EPR STEERING AND THE STEERING ELLIPSOID

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Joint work with

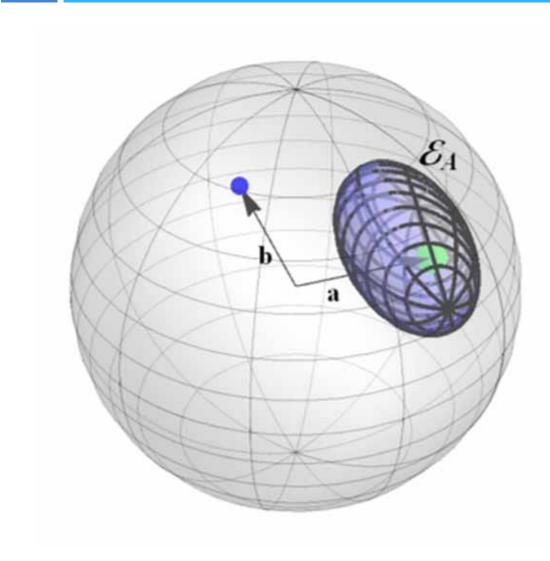
- Michael Hall
- Malcolm Anderson
- Marcin Zwierz
- Howard Wiseman

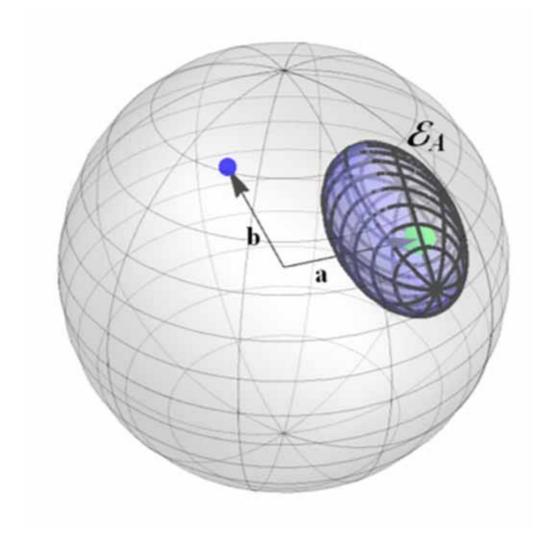
□ Based on: JOSA B, Vol. 32, Issue 4, pp. A40-A49 (2015)

Way of visualising two-qubit state
$$\rho = \frac{1}{4}(I + \boldsymbol{a}.\boldsymbol{\sigma} \otimes I + I \otimes \boldsymbol{b}.\boldsymbol{\sigma} + \sum_{ij} T_{ij}\sigma_i \otimes \sigma_j)$$

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$$\rho = \frac{1}{4}(I + a.\sigma \otimes I + I \otimes b.\sigma + \sum_{ij} T_{ij}\sigma_i \otimes \sigma_j)$$
Bob performs POVM on his qubit {*F_f*}
Upon getting outcome "*f*", Alice is steered to
$$\rho_f = \frac{\operatorname{tr}_B(\rho I \otimes F_f)}{\operatorname{tr}(\rho I \otimes F_f)} = \frac{1}{2}\left(I + \left[\frac{a + Tf}{1 + b.f}\right].\sigma\right)$$

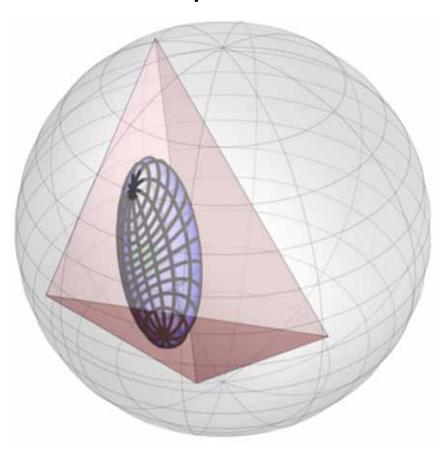
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The steering ellipsoid (for Alice) is generated by varying over all {*F_f*}





- Reconstruct state (up to local Us)
- Rank of ellipsoid
- □ Volume $V > V^* > 0$
- Geometric condition
 for entanglement

The quantum steering ellipsoid gives us an iff condition for a non-separable state



The quantum steering ellipsoid gives us an iff condition for entanglement

Can it say anything about other forms of nonlocality?

□ Werner state

$$\rho = (1-p)\frac{I}{4} + p|\psi^{-}\rangle\langle\psi^{-}|$$

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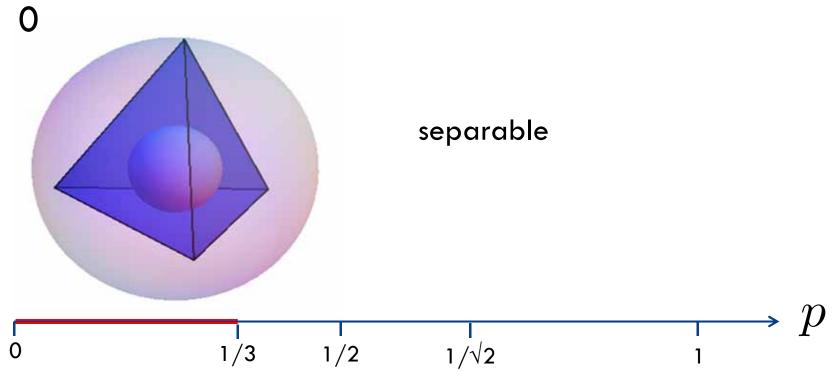
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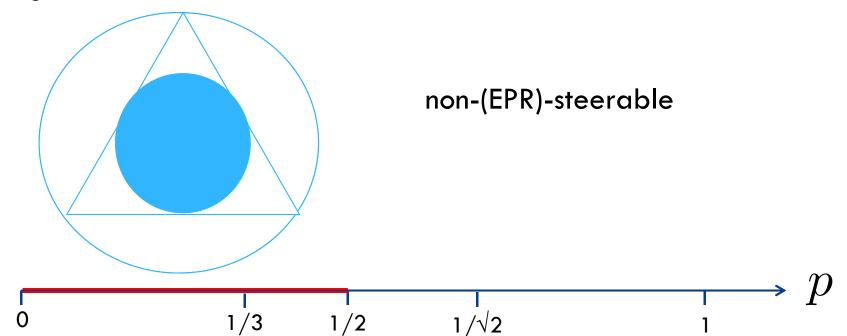
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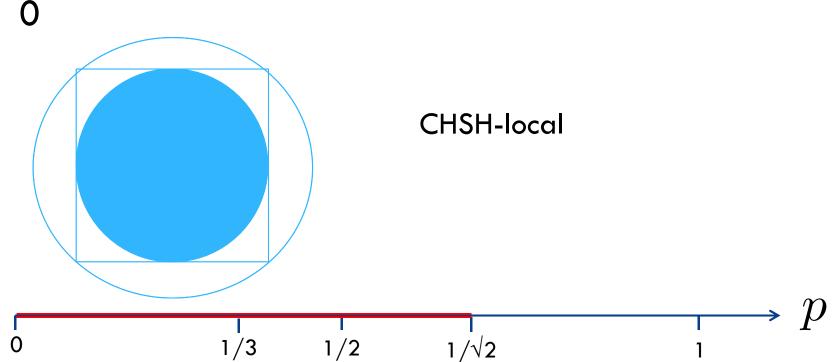
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 $\square \operatorname{Recall} \rho = \frac{1}{4} (I + \boldsymbol{a}.\boldsymbol{\sigma} \otimes I + I \otimes \boldsymbol{b}.\boldsymbol{\sigma} + \sum_{ij} T_{ij}\sigma_i \otimes \sigma_j)$

• Set
$$a = b = 0$$

$$\rho = \frac{1}{4} \left(I + \sum_{ij} T_{ij} \sigma_i \otimes \sigma_j \right)$$

$\square \text{ Mixture of Bell states} \\ \rho = \frac{1}{4} (I + \sum_{i=1}^{3} t_i \sigma_i \otimes \sigma_i)$

Mixture of Bell states
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Quantum steering ellipsoid
 Semiaxes |t_i|
 Aligned along x,y,z - axes
 Centred at 0

A Mixture of Bell states
$$\rho = \frac{1}{4} (I + \sum_{i=1}^{3} t_i \sigma_i \otimes \sigma_i)$$

- Quantum steering ellipsoid
 - Semiaxes $|t_i|$
 - Aligned along x,y,z axes
 - Centred at 0
 - Does not possess nice geometric conditions for LHS or LHV models (but if separable still obeys nested tet cond)

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 LHS model for all projective measurements
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$$\Box \text{ Each outcome } f = \pm 1 \text{ occurs with } p_f = \frac{1}{2}$$

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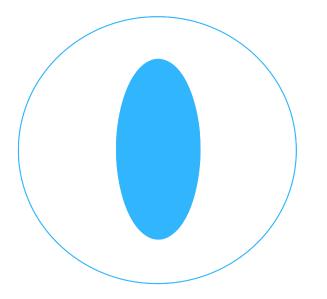
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$$p_f \boldsymbol{b}_f = \int P(\boldsymbol{n}) p(f|\boldsymbol{f}, \boldsymbol{n}) \boldsymbol{n} \, d^2 \boldsymbol{n} = \frac{1}{2} T \boldsymbol{f}$$
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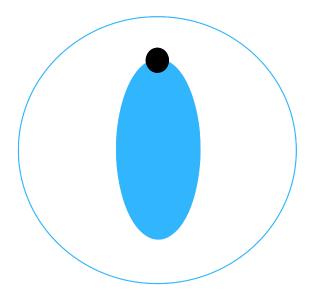
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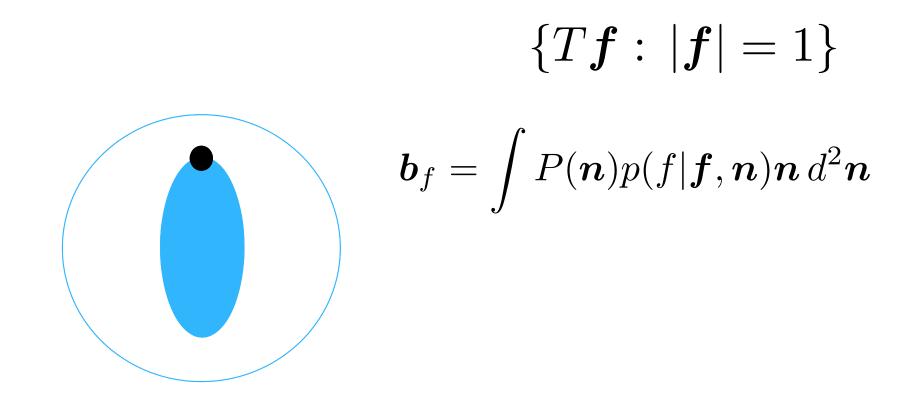
$$\boldsymbol{b}_{f} = \int P(\boldsymbol{n}) p(f|\boldsymbol{f}, \boldsymbol{n}) \boldsymbol{n} \, d^{2}\boldsymbol{n} = T\boldsymbol{f}$$
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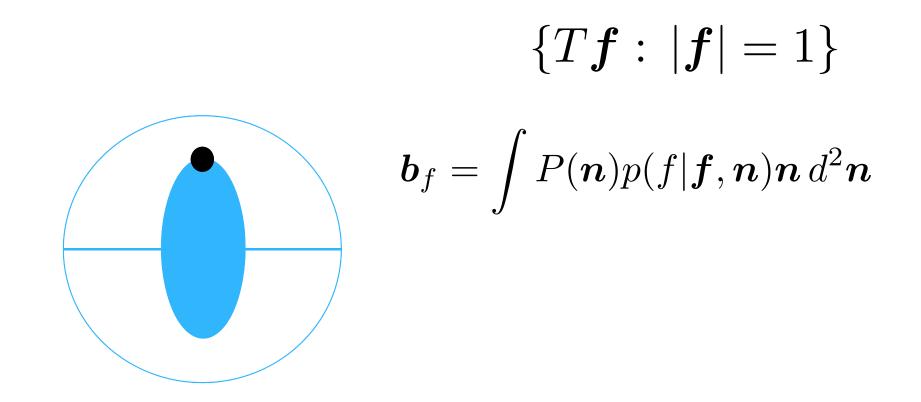
$\{Tf: |f| = 1\}$

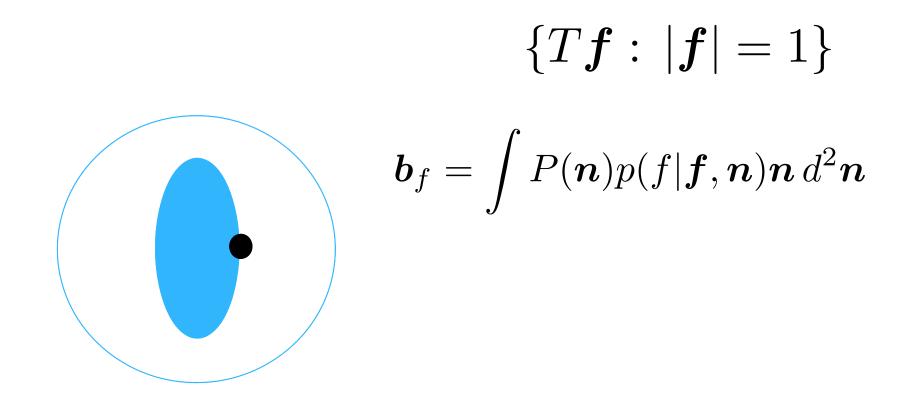


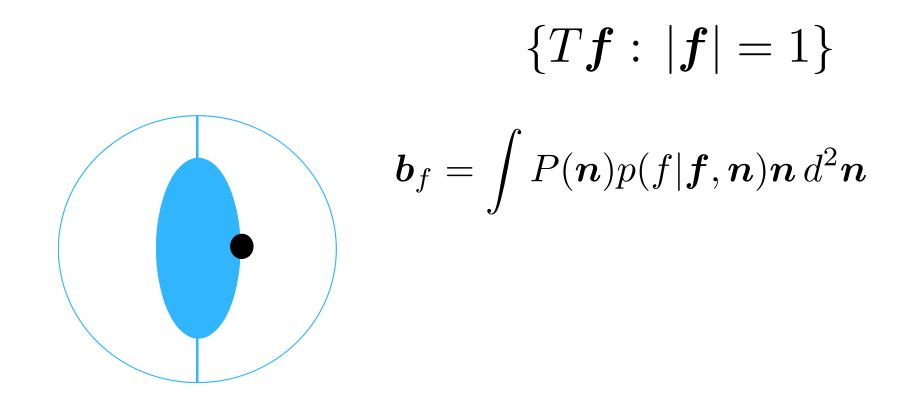
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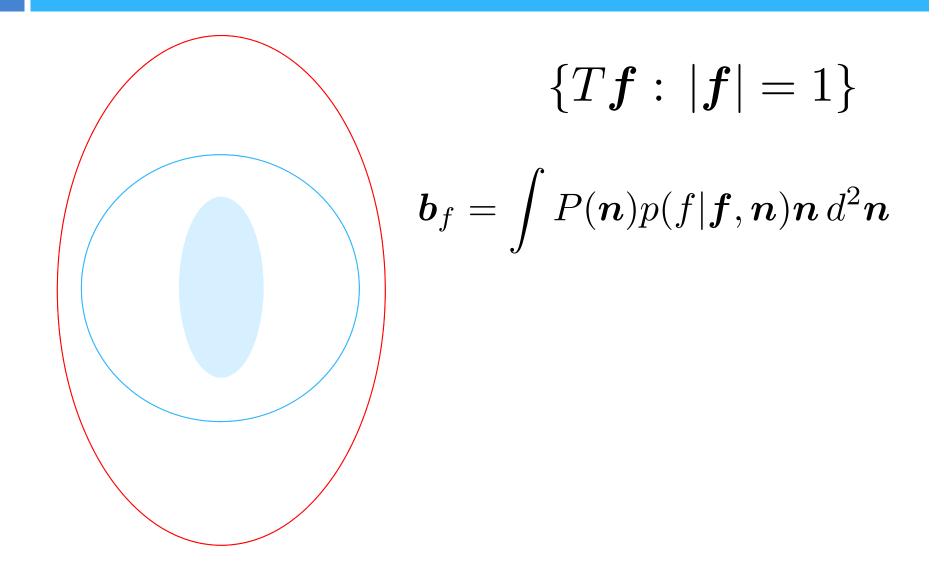


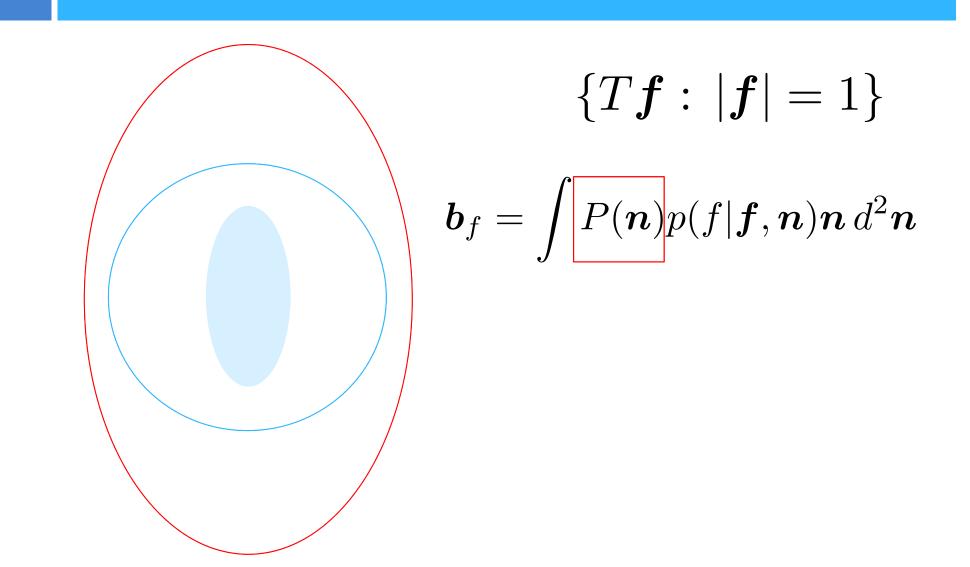


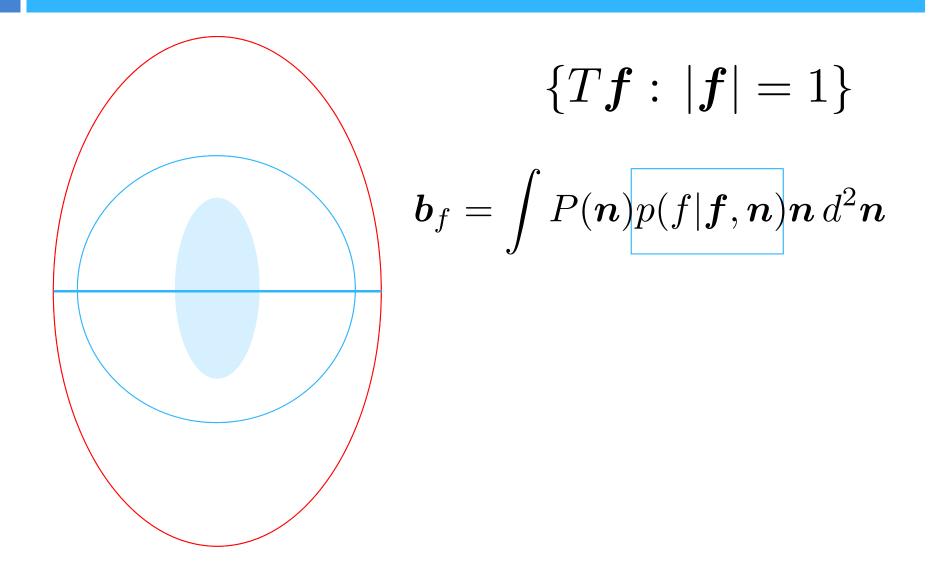










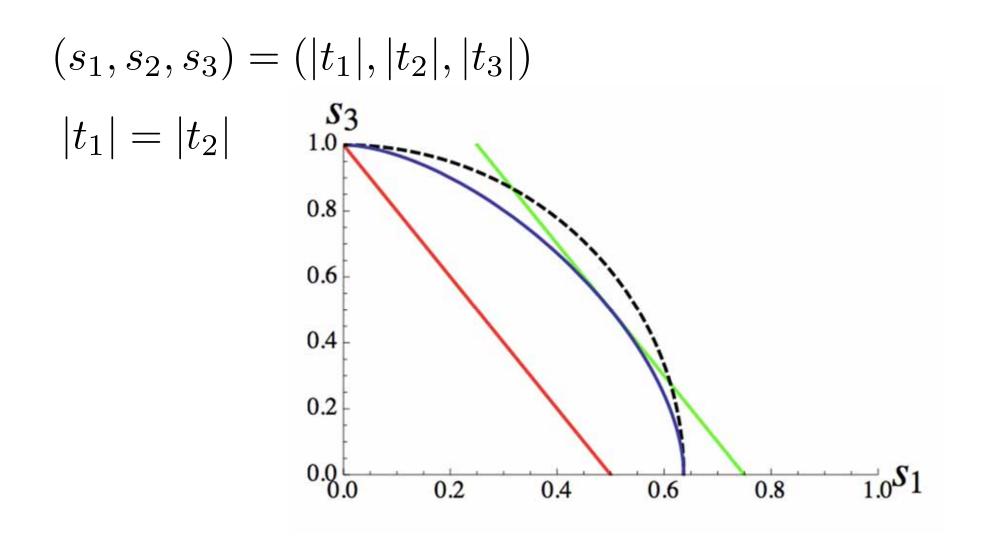


$$\{Tm{f}: |m{f}| = 1\}$$

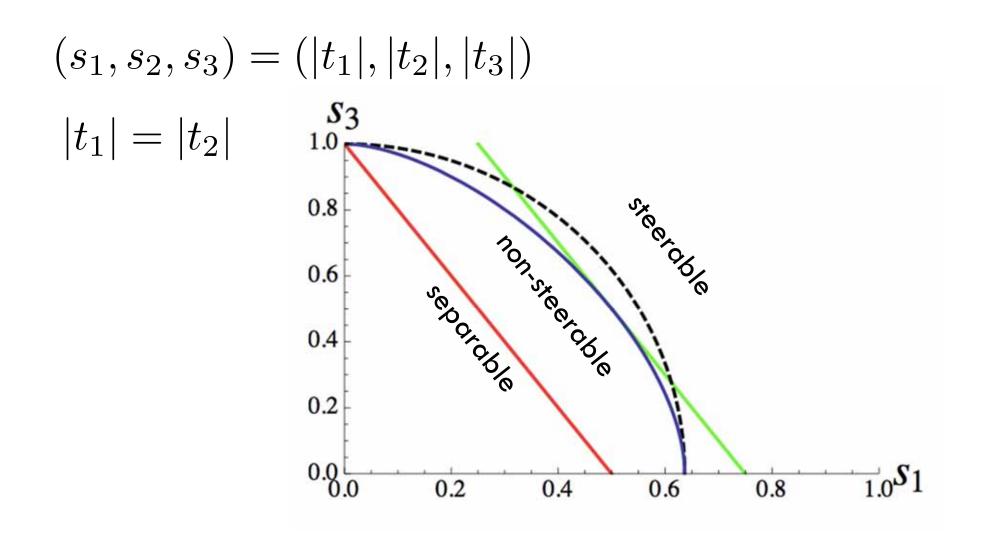
 $m{b}_f = \int P(m{n})p(f|m{f},m{n})m{n}\,d^2m{n}$
 $P(m{n}) \propto r(heta,\phi)^4$
 $ext{hemi}(m{f}) = \{m{n}: m{n}T^{-1}m{f} \ge 0\}$

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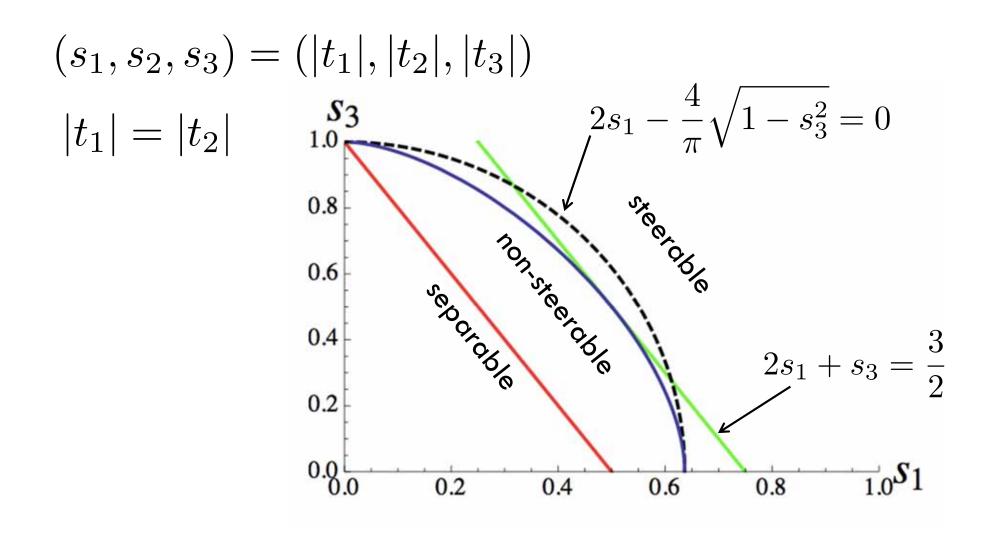
Surface of EPR-steerable T-states

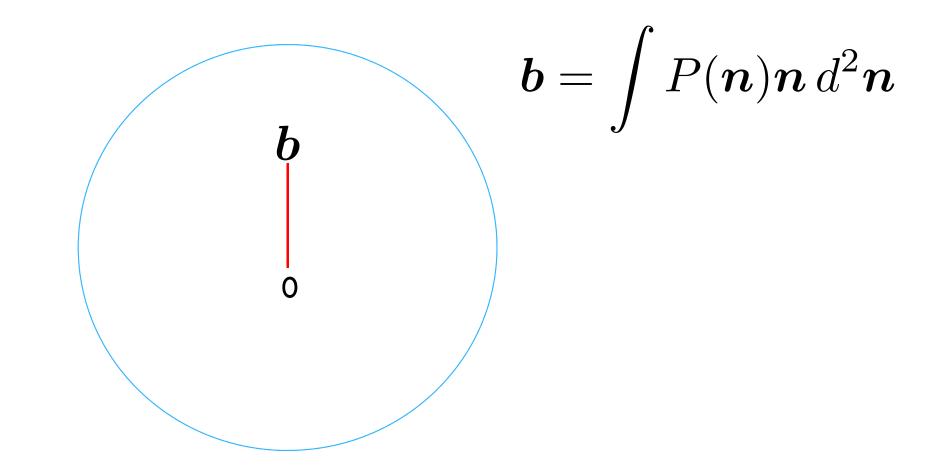


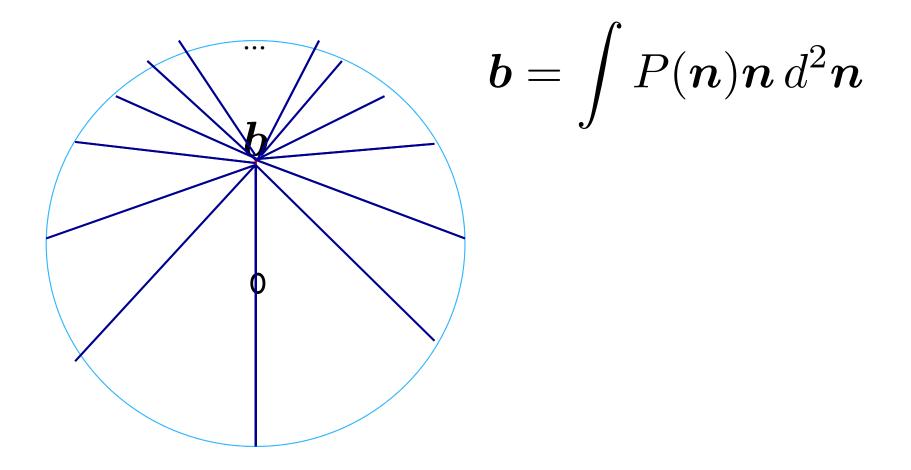
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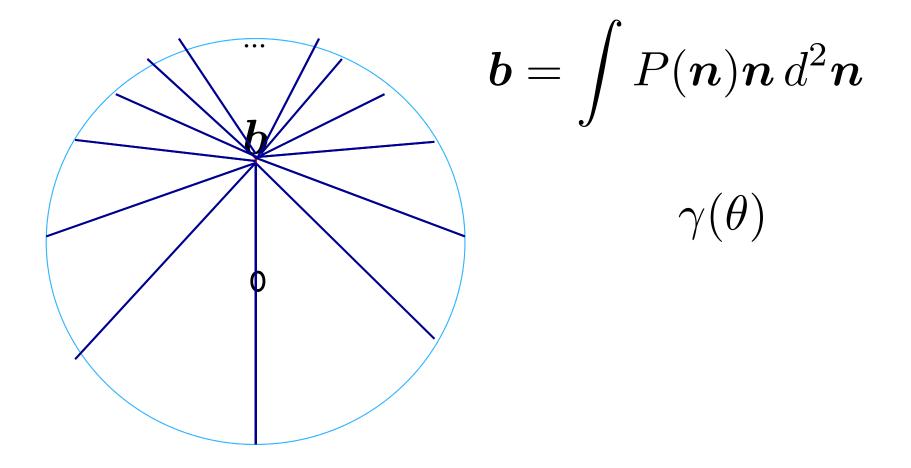


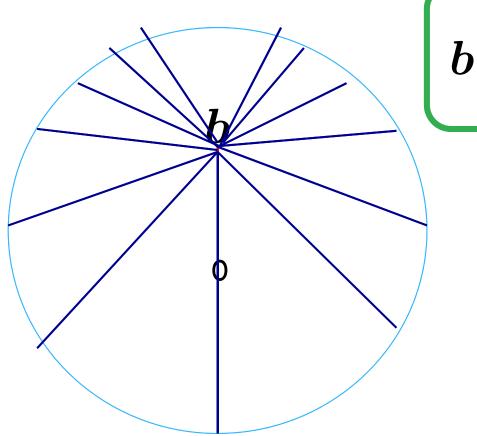
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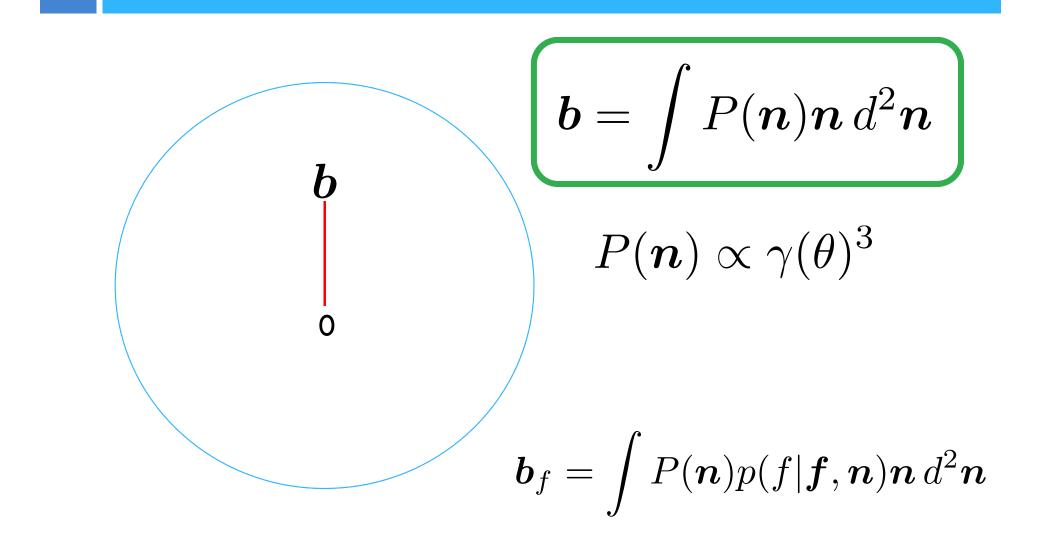


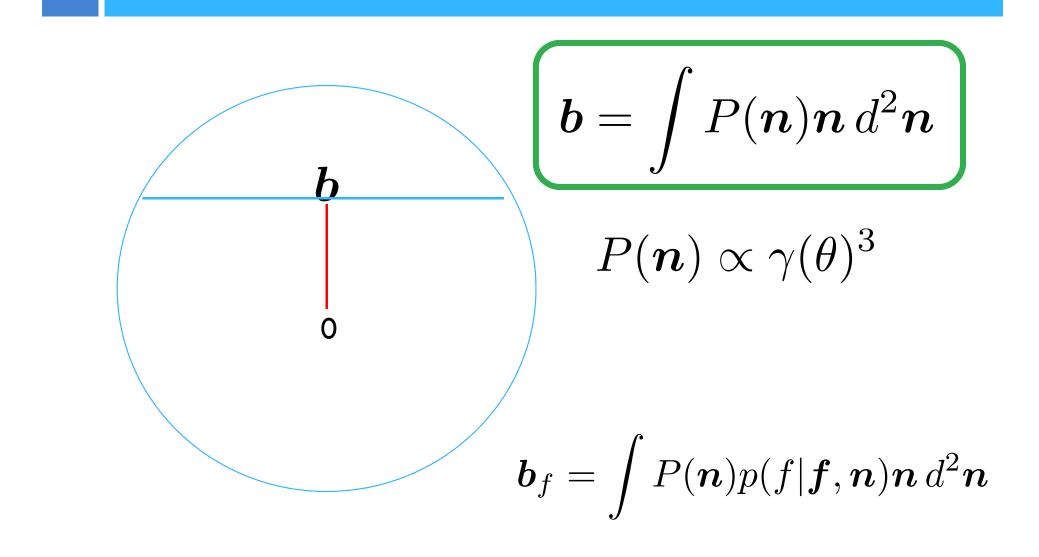




$$\boldsymbol{b} = \int P(\boldsymbol{n})\boldsymbol{n} \, d^2 \boldsymbol{n}$$

$$P(oldsymbol{n}) \propto \gamma(heta)^3$$





Further work

- For two-qubit states with maximally mixed reduced states, prove our LHS model based on the QSE is optimal
- Can QSE be useful for developing LHS models for general two-qubit states?
- What about LHS models for POVMs? Difficult already for Werner states, but QSE is depiction of all steered states, including POVM

Thanks!