaneously



$^{133}\text{Cs} 6\text{S}_{1/2}(F=1) \rightarrow 5\text{P}_{3/2}(F=2) \rightarrow 5\text{D}_{5/2}(F=3)$

# Some intriguing spectral features (I)

Line narrowing

# Doppler-free two-color spectroscopy of the $6_2S_{1/2}$ - $8_2S_{1/2}$ cesium transition using semiconductor diode lasers

C. Fort<sup>1</sup>, M. Inguscio<sup>1</sup>, P. Raspollini<sup>1</sup>, F. Baldes<sup>2</sup>, A. Sasso<sup>2</sup>

Appl. Phys. B 61, 467–472 (1995)

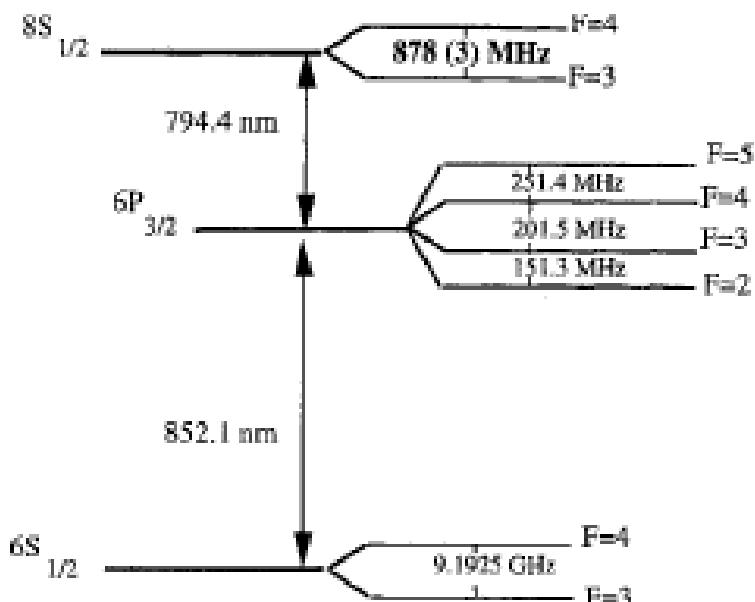


Fig. 4. Energy-level diagram of the relevant cesium states and the relative transitions involved in this experiment

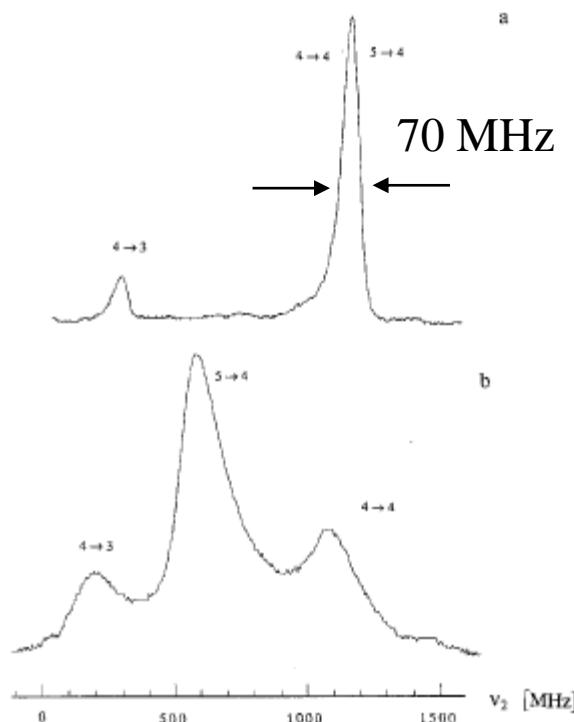


Fig. 6a,b. Two-color transition with DL1 locked onto the  $F = 4 - F = 5$  cross-over resonance. The hyperfine structure of the final state  $8S$  is partially resolved with the counterpropagating scheme (a) while it is fully resolved with copropagating laser beams (b)

# Doppler-free two-color spectroscopy of the $6_2S_{1/2}$ – $8_2S_{1/2}$ cesium transition using semiconductor diode lasers

C. Fort<sup>1</sup>, M. Inguscio<sup>1</sup>, P. Raspollini<sup>1</sup>, F. Baldes<sup>2</sup>, A. Sasso<sup>2</sup>

Appl. Phys. B 61, 467–472 (1995)

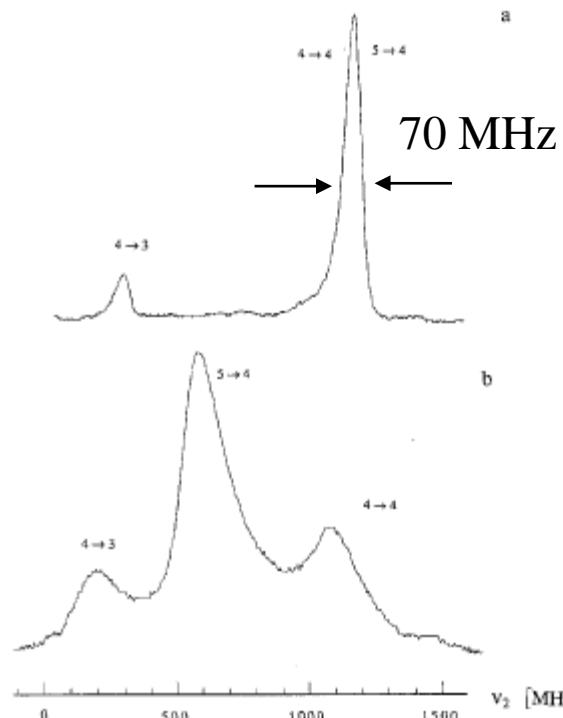
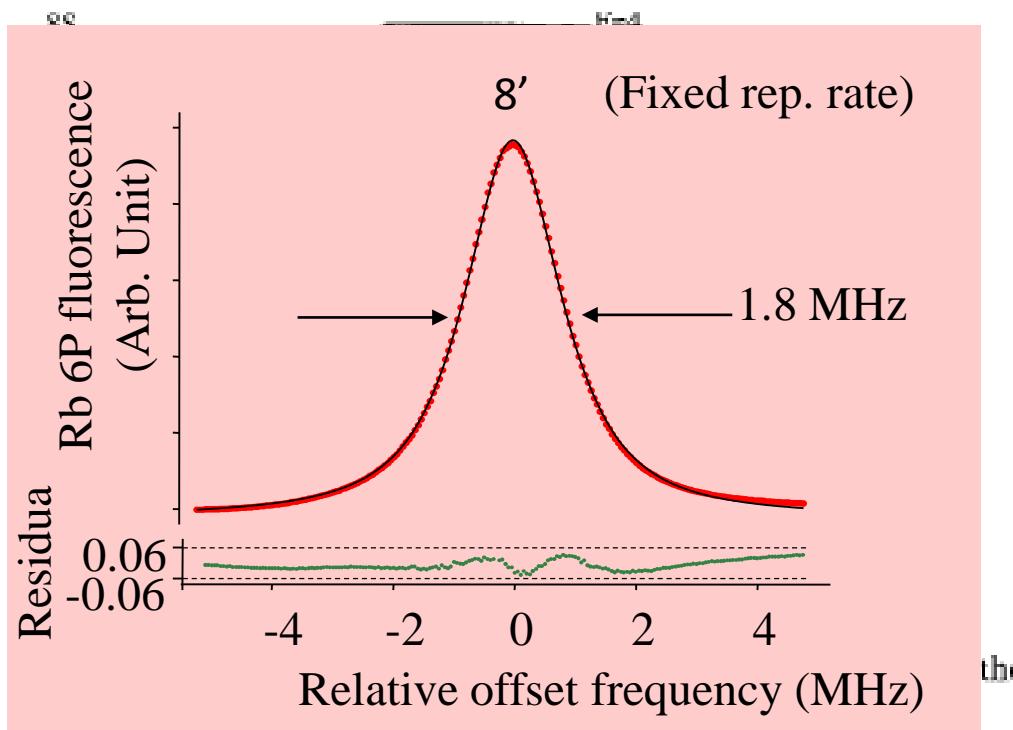
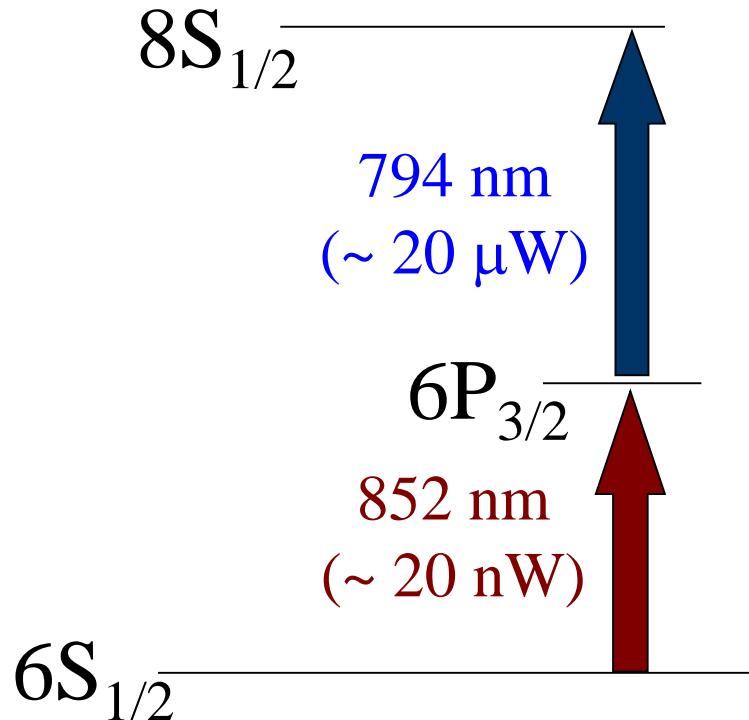


Fig. 6a,b. Two-color transition with DL1 locked onto the  $F = 4$ – $F = 5$  cross-over resonance. The hyperfine structure of the final state  $8S$  is partially resolved with the counterpropagating scheme (a) while it is fully resolved with copropagating laser beams (b)



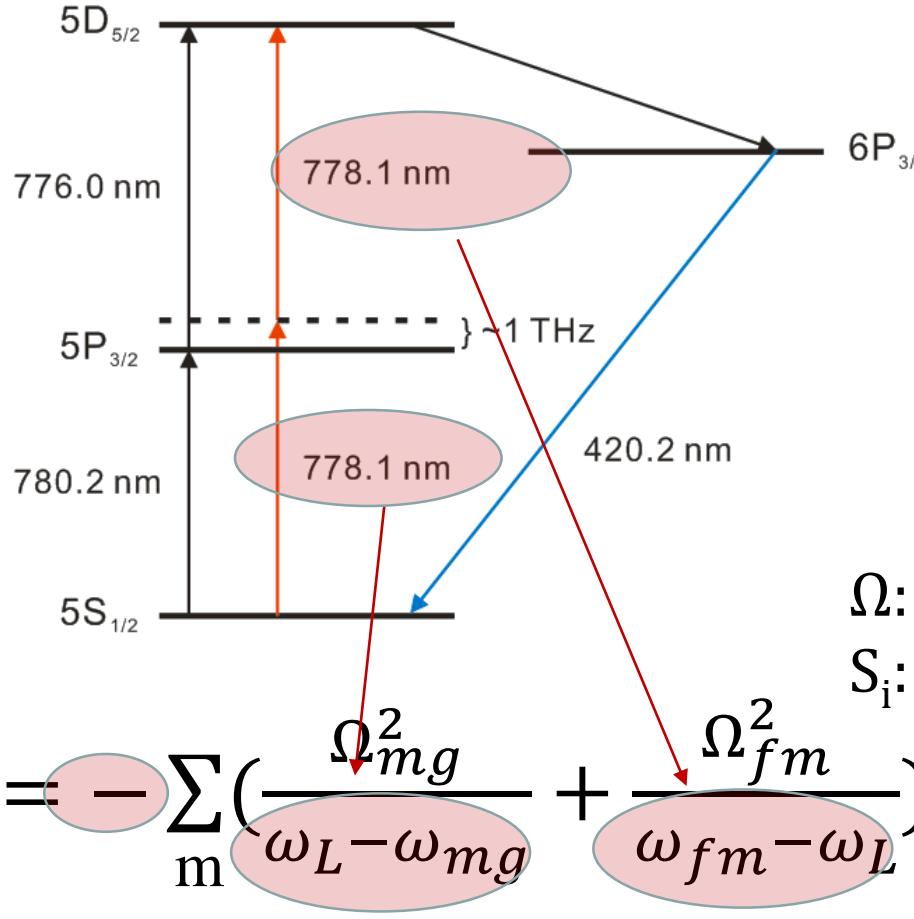
The intermediate state is  
**transparent?**

## Some intriguing spectral features (II)

Multi-pathway AC stark shift

# CW-clocks

$^{85}\text{Rb}$

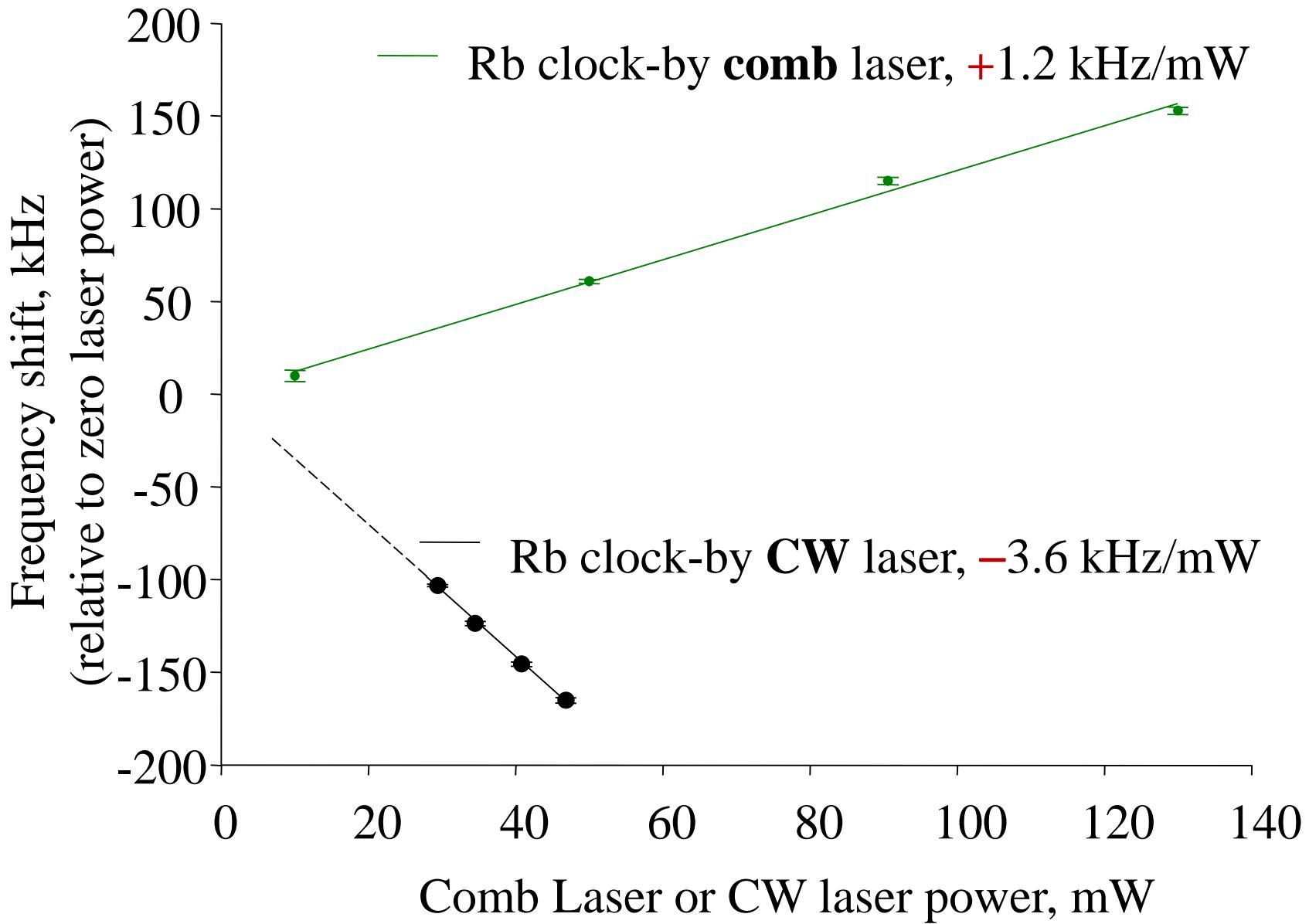


$\Omega$ : Rabi frequency  
 $S_i$ : frequency shift

$$S_{fg} = S_f - S_g = - \sum_m \left( \frac{\Omega_{mg}^2}{\omega_L - \omega_{mg}} + \frac{\Omega_{fm}^2}{\omega_{fm} - \omega_L} \right)$$

Single pathway: Ladder direct two-photon yields red shift

Most concerned issue in clock !!

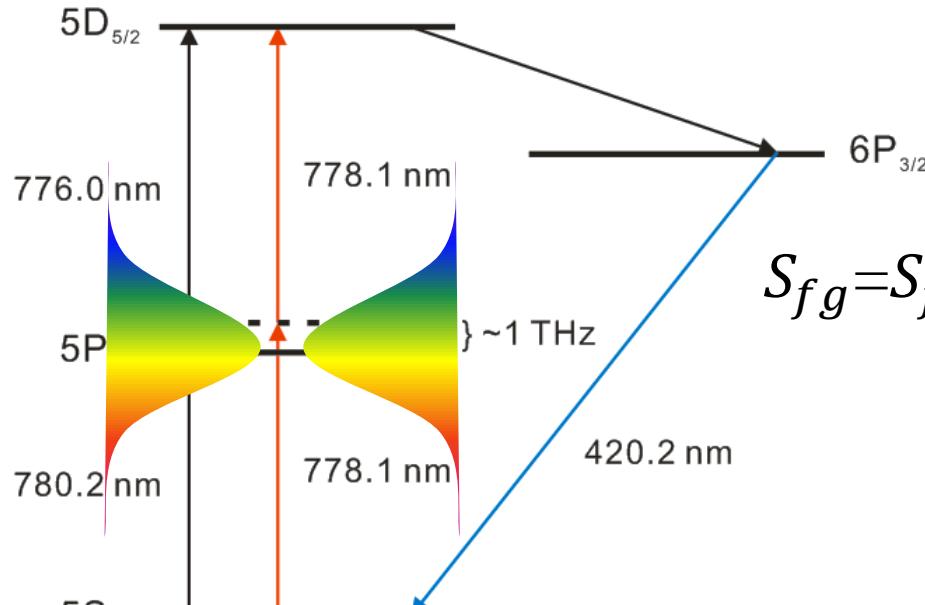


# CW-clocks

## Multi-paths induced AC stark shift



$^{85}\text{Rb}$



$$S_{fg} = S_f - S_g = - \sum_m \left( \frac{\Omega_{mg}^2}{\omega_L - \omega_{mg}} + \frac{\Omega_{fm}^2}{\omega_{fm} - \omega_L} \right)$$



$$S_{fg} = S_f - S_g = - \sum \text{all comb modes} \sum_m \left( \frac{\Omega_{mg}^2}{\omega_L - \omega_{mg}} + \frac{\Omega_{fm}^2}{\omega_{fm} - \omega_L} \right)$$

Question: Would it be possible to quantum controlling the AC stark shift to be zero by pulse shaping?



陽明交通大學

寺西慶哲 教授

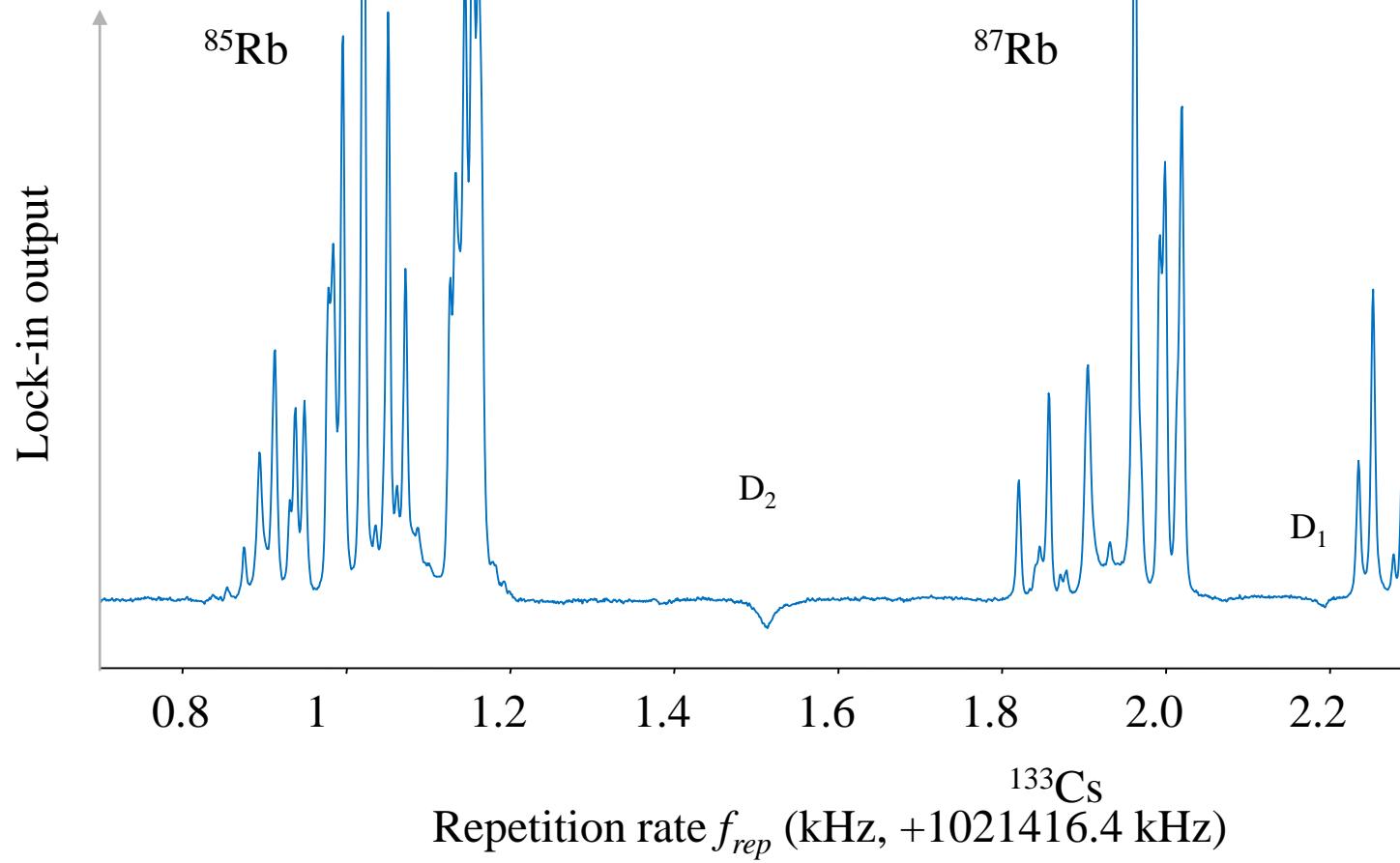
•物理研究所

Most concerned issue in clock !!

## Some intriguing spectral features (III)

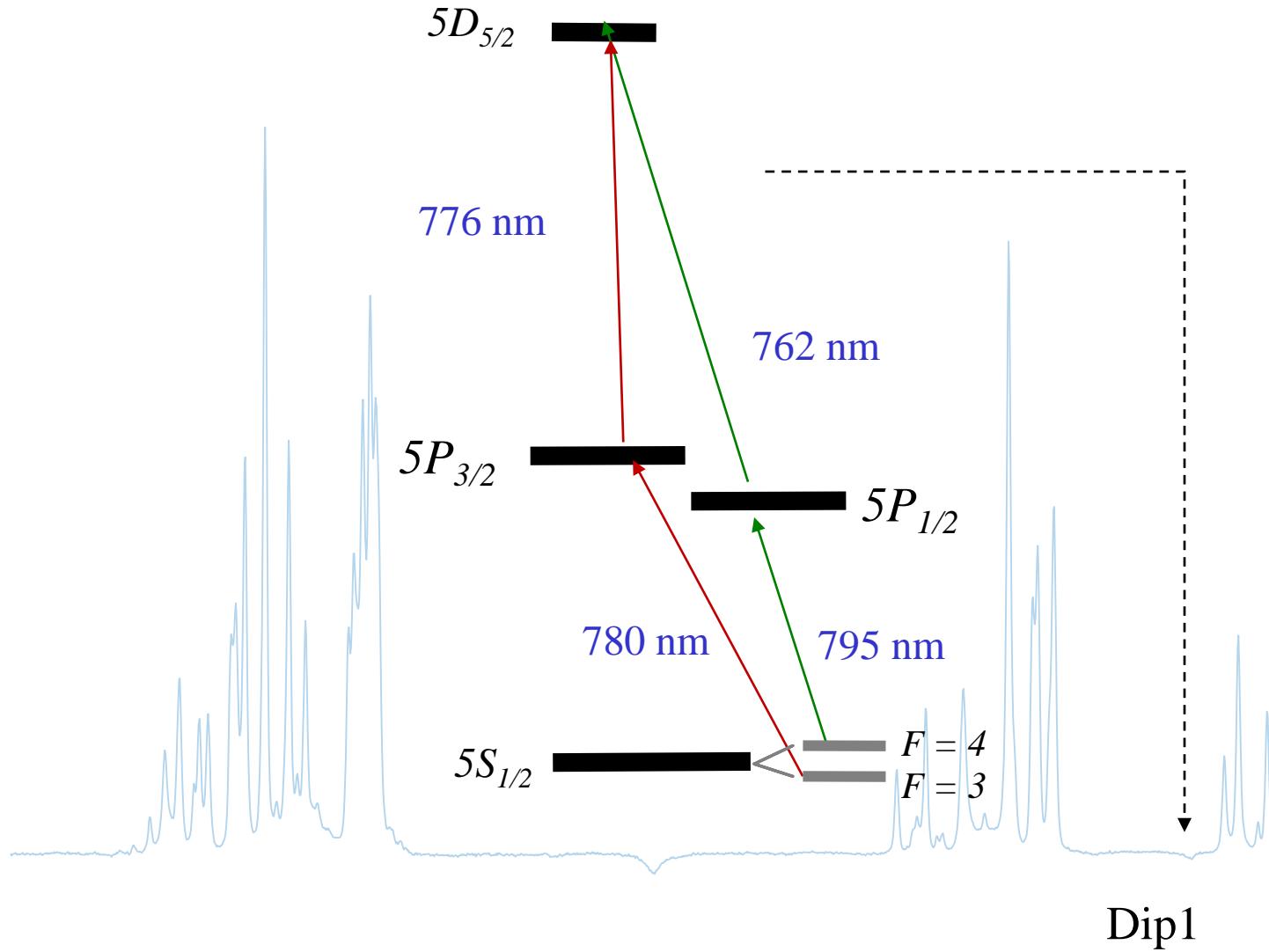
Quantum interference by different comb modes

Why dips ?



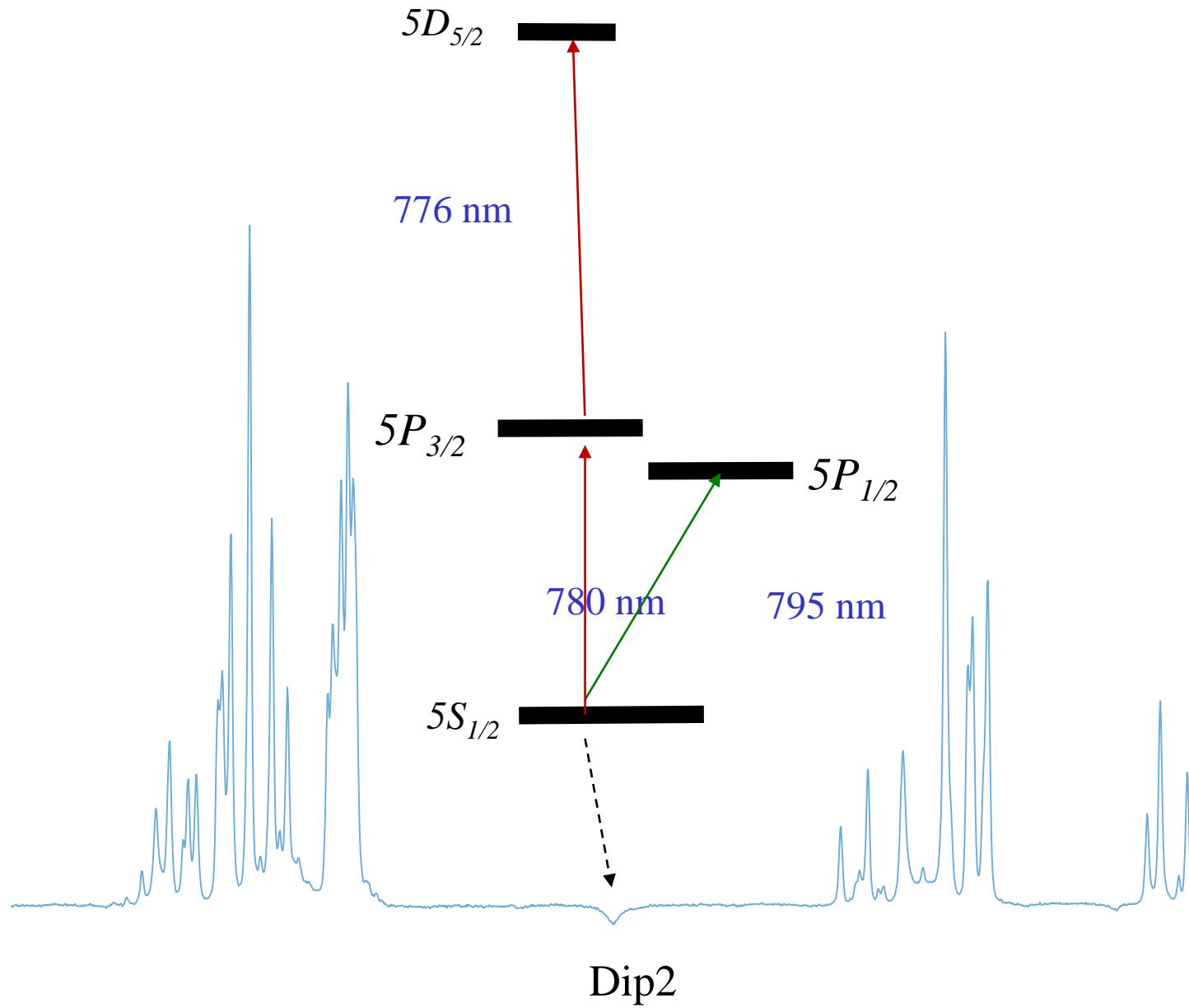
# Two-pathway quantum interference

$^{87}\text{Rb}$



# V type EIT interference

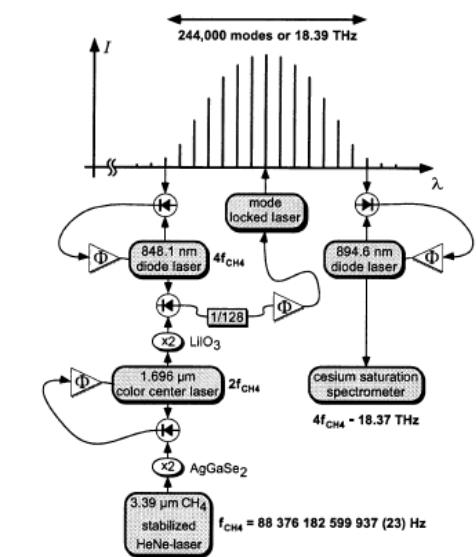
$^{87}\text{Rb}$



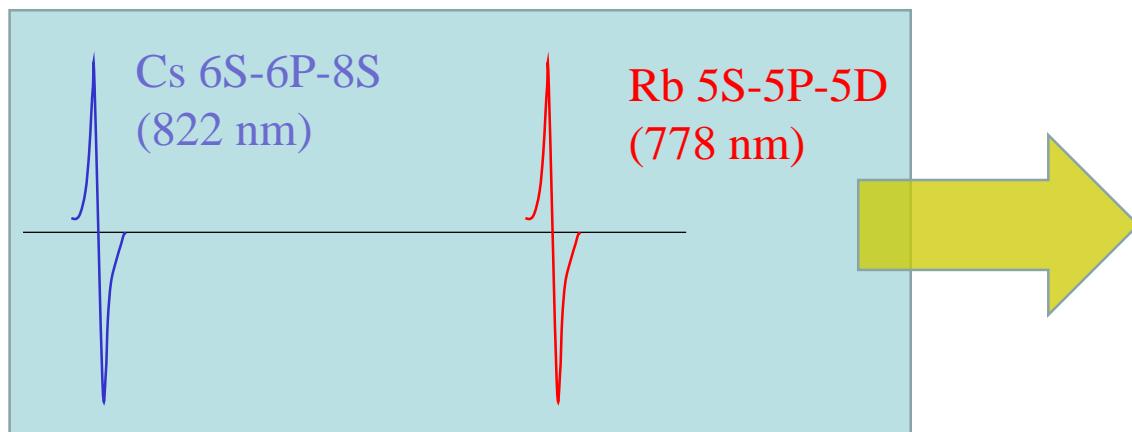
# Application:

## Direct comb clocks

# Ted Hansch



They all need a Cs clock to lock rep. rate!!

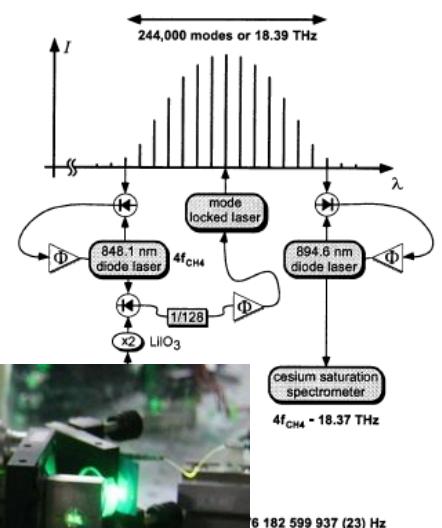
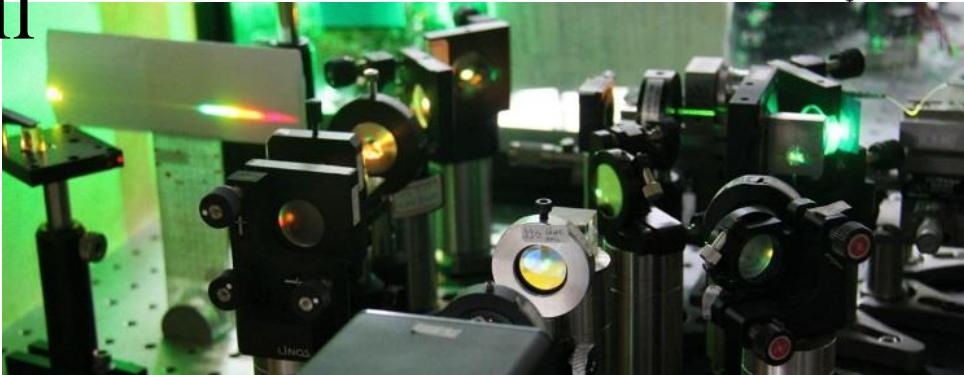


We locked two parameters directly with two spectra (without Cs clock)

# Ted Hansch



Jan Hall

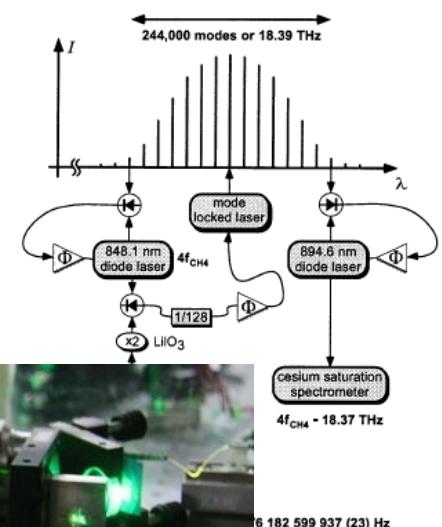
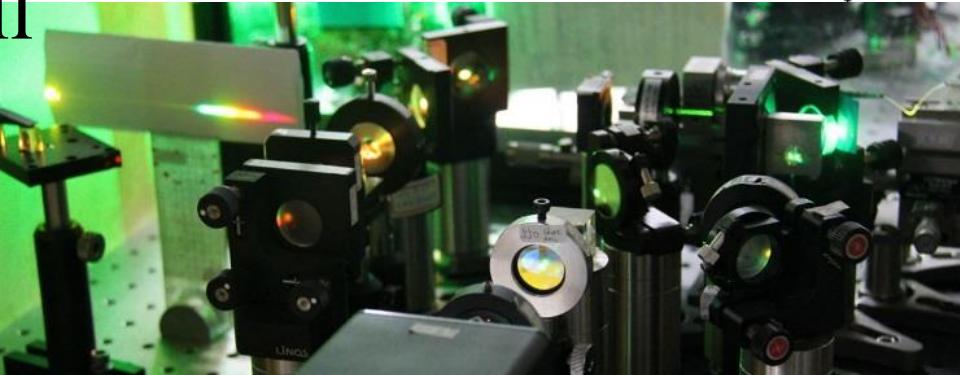


They all need a Cs clock to lock rep. rate!!

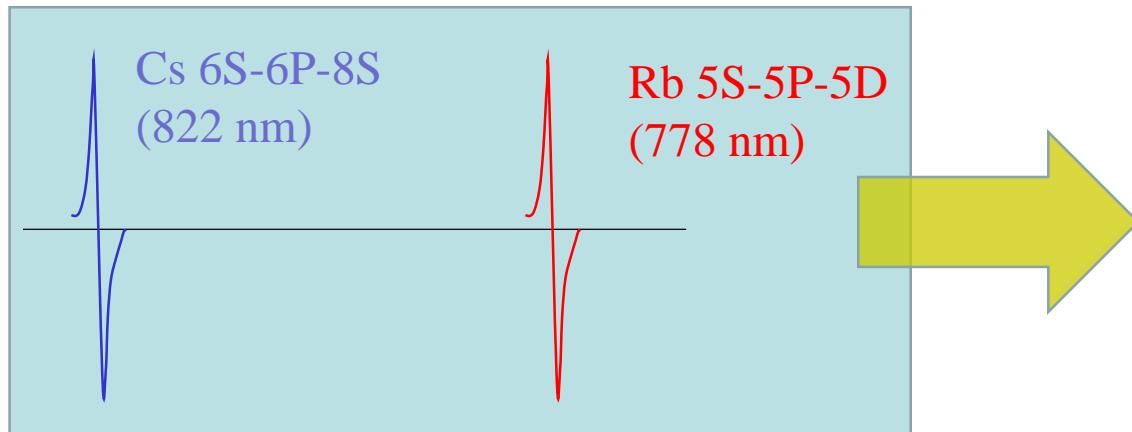
# Ted Hansch



Jan Hall



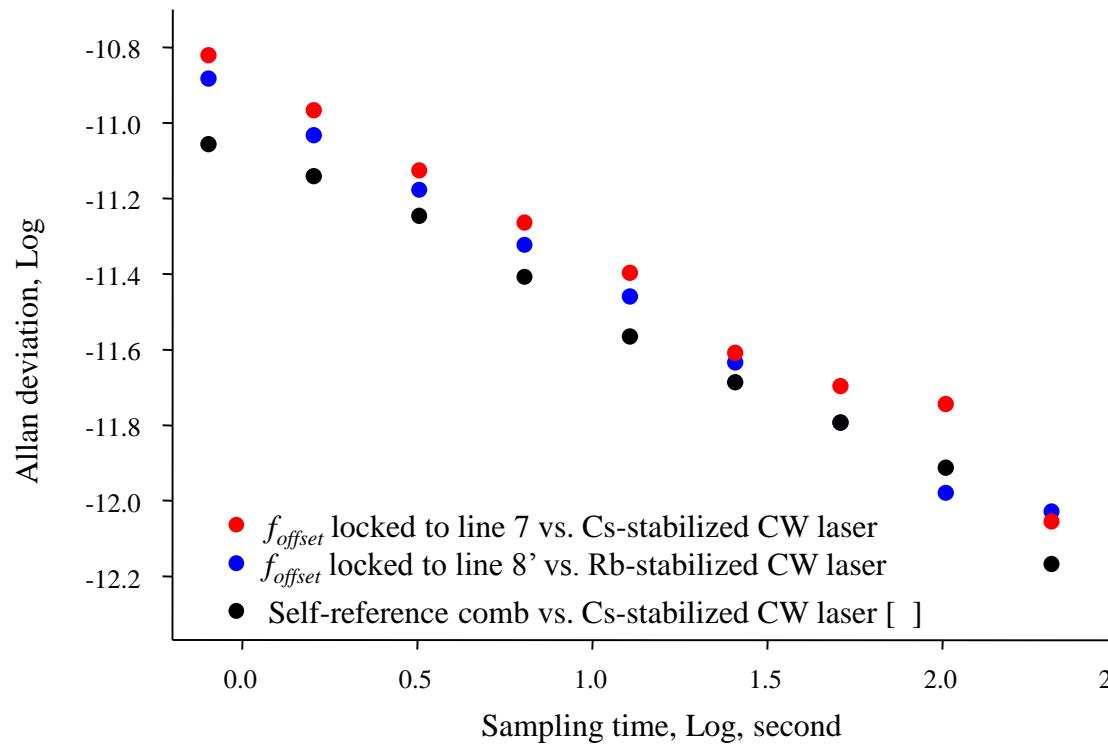
They all need a Cs clock to lock rep. rate!!



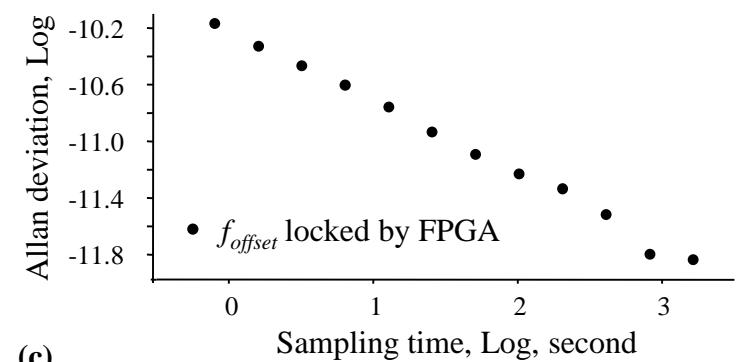
We locked two parameters directly with two spectra (without Cs clock)

Now we need merely one 6-cm Cs/Rb mixed cell to build up comb clocks

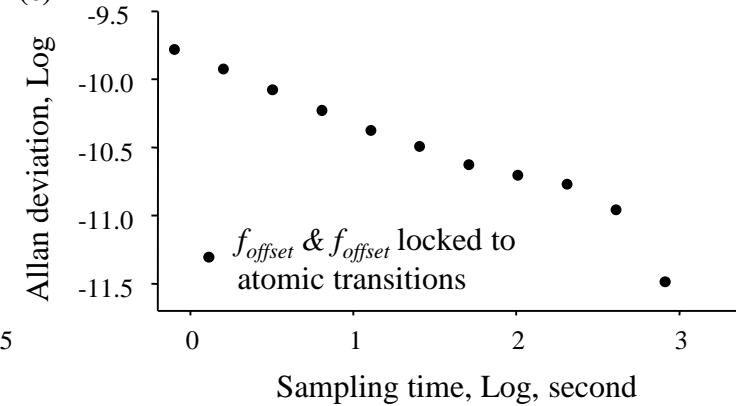
(a)



(b)

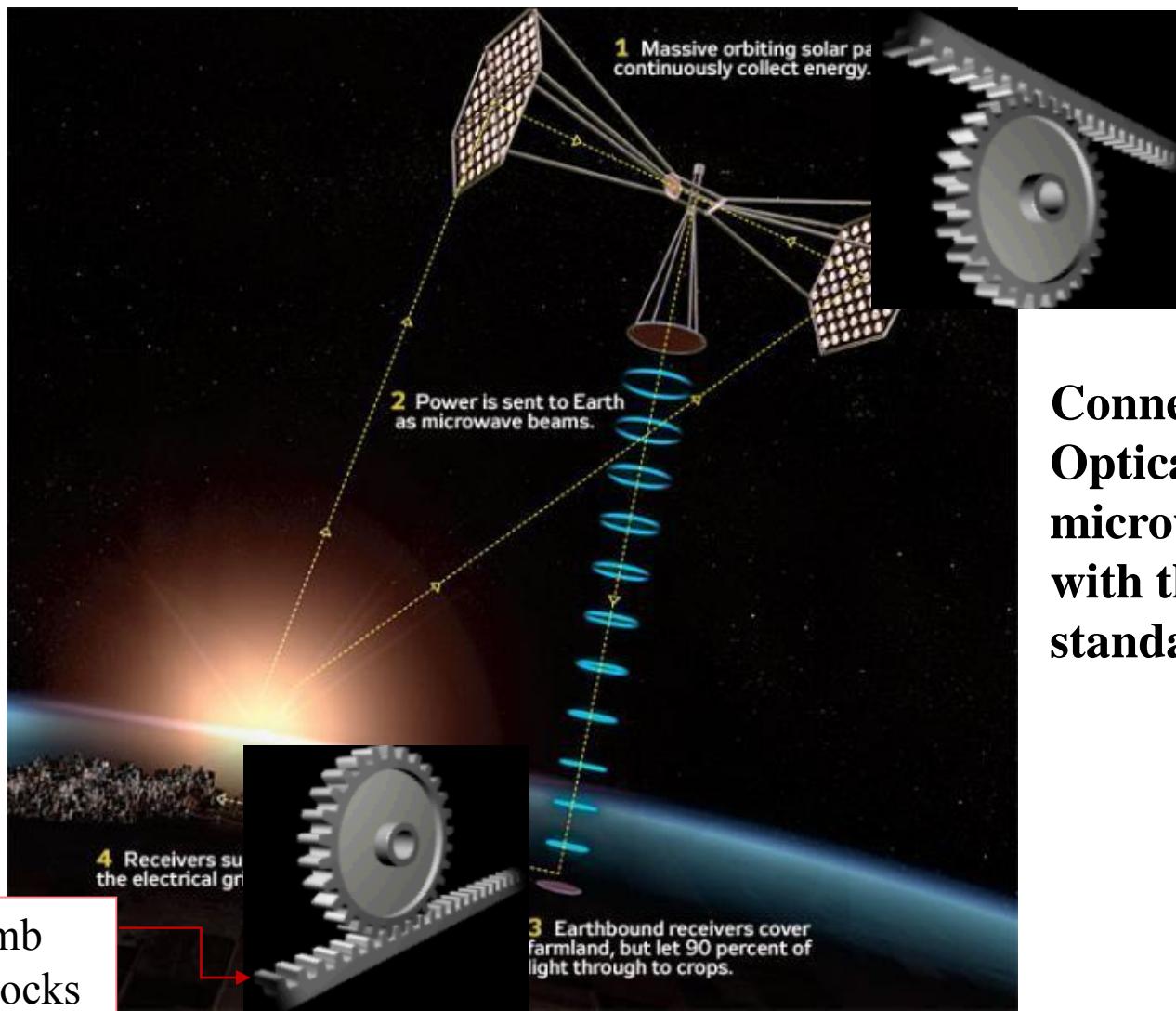


(c)



## Future direction (1)

## Comb laser in artificial satellite



Ground comb  
laser and clocks

**Connects both  
Optical and  
microwave clocks  
with the ground  
standards**

## Future direction (2)

# You can entangle comb photon

PRL 107, 030505 (2011)

PHYSICAL REVIEW LETTERS

week ending  
15 JULY 2011

## Parallel Generation of Quadripartite Cluster Entanglement in the Optical Frequency Comb

Matthew Pysher,<sup>1</sup> Yoshichika Miwa,<sup>2</sup> Reihaneh Shahrokhshahi,<sup>1</sup> Russell Bloomer,<sup>1</sup> and Olivier Pfister<sup>1,\*</sup>

# You can squeeze comb photon

PRL 108, 083601 (2012)

PHYSICAL REVIEW LETTERS

week ending  
24 FEBRUARY 2012

## Generation and Characterization of Multimode Quantum Frequency Combs

Olivier Pinel,<sup>1</sup> Pu Jian,<sup>1</sup> Renné Medeiros de Araújo,<sup>1</sup> Jinxia Feng,<sup>1,2</sup> Benoît Chalopin,<sup>1,3</sup>  
Claude Fabre,<sup>1,\*</sup> and Nicolas Treps<sup>1</sup>

# You can apply comb photon on quantum computation

PRL 101, 130501 (2008)

 Selected for a [Viewpoint](#) in *Physics*  
PHYSICAL REVIEW LETTERS

week ending  
26 SEPTEMBER 2008

# One-Way Quantum Computing in the Optical Frequency Comb

Nicolas C. Menicucci,<sup>1,2</sup> Steven T. Flammia,<sup>3</sup> and Olivier Pfister<sup>4</sup>

You can use comb laser to entangle atomic Qubits

PRL 104, 140501 (2010)

 Selected for a Viewpoint in *Physics*  
PHYSICAL REVIEW LETTERS

week ending  
9 APRIL 2010



## Entanglement of Atomic Qubits Using an Optical Frequency Comb

D. Hayes,\* D. N. Matsukevich, P. Maunz, D. Hucul, Q. Quraishi, S. Olmschenk, W. Campbell, J. Mizrahi, C. Senko, and C. Monroe

## Acknowledgement:

Telecommun. Labs. (中華電信研究所)

Ministry of Science and Technology (科技部)

partner in next step:



陽明交通大學

寺西慶哲 教授

•物理研究所



中央研究院

張銘顯 研究員

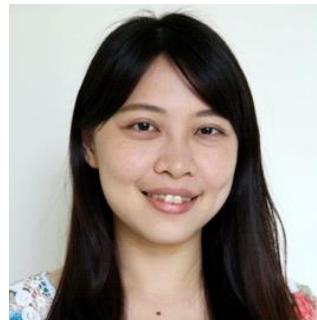
原子分子研究所



中央大學

陳彥宏 教授

•光電研究所



中央大學

李依珊 教授

•電機研究所