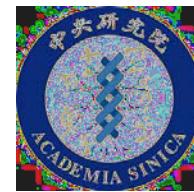
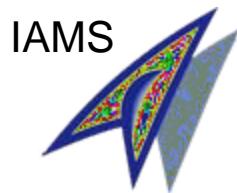


Toward generating synthetic gauge potentials for a Bose-Einstein condensate in a toroidal trap



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APCWQIS 2014, December 15, 2014, NCKU

Outline

- introduction of synthetic vector gauge potentials for ultracold neutral atoms: synthetic magnetic field and spin-orbit coupling
- experiment direction:
 - (i) generate synthetic magnetic flux Φ_B in a toroidal trap
 - (ii) detection atomic circulations induced by Φ_B
- We have achieved a ^{87}Rb Bose-Einstein condensate (BEC) with up to 3×10^5 atoms in a crossed optical dipole trap
- current experimental work and plan

Introduction: ultracold quantum gases

ultracold atoms: degenerate, non-classical gases

$$\frac{N}{V} \lambda_{dB}^3 \geq 1$$

$\frac{N}{V}$: density λ_{dB} : de Broglie wave length

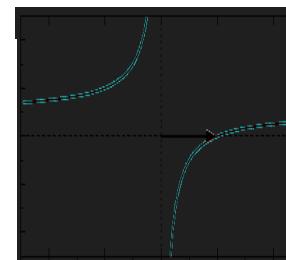
(1) cold and dilute: contact interaction

$$V(r - r') = \frac{4\pi\hbar^2 a}{m} \delta(r - r')$$

a : s-wave scattering length

(2) tunable interaction: Feshbach resonance

a vs. magnetic field B



(3) nearly disorder free

precisely controlled magnetic and optical potentials

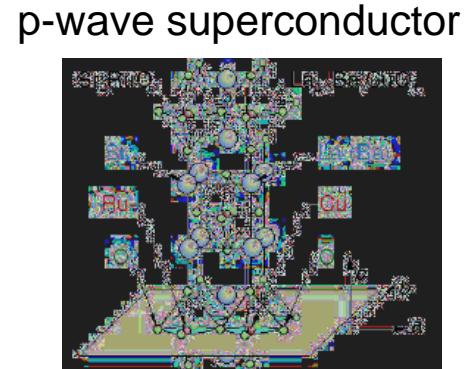
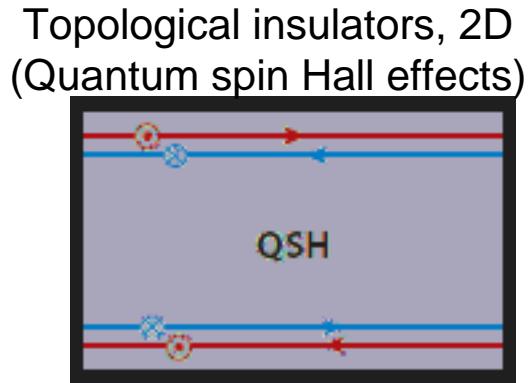
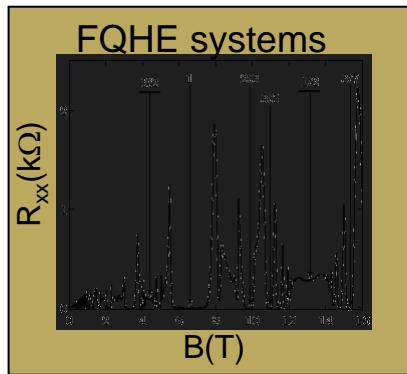
Zeeman shift

AC stark shift

→ ideal for quantum simulation:
model systems for condensed-matter physics

Quantum simulations of synthetic gauge potentials

to “charge” neutral atoms by creating a “ synthetic vector gauge potential A^* ”



- new approach to generate large B^* to study quantum-Hall physics

2D system and $v = N_{2D}/N_v \leq 1$
 $N_{2D} = \text{atom\#}, N_v = \# \text{ of flux quanta}$

- bosonic $v = 1$ state: w/ binary contact interaction,
nonabelian, for topological quantum computation

Ref: N. R. Cooper, 2008

- Spin-dependent $\vec{A}^*(\vec{\sigma})$: spin-orbit coupling
TR preserved topological insulators, topological superconductors:
nonabelian gauge potentials $\rightarrow [A_i^*, A_j^*] \neq 0$

Introduction of synthetic gauge potentials

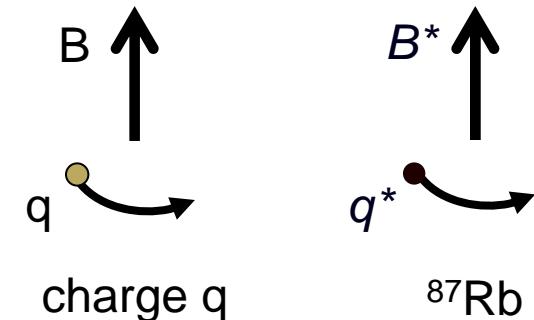
- Optically induced vector gauge potential A^* for neutral atoms:

$$H = \frac{(p - q^* A^*)^2}{2m^*} + V(x)$$

→ synthetic electric and magnetic fields

$$E^* = -\frac{\partial A^*}{\partial t}, B^* = \nabla \times A^*$$

- Create synthetic field B^* for neutral atoms:
effective Lorentz force $F = qv \times B$
to simulate charged-particles in real magnetic fields



- Light-induced potential to generate B^* in lab frame, no rotation of trap:
 - steady B^* , not metastable
 - easy to add optical lattices

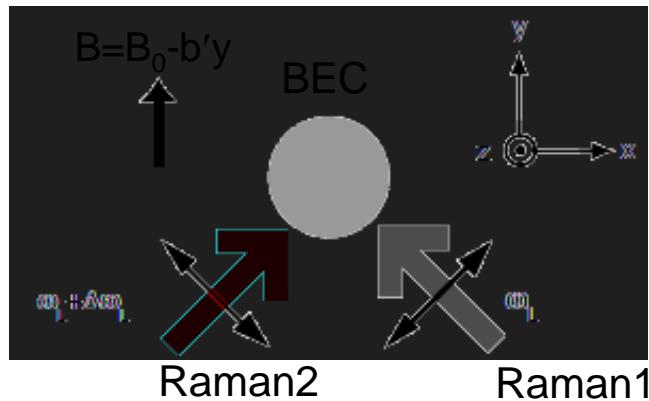
↓

B^* in rotating frame:
Coriolis force \leftrightarrow Lorentz force

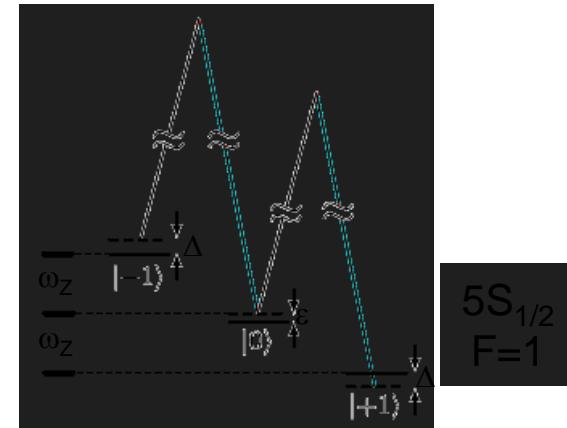
rotation: technical limit on B^*

Synthetic vector potentials (I): synthetic magnetic field B^*

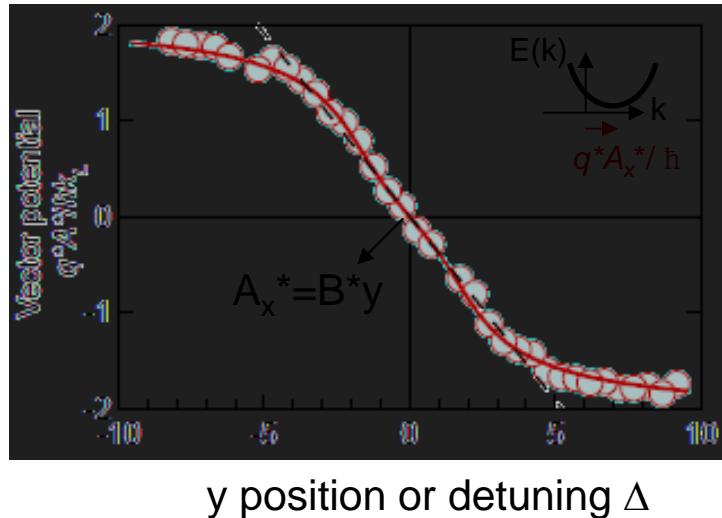
a. Raman-dressed BEC



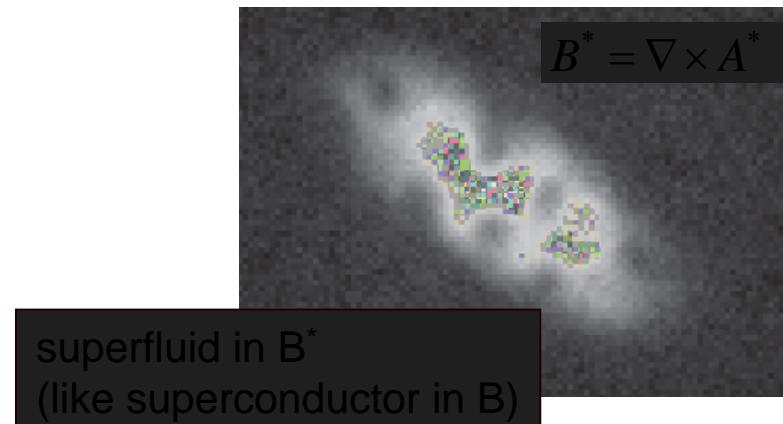
b. Level diagram



c. Vector potential A_x^* vs. position y



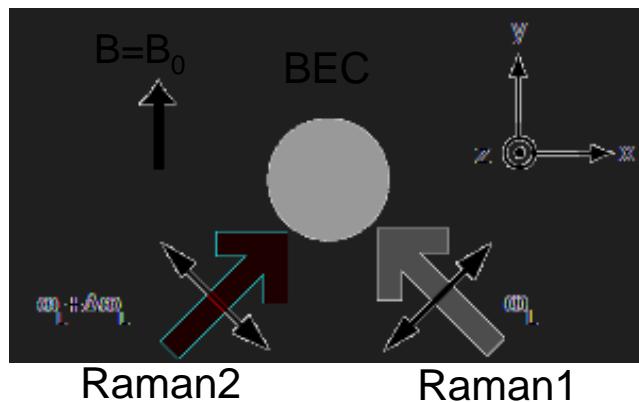
d. Synthetic magnetic field B^* , $A_x^* = B^*y$



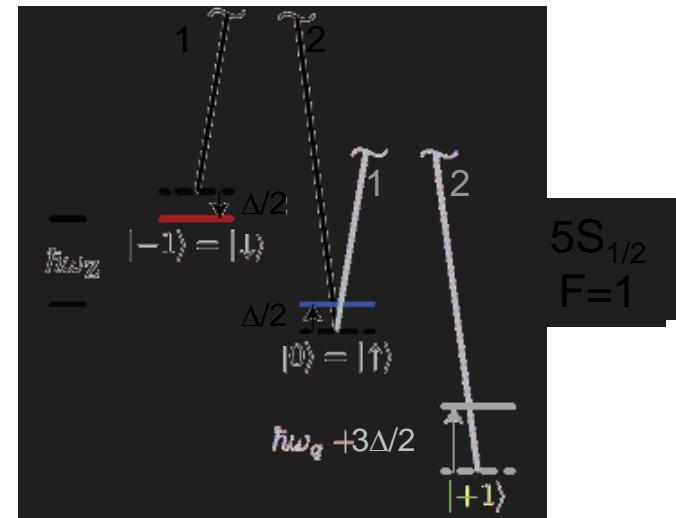
Reference: Y.-J. Lin et al., Nature **462**, 628 (2009),
Y.-J. Lin et al., PRL **102**, 130401 (2009).

Synthetic vector potentials (II): spin-orbit coupling

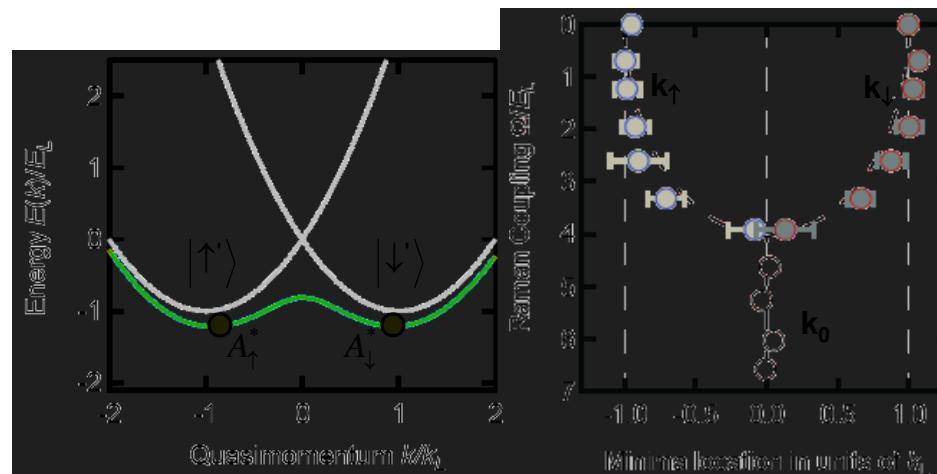
a. Raman-dressed BEC



b. Level diagram



c. Spin dependent A^* : spin-orbit coupling



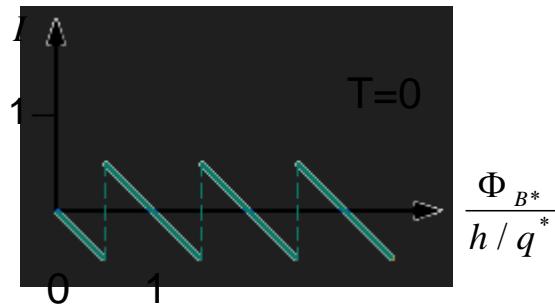
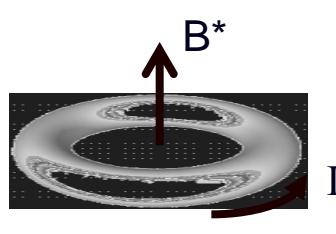
Reference: Y.-J. Lin, K.J.-Garcia and Ian Spielman, Nature **471**, 83 (2011).

Outline

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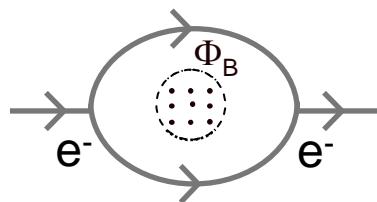
Atomic currents vs. magnetic flux

Equilibrium atomic currents I vs. B^* : expect Aharonov-Bohm oscillation



Aharonov-Bohm effect:

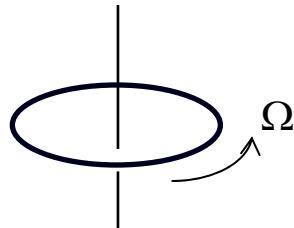
e^- in a 2-arm interferometer enclosing a magnetic flux Φ_B



$$\text{probability} \propto 1 + \cos\left(2\pi \frac{\Phi_B}{h/e}\right)$$

Atomic currents: analog between rotation and B^*

B^* = rotation and go to the rotating frame



$$L=2\pi r_0$$

1D ring, radius r_0 , rotating velocity Ω

rotation

B^*

velocity in the rotating frame \leftrightarrow velocity of q^* in B^*

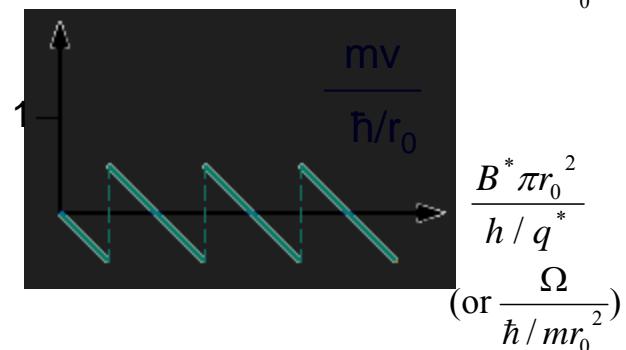
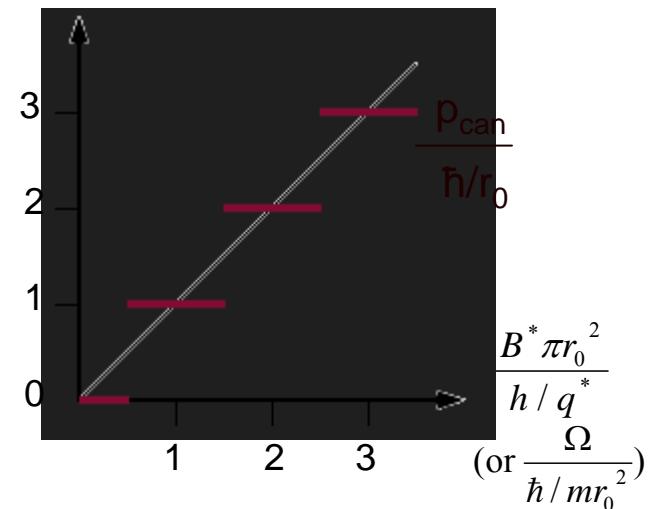
$$m\vec{v} = \vec{p}_{can} - m\Omega r_0 \hat{\phi} \leftrightarrow \vec{p}_{can} - \frac{q^* B^*}{2} r_0 \hat{\phi}$$

v = equilibrium current

= nonclassical rotational inertia

\rightarrow a measure of SF

P_{can} and mv vs. B^* (or rotation)



Traditional methods to create B^* : rotation

rotating neutral atoms

$$F_{\text{Coriolis}} = 2m\Omega v_{\text{rot}}$$

$$\Omega \leftrightarrow qB/2m$$

charge q in B

$$F_{\text{Lorentz}} = qvB$$

$$H_{\text{rot}} = \frac{\hbar^2}{2m} \left[\left(k_x - \frac{m\Omega y}{\hbar} \right)^2 + \left(k_y + \frac{m\Omega x}{\hbar} \right)^2 \right] + V(r)$$

$$V(r) = \frac{1}{2} m(\omega^2 - \Omega^2) r^2$$

$$H_B = \frac{\hbar^2}{2m} \left[\left(k_x - \frac{qBy}{2\hbar} \right)^2 + \left(k_y + \frac{qBx}{2\hbar} \right)^2 \right] + V(r)$$

$$V(r) = \frac{1}{2} m\omega^2 r^2$$

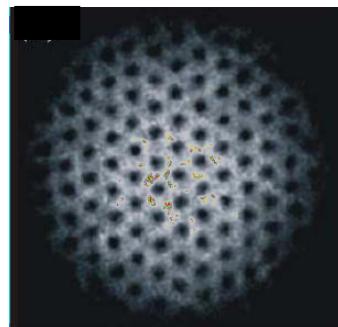
w/ mean field interaction

N_v vortices, $L/N=N_v/2$ (large N_v)

N_v vortices or flux quanta

(one vortex $\leftrightarrow \Phi_0 = h/q$)

rotating neutral BEC (experiment)



$\Omega/\omega = 0.975$, $R \sim 30\mu\text{m}$

Coddington et al., JILA, 2004

Generating magnetic flux

transfer angular momentum, modify dispersion $E(\ell)$

- Raman beams w/ OAM (orbital angular momentum), transfer $\Delta\ell$ to atoms
- shift of dispersion minimum: ℓ_{\min} (controlled by detuning)

$$1D: H = \frac{\hbar^2}{2mr_0^2} \left(\frac{\partial}{i\partial\phi} - \ell_{\min} \right)^2$$

$$\vec{p}_{can} - q^* A^* \hat{\phi}$$

$$\underline{r_0 q^* A^* = \hbar \ell_{\min}}$$

$$\text{flux } \Phi_B = \oint A^* dl = \hbar \ell_{\min}$$

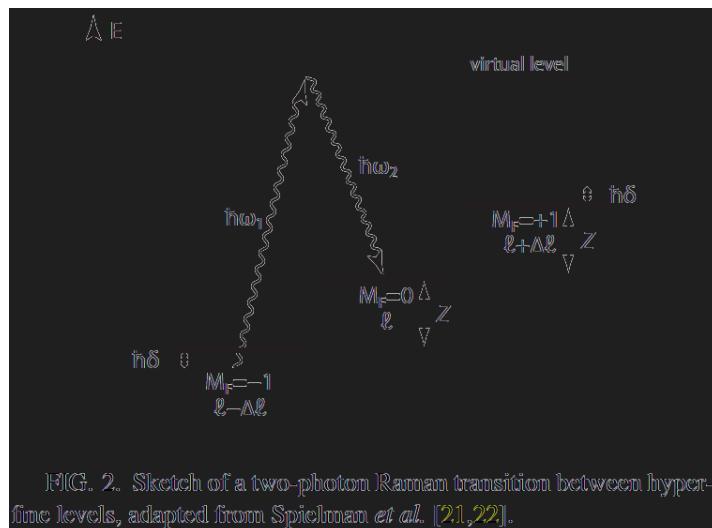
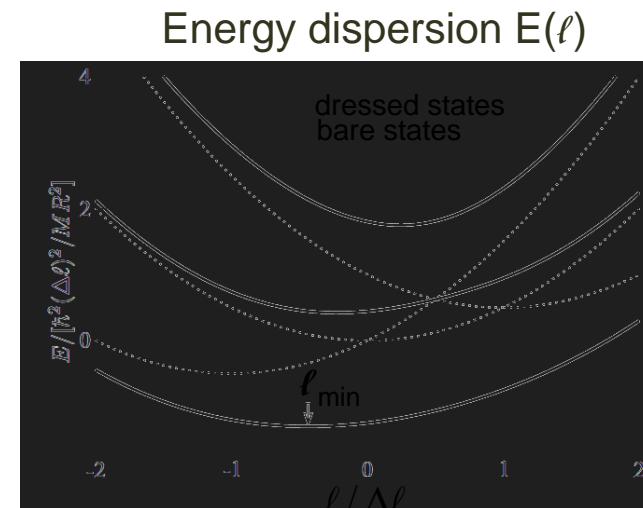


FIG. 2. Sketch of a two-photon Raman transition between hyperfine levels, adapted from Spielman *et al.* [21,22].



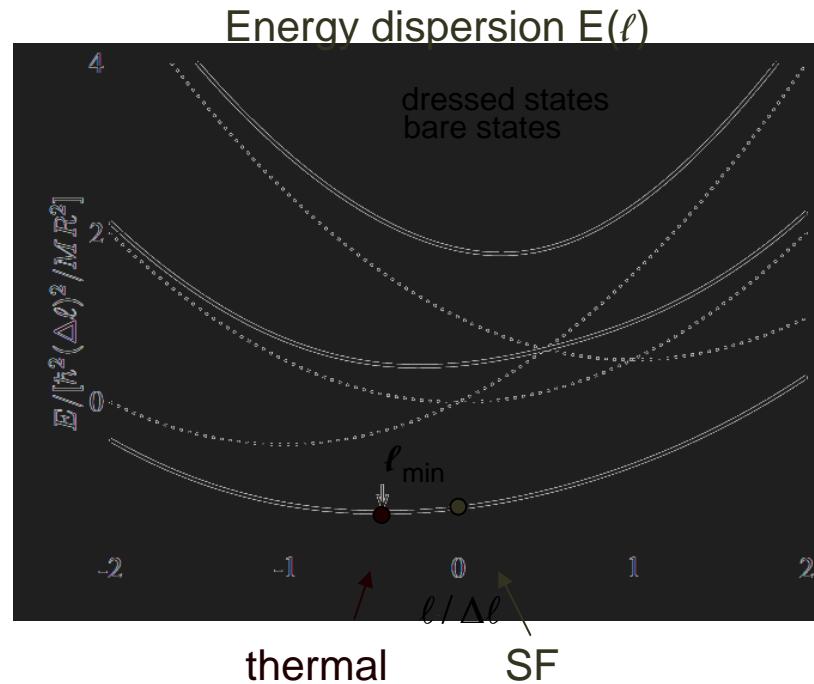
Ref: Lin et al., PRL (2009)

S.T.John, Z. Hadzibabic and N. R. Cooper, PRA (2011) (adapted figure source)

Detecting dressed atoms' angular momentum

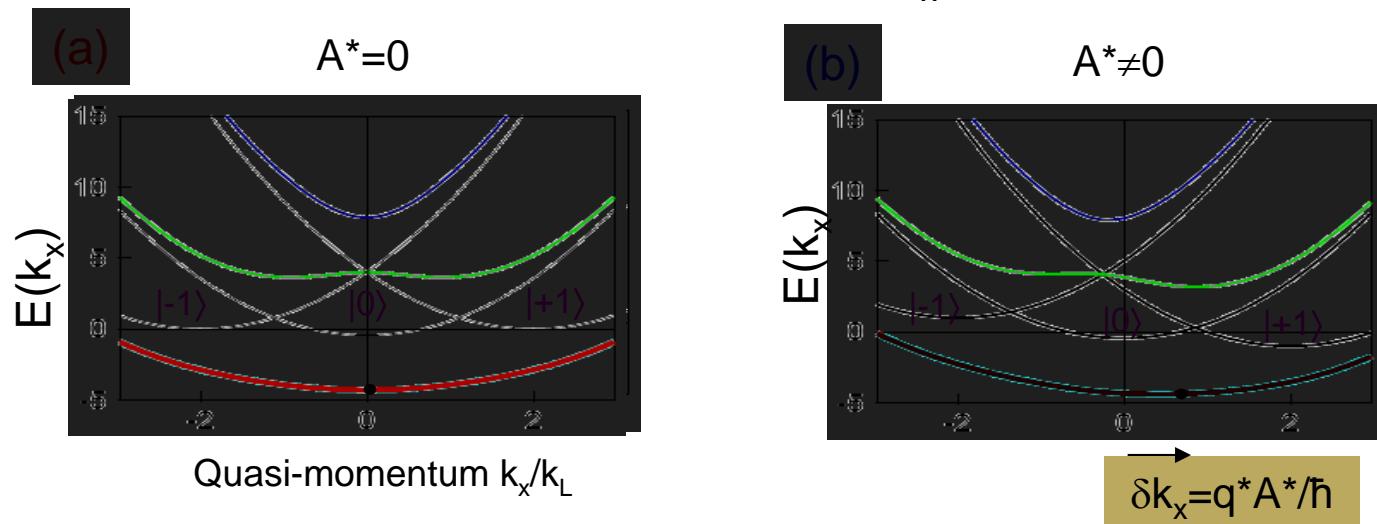
dressed state: $| -1, \ell - \Delta\ell \rangle, | 0, \ell \rangle, | +1, \ell + \Delta\ell \rangle$
spin $m_F = 0, \pm 1$

- dressed state $\langle \ell \rangle$ = population weighted average with m_F
= group velocity $\propto dE/d\ell$
- thermal atoms' ensemble average: $\ell = \ell_{\min}, v=0$
SF: $\ell=0, v \neq 0$

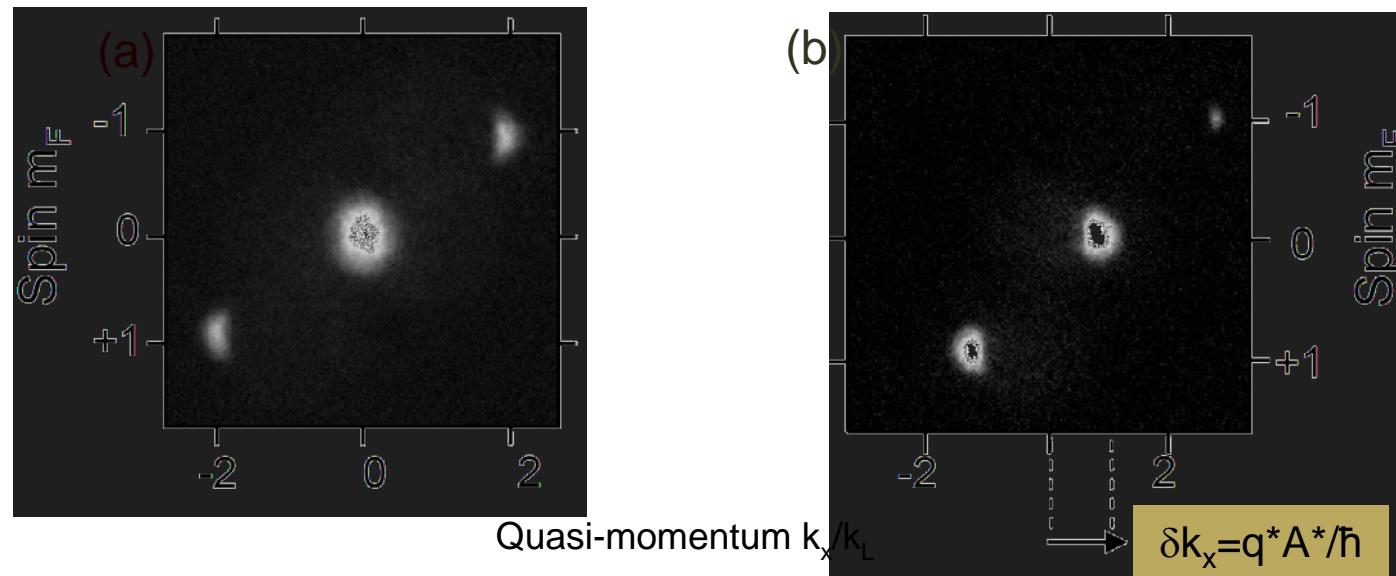


dressed state with uniform A^*

Energy dispersion $E(k_x)$



Spin-resolved Time-of-Flight (TOF) images

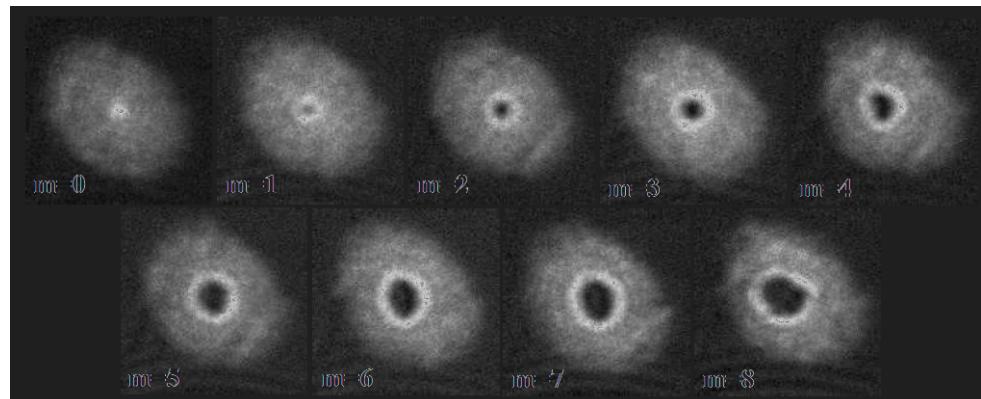
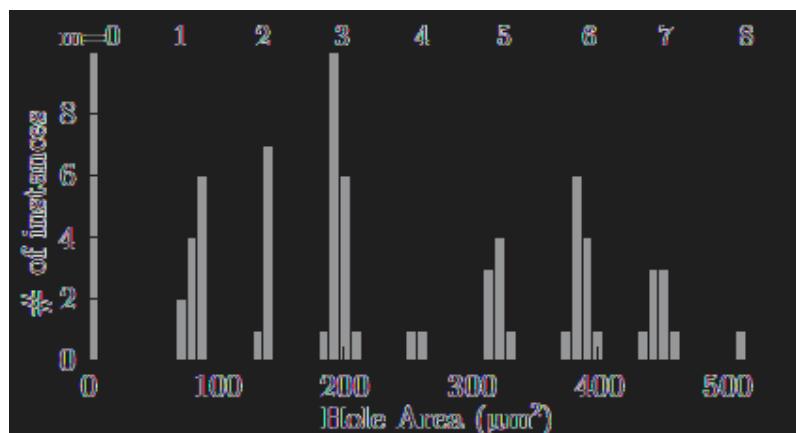


Detection of angular momentum ℓ

angular momentum: from hole size in TOF images

circulation of stirred BEC determined by the hole sizes in 10ms TOF
($\ell \rightarrow m$)

ramp down → release → TOF

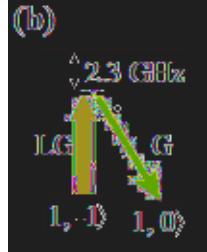
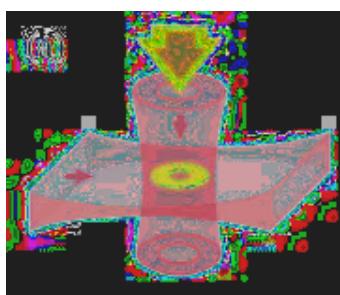


Ref: Murray et al, arxiv cond-mat.quat-gas 1309.2257v1
experiment: NIST Na ring BEC

Setup: ring trap geometry

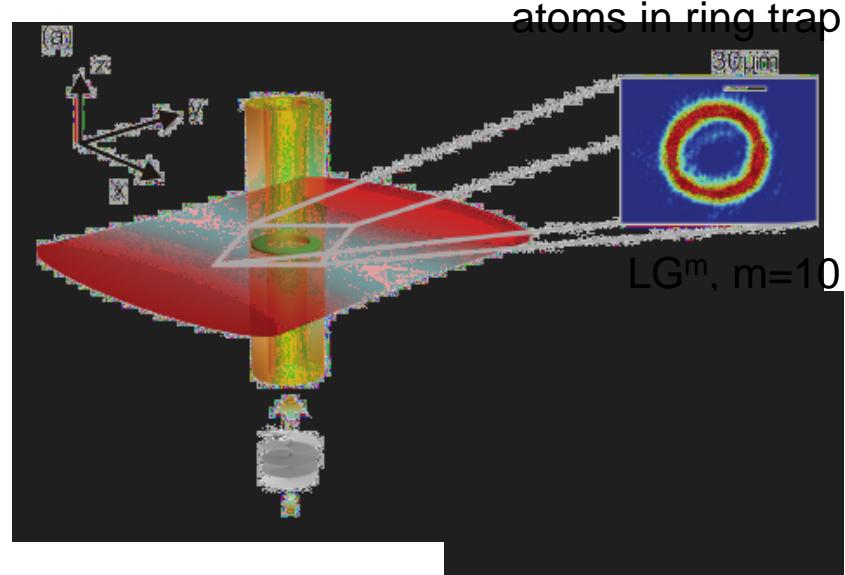
Typical setup:
light sheet + ring beam

NIST Na ring BEC



Ref: Ramanthan et al., PRL (2011)

Cambridge group Rb ring BEC



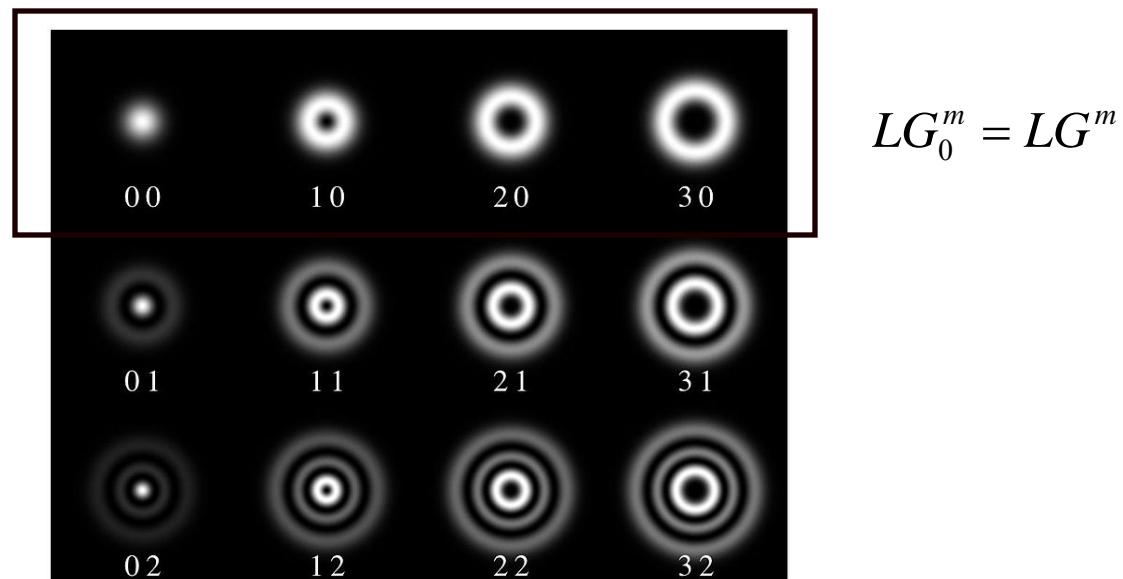
Ref: Z. Hadzibabic group, PRA (2012)
PRL(2013)

Laguerre-Gaussian beam LG^m

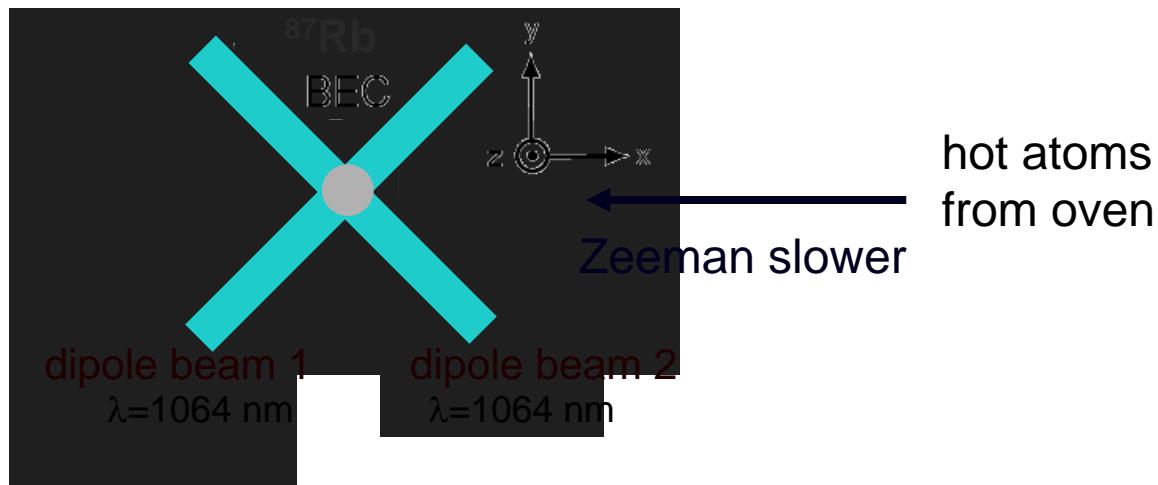
Laguerre-Gaussian beams LG_n^m m= phase winding
n= radial index



$$LG_0^m = LG^m \propto r^m e^{-r^2/w^2} e^{im\phi}$$



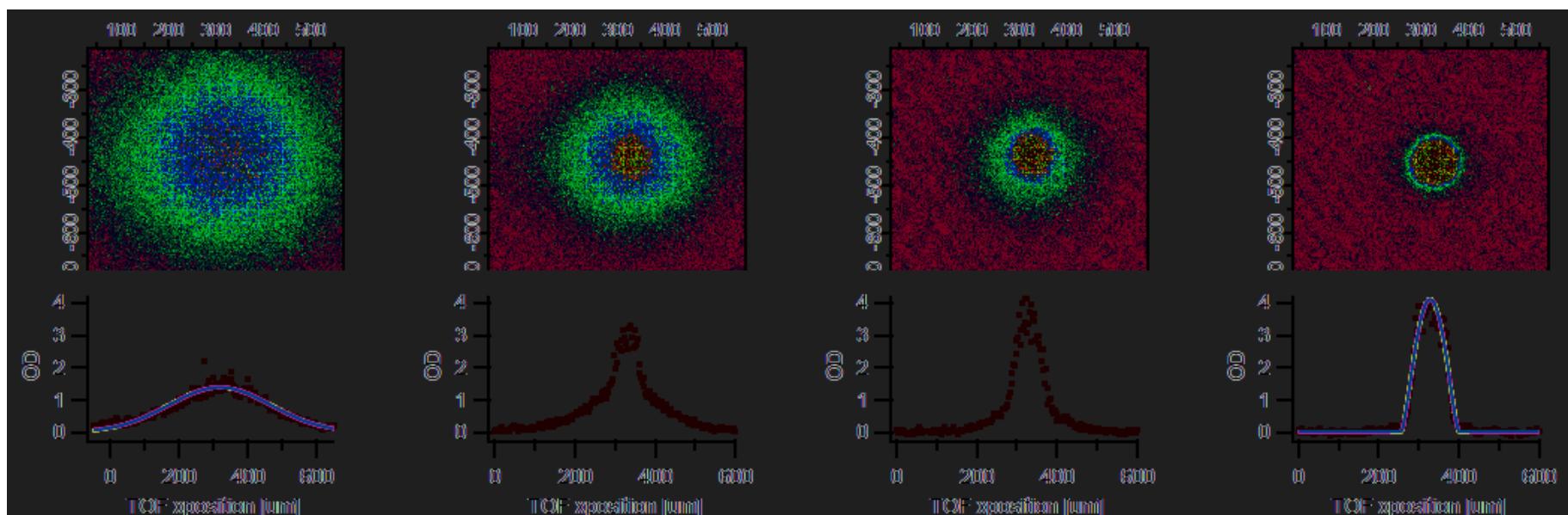
Experiment: BEC production



- load Magneto-Optical trap (MOT) from Zeeman slower: $\sim 8 \times 10^8$ atoms in 3 s
- polarization gradient cooling $>\sim 10$ ms to $\sim 8 \mu\text{K}$
- rf-evaporative cooling in a quadrupole magnetic trap for 3.5 s, $|F=1, m_F = -1\rangle$
- crossed optical dipole trap + weak magnetic trap:
evaporate in hybrid potential for $\sim 7\text{s}$
→ ramp off magnetic gradient in 2.5s
→ BEC in crossed dipole trap: 3×10^5 atoms
- total cycle time ~ 20 s

Thermal atoms to BEC

20ms Time of Flight



thermal

bimodal

cooling

BEC
atom number
 $N \sim 10^5$

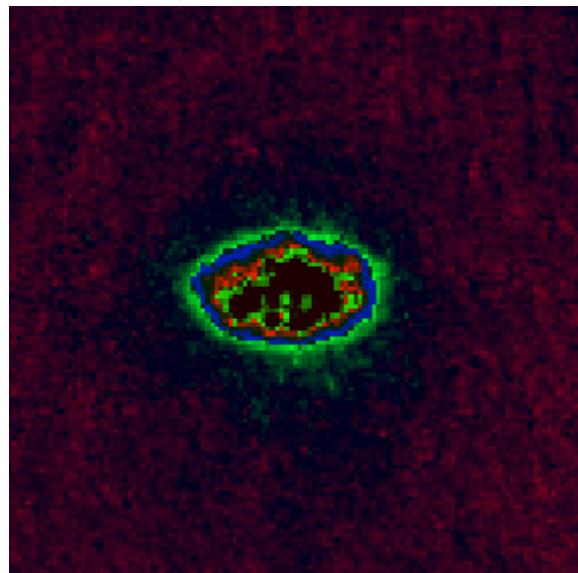


Optimized BEC

BEC number $\sim 3 \times 10^5$

TOF= 30ms

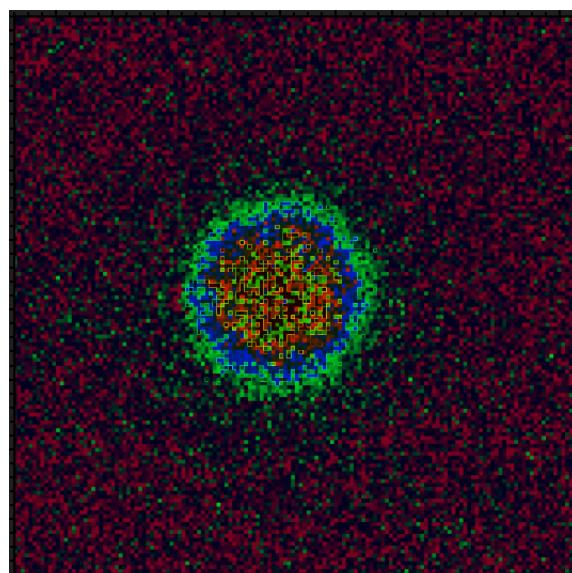
y



x

TOF= 20ms

z



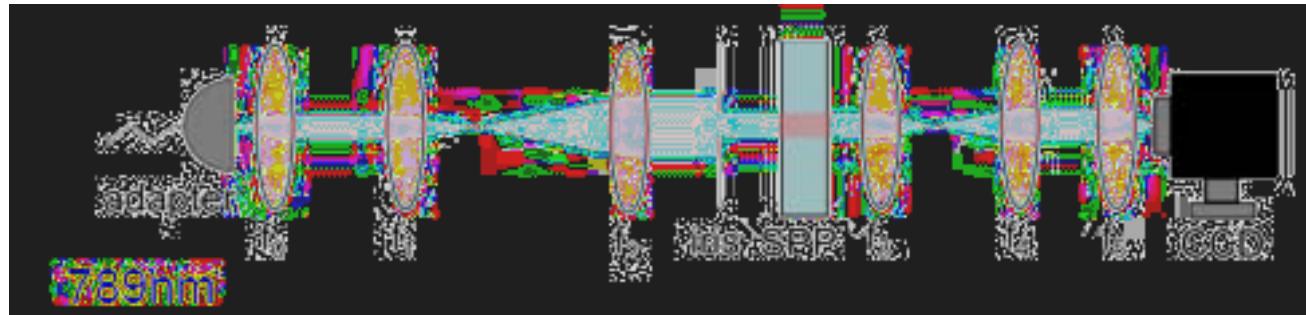
x

100um

Experiment: generating magnetic flux

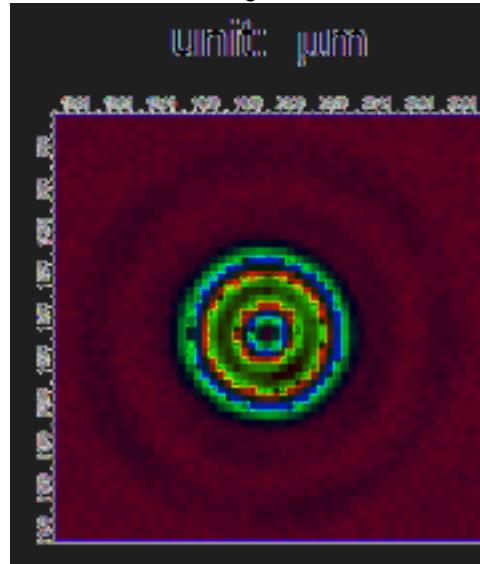
Raman beam with OAM= $\ell\hbar$

spiral phase plate (SPP): generate phase winding $e^{i\ell\phi}$

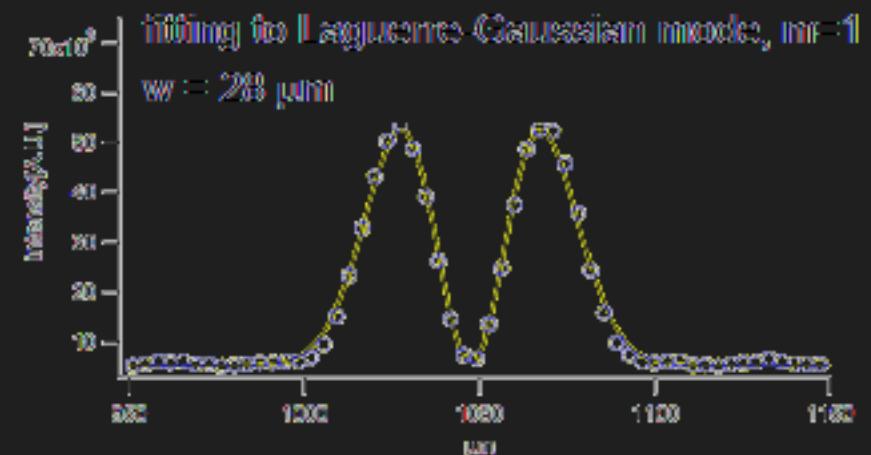


CCD image

ring radius $r_0 \sim 20\text{ }\mu\text{m}$ $\ell=1$



$$\mathcal{LC}_{\ell=1}^{\text{ring}} := \mathcal{LC}_r^{\text{ring}} \ll r^m \approx r^2 \ln^2 \ell \approx \mu\text{m}^2$$

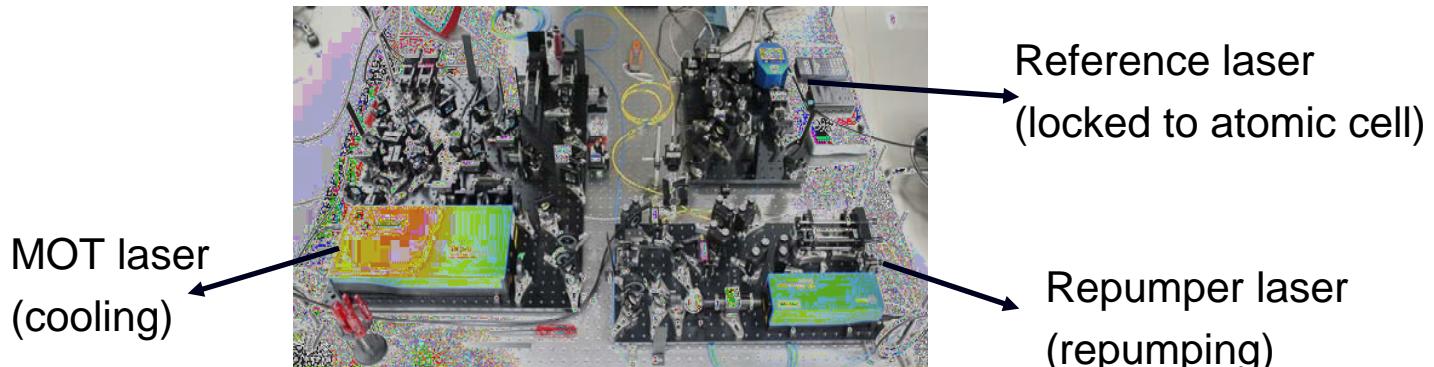


Ongoing work

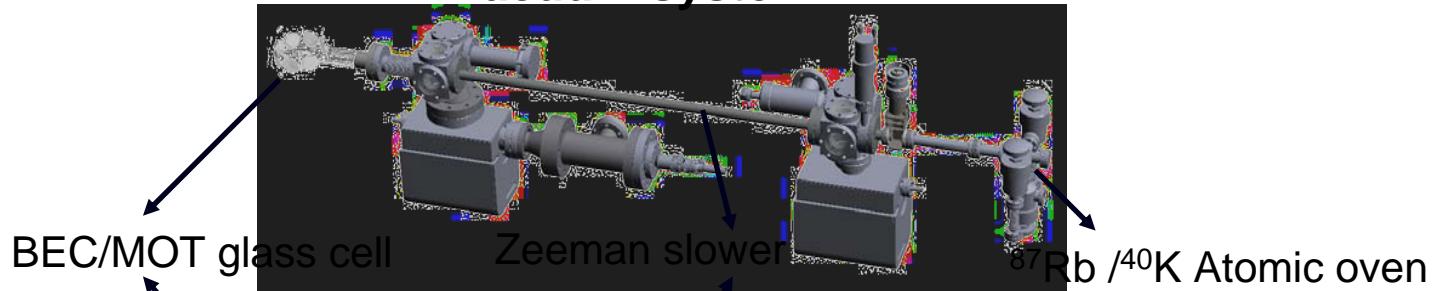
- loading atoms into the ring trap: crossed light sheet beam+ ring beam
- making smooth ring-shaped trapping potentials, and combining multiple beams
- stability of the experiment
- finite temperature effects, detection of thermal atoms

Experimental setup

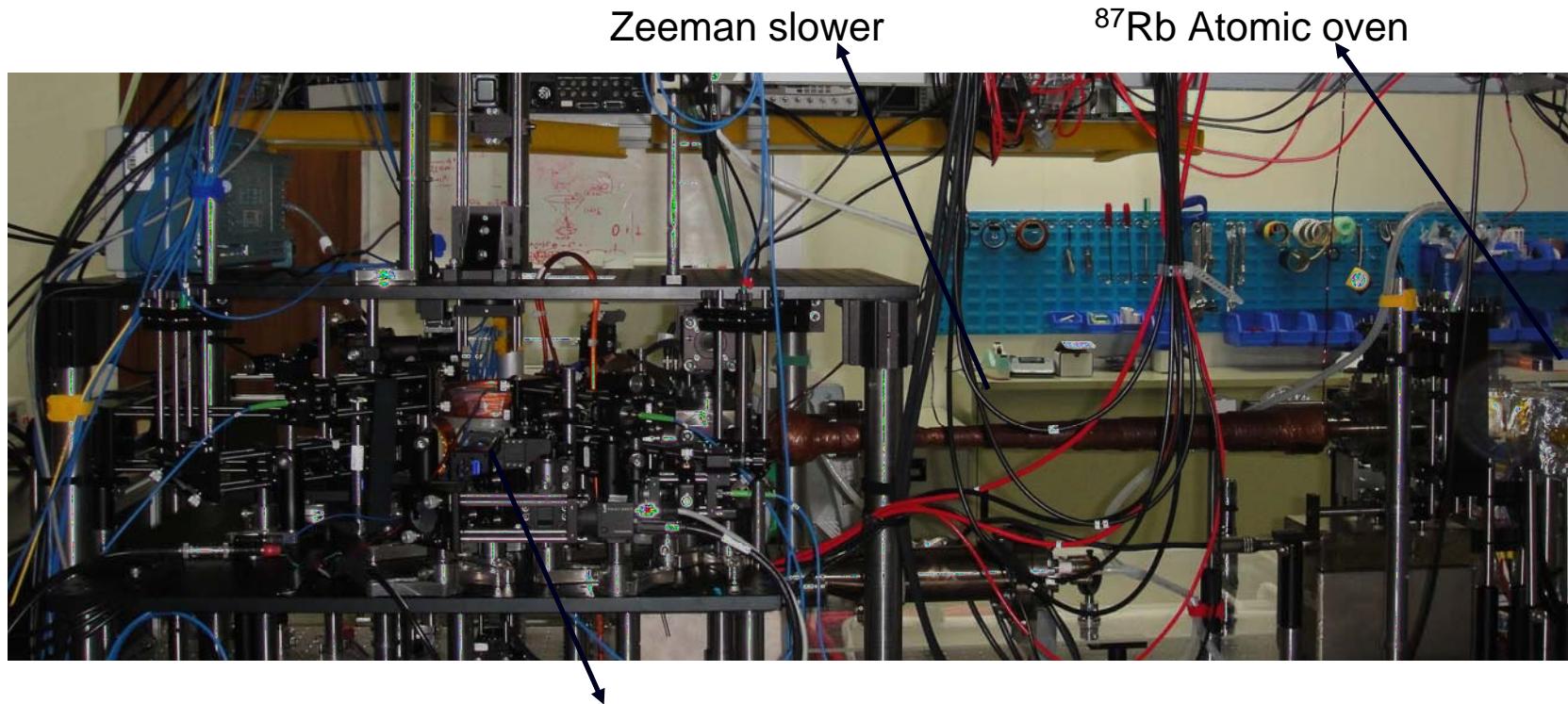
Diode lasers for laser cooling



Vacuum system



Experimental setup



Acknowledgements

NIST

Karina Jimenez-Garcia

Robert Compton

James Trey Porto

William Phillips

Ian Spielman

IAMS

Cheng-An Chen

Pan-Pan Huang

Jung-Bin Wang

Chin-Yeh Yu

Gergely Imreh

IAMS



科 技 部

Ministry of Science and Technology

Summary

- synthetic vector gauge potentials for ultracold neutral atoms
- have achieved ^{87}Rb BEC with up to 3×10^5 atoms in a crossed dipole trap
- next: produce BEC in a ring trap
- toward generating atomic circulation from a magnetic flux Φ_B and A^* by Raman beams with OAM