

*General analytical solution to
quantum states
of
open nanoelectronic systems*

Matisse Wei Yuan Tu

Collaborators:

Wei Min Zhang, Amnon Aharony, Ora Entin Wohlman
Jiang Heng Liu, and Avraham Schiller

in preparation

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I. Introduction and motivation

quantum states

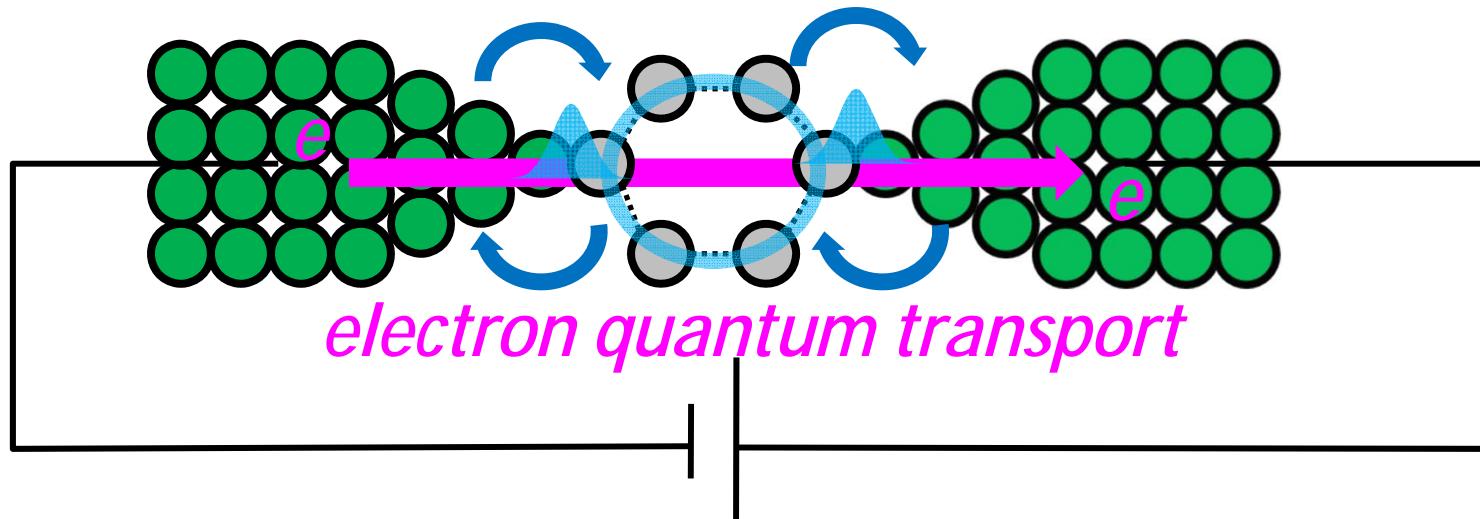
preparation

manipulation

detection

quantum information processing

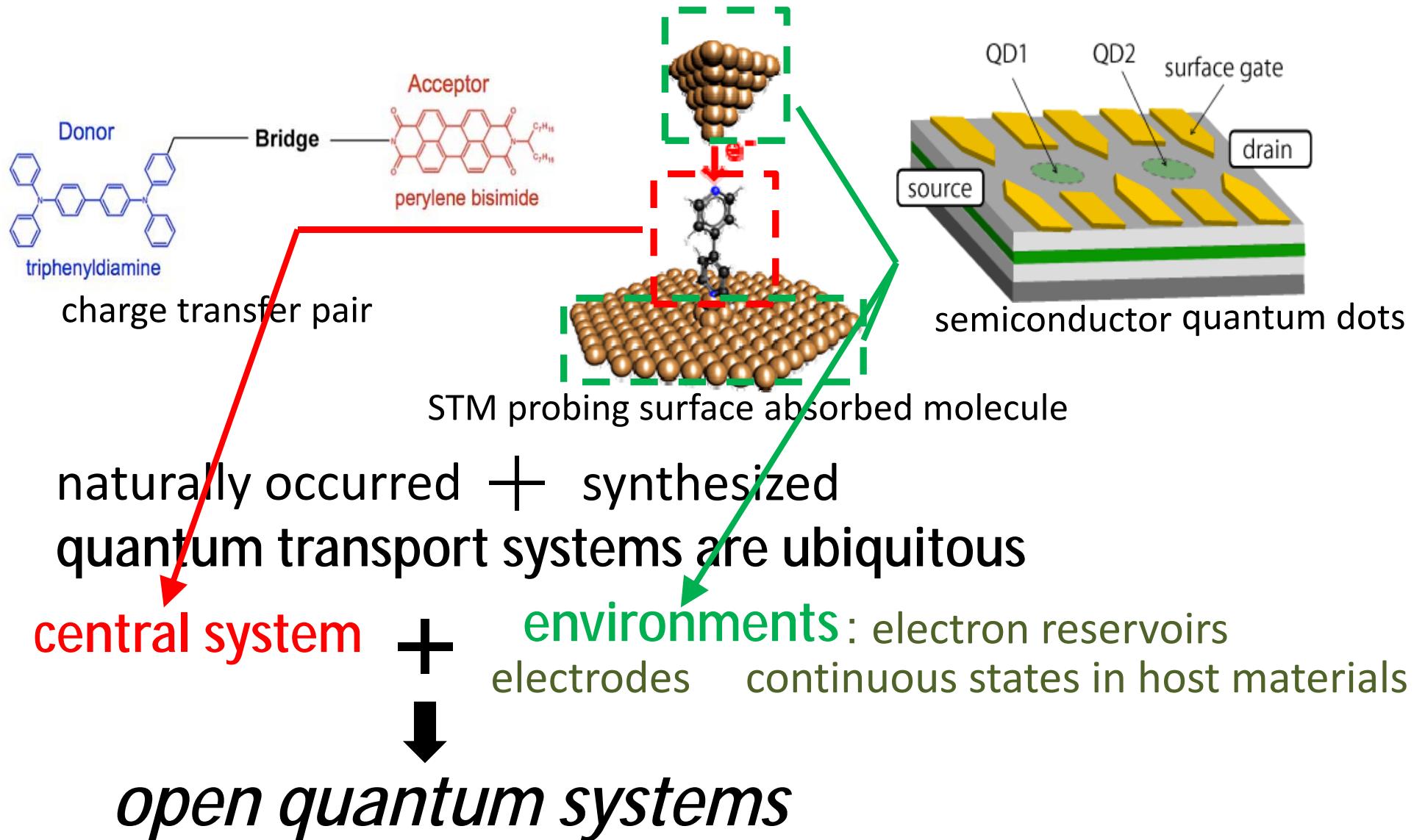
nanoelectronic systems



- various tunable couplings
- adjustable energetic structures
- useful as quantum devices

I. Introduction and motivation

quantum transport and nanoelectronics

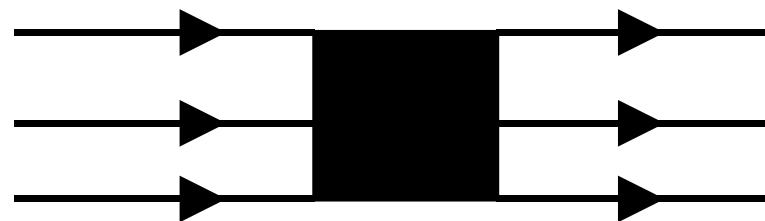


I. Introduction and motivation

theories for nanoelectronic transport setups

- *Landauer Büttiker formalism*

- in terms of scattering amplitudes
- targeting at in-out relation
- the central system as a black box



- *Schwinger Keldysh nonequilibrium Green function technique*

- in terms of Green functions
- targeting at the transport currents
- the state of the central system is averaged

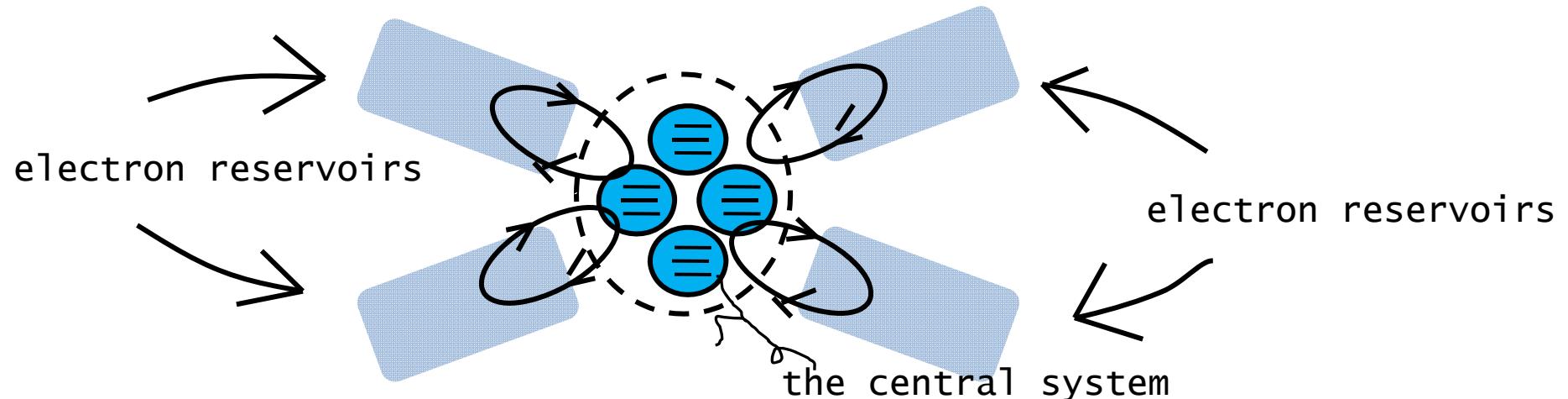


- *Feynman Vernon's influence functional theory*

- in terms of the reduced density matrix
- targeting at the quantum state of open systems
- fully characterizing the state of the central system

I. Introduction and motivation

modeling of quantum transport devices



$$H(t) = H_{dq}(t) + H_{rev}(t) + H_T(t)$$

$$\rightarrow H_{qd}(t) = \sum_{ij} \epsilon_{ij}(t) d_i^\dagger d_j$$

the central system

$$\rightarrow H_{rev}(t) = \sum_{\alpha k} \epsilon_{\alpha k}(t) c_{\alpha k}^\dagger c_{\alpha k}$$

electron reservoirs

$$\rightarrow H_T(t) = \sum_{i \alpha k} [V_{i \alpha k}(t) d_i^\dagger c_{\alpha k} + V_{i \alpha k}^*(t) c_{\alpha k}^\dagger d_i]$$

electron exchange

I. Introduction and motivation

exact master equation for a class of nanoelectronic systems

- equation of motion for the reduced density operator

$$\frac{d}{dt}\hat{\rho}(t) = -i[\hat{H}_{\text{dq}}(t), \hat{\rho}(t)] + \sum_{\alpha} [\mathcal{L}_{\alpha}^{+}(t) + \mathcal{L}_{\alpha}^{-}(t)]\hat{\rho}(t)$$

- transport currents

$$I_{\alpha}(t) = \frac{e}{2\hbar} \text{tr}_s [(\mathcal{L}_{\alpha}^{+}(t) - \mathcal{L}_{\alpha}^{-}(t))\hat{\rho}(t)]$$

non-unitary processes due to
electron exchanges

Starting theory:

M.W.Y. Tu and W.M. Zhang, PRB **78**, 235311 (2008);

M.W.Y. Tu, M.T. Lee, and W.M. Zhang, QIP (Springer) **8**, 631 (2009).

J.S. Jin, M.W.Y. Tu, W. M. Zhang, and Y.J. Yan, NJP, **12**, 083013 (2010).

Applications:

M.W.Y. Tu, W.M. Zhang, and J.S. Jin, Phys. Rev. B **83**, 115318 (2011)

H.N. Xiong, W. M. Zhang, M.W.Y. Tu, and D. Braun, PRA **86**, 032107(2012).

W.M. Zhang, L.P. Yuan, H.N. Xiong, M.W.Y. Tu, and F. Nori, PRL **109**, 170402(2012).

M.W.Y. Tu, W.M. Zhang, J.S. Jin, O. Entin-Wohlman, and A. Aharony, PRB **86**, 115453(2012).

M.W.Y. Tu, W. M. Zhang, and F. Nori, PRB **86**, 195403(2012).

J.S. Jin, M.W.Y. Tu, N.E. Wang and W.M. Zhang, J. Chem. Phys. **139**, 064706 (2013).

M.W.Y. Tu, A. Aharony, W.M. Zhang, O. Entin-Wohlman, PRB **90**, 165422 (2014).

↳ applied only using initial empty state, **not full use** of this master equation
numerically solving master equation is difficult, limited to small dimension

general analytical solution ??

useful for other more applications

II. Quantum states of open electronic systems

general analytical solution to exact master equation

$$\langle Pl^{(n)} | \hat{\rho}(t) | Pr^{(n)} \rangle = \det [1_D - v(t)] \sum_{m=0}^D \sum_{Pa^{(m)}} \sum_{Pb^{(m)}} \left[J_{a^{(m)}, b^{(m)}}^{l^{(n)}, r^{(n)}}(t) \langle Pa^{(m)} | \hat{\rho}(t_0) | Pb^{(m)} \rangle \right]$$

non unitary evolutions

initial state
environments initially
at thermal equilibria

later state

$l^{(n)}$: n-electron
 $r^{(n)}$: configurations

$a^{(m)}$: m-electron
 $b^{(m)}$: configurations

notations for configurations

D : number of available orbitals in the system

$S = \{1, 2, \dots, D\}$: set of available orbitals in the system

$\mathbf{i}^{(n)} = \{i_1, \dots, i_n\}$: set of n occupied orbitals as a configuration

notations for multi-electron states

$|P\mathbf{i}^{(n)}\rangle = d_{P\mathbf{i}_n^{(n)}}^\dagger \cdots d_{P\mathbf{i}_1^{(n)}}^\dagger |0\rangle$: an n -electron state

$P\mathbf{i}_n^{(n)} > \dots > P\mathbf{i}_2^{(n)} > P\mathbf{i}_1^{(n)} \in \mathbf{i}^{(n)}$

↑ fixing ordering to avoid fermion exchange sign problem

II. Quantum states of open electronic systems

elementary processes behind the open system evolutions

- single particle propagations

$$\frac{\partial}{\partial t} \mathbf{u}(t, \tau) + i\epsilon(t)\mathbf{u}(t, \tau) + \int_{\tau}^t ds \sum_{\alpha} g_{\alpha}(t, s)\mathbf{u}(s, \tau) = 0$$

- statistical actions from reservoirs on filling the central system

$$v(t) = \int_{t_0}^t d\tau \int_{t_0}^t d\tau' \mathbf{u}(t, \tau) \sum_{\alpha} \tilde{g}_{\alpha}(\tau, \tau') [\mathbf{u}(t, \tau')]^\dagger = \langle d_j^\dagger(t) d_i(t) \rangle$$

initial
empty state

$$g_{\alpha}(t_1, t_2) = \int \frac{d\omega}{2\pi} \Gamma_{\alpha}(\omega, t_1, t_2) e^{-i\omega(t_1 - t_2)}$$

dissipation kernel

$$\tilde{g}_{\alpha}(t_1, t_2) = \int \frac{d\omega}{2\pi} f_{\alpha}(\omega) \Gamma_{\alpha}(\omega, t_1, t_2) e^{-i\omega(t_1 - t_2)}$$

fluctuation kernel

fermi function

$$[\Gamma_{\alpha}(\omega, t_1, t_2)]_{ij} = 2\pi \sum_{k \in \alpha} V_{i\alpha k}(t_1) \exp \left[-i \int_{t_2}^{t_1} ds \Delta_{\alpha k}(s) \right] V_{\alpha k j}(t_2) \delta(\omega - \epsilon_{\alpha k}^0)$$

spectral density

tunneling amplitudes

$$\epsilon_{\alpha k}(t) = \epsilon_{\alpha k}^0 + \Delta_{\alpha k}(t) : \text{time-dependent energies of reservoir states}$$

II. Quantum states of open electronic systems *propagating function in Fock space*

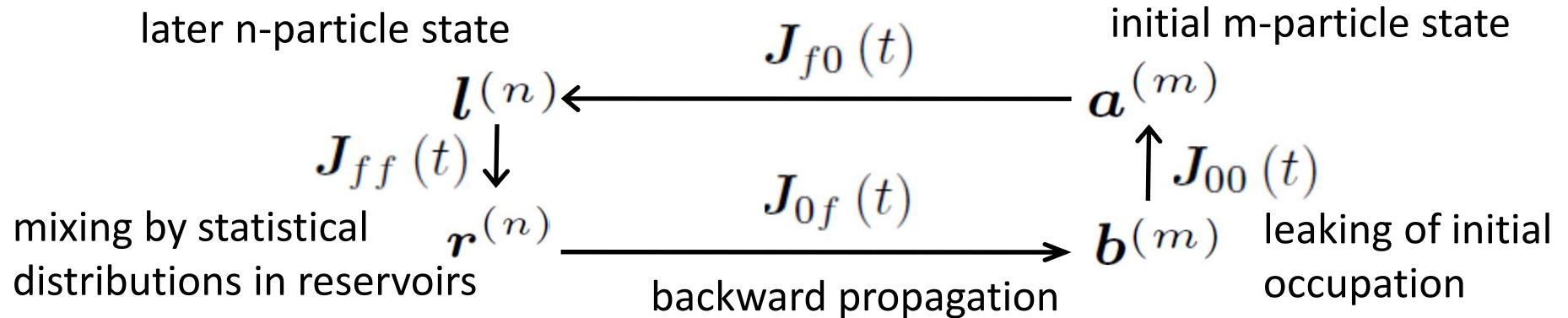
$$\mathcal{J}_{\mathbf{a}^{(m)}, \mathbf{b}^{(m)}}^{\mathbf{l}^{(n)}, \mathbf{r}^{(n)}}(t) = (-1)^m \det \begin{pmatrix} [\mathbf{J}_{00}(t)]_{\mathbf{b}^{(m)}}^{\mathbf{a}^{(m)}} & [\mathbf{J}_{0f}(t)]_{\mathbf{b}^{(m)}}^{\mathbf{r}^{(n)}} \\ [\mathbf{J}_{f0}(t)]_{\mathbf{l}^{(n)}}^{\mathbf{a}^{(m)}} & [\mathbf{J}_{ff}(t)]_{\mathbf{l}^{(n)}}^{\mathbf{r}^{(n)}} \end{pmatrix}$$

$$\begin{aligned} \left([\mathbf{J}_{00}(t)]_{\mathbf{b}^{(m)}}^{\mathbf{a}^{(m)}} \right)_{ij} &= [\mathbf{J}_{00}(t)]_{P\mathbf{b}_i^{(m)} P\mathbf{a}_j^{(m)}} \quad \left([\mathbf{J}_{0f}(t)]_{\mathbf{b}^{(m)}}^{\mathbf{r}^{(n)}} \right)_{ij} = [\mathbf{J}_{0f}(t)]_{P\mathbf{b}_i^{(m)} P\mathbf{r}_j^{(n)}} \\ \left([\mathbf{J}_{f0}(t)]_{\mathbf{l}^{(n)}}^{\mathbf{a}^{(m)}} \right)_{ij} &= [\mathbf{J}_{f0}(t)]_{P\mathbf{r}_i^{(n)} P\mathbf{b}_j^{(m)}} \quad \left([\mathbf{J}_{ff}(t)]_{\mathbf{l}^{(n)}}^{\mathbf{r}^{(n)}} \right)_{ij} = [\mathbf{J}_{ff}(t)]_{P\mathbf{l}_i^{(n)} P\mathbf{r}_j^{(n)}} \end{aligned}$$

$$\begin{aligned} \mathbf{J}_{f0}(t) &= [\mathbf{J}_{0f}(t)]^\dagger = (\mathbf{1}_D - \mathbf{v}(t))^{-1} \mathbf{u}(t) & \mathbf{J}_{ff}(t) &= (\mathbf{1}_D - \mathbf{v}(t))^{-1} - \mathbf{1}_D \\ \mathbf{J}_{00}(t) &= \mathbf{u}^\dagger(t) (\mathbf{1}_D - \mathbf{v}(t))^{-1} \mathbf{u}(t) - \mathbf{1}_D & \mathbf{u}(t) &= \mathbf{u}(t, t_0) \end{aligned}$$

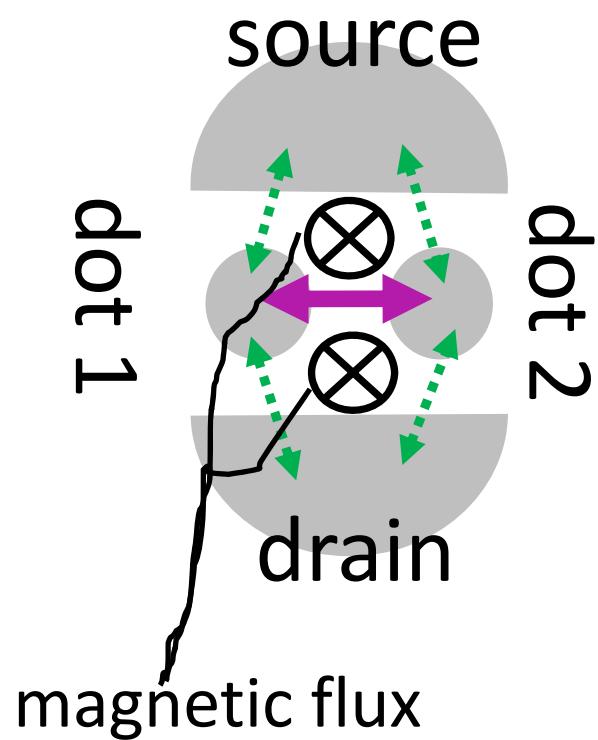
physical interpretation of the solution

forward propagation



III. Example of application

Aharonov Bohm interferometer with double quantum dot molecule



$$H_{qd} = \epsilon_{11} d_1^\dagger d_1 + \epsilon_{22} d_2^\dagger d_2 - t_c(d_1^\dagger d_2 + d_2^\dagger d_1)$$

zero electron state: empty



$$|0, 0\rangle = |0\rangle$$

one electron states

-- local charge states



$$|1, 0\rangle = |1\rangle$$



$$|0, 1\rangle = |2\rangle$$

-- molecular states

$$|b\rangle = \sqrt{1/2}[|1\rangle + |2\rangle]$$

bonding state (BS)

$$|a\rangle = \sqrt{1/2}[|1\rangle - |2\rangle]$$

anti-bonding state (AS)

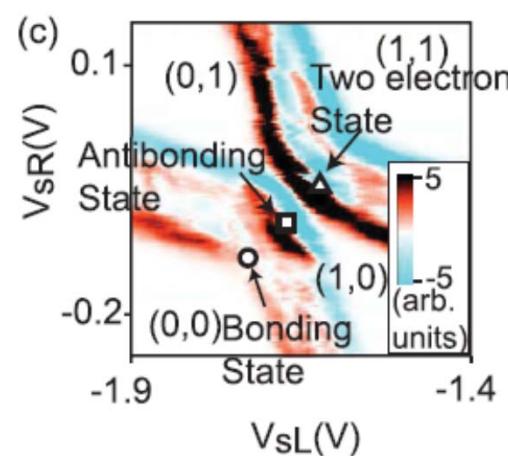
two electron states



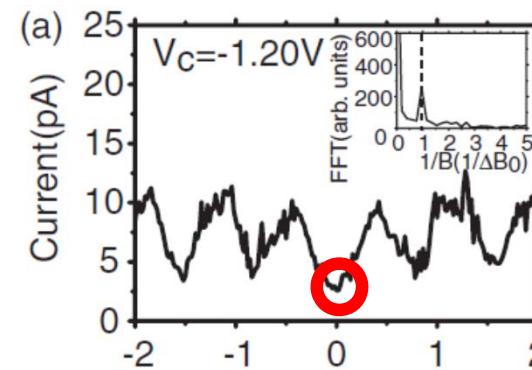
$$|1, 1\rangle = |3\rangle$$

III. Example of application

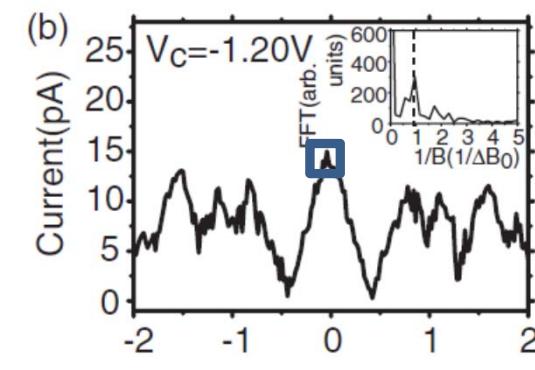
--verifying working regimes



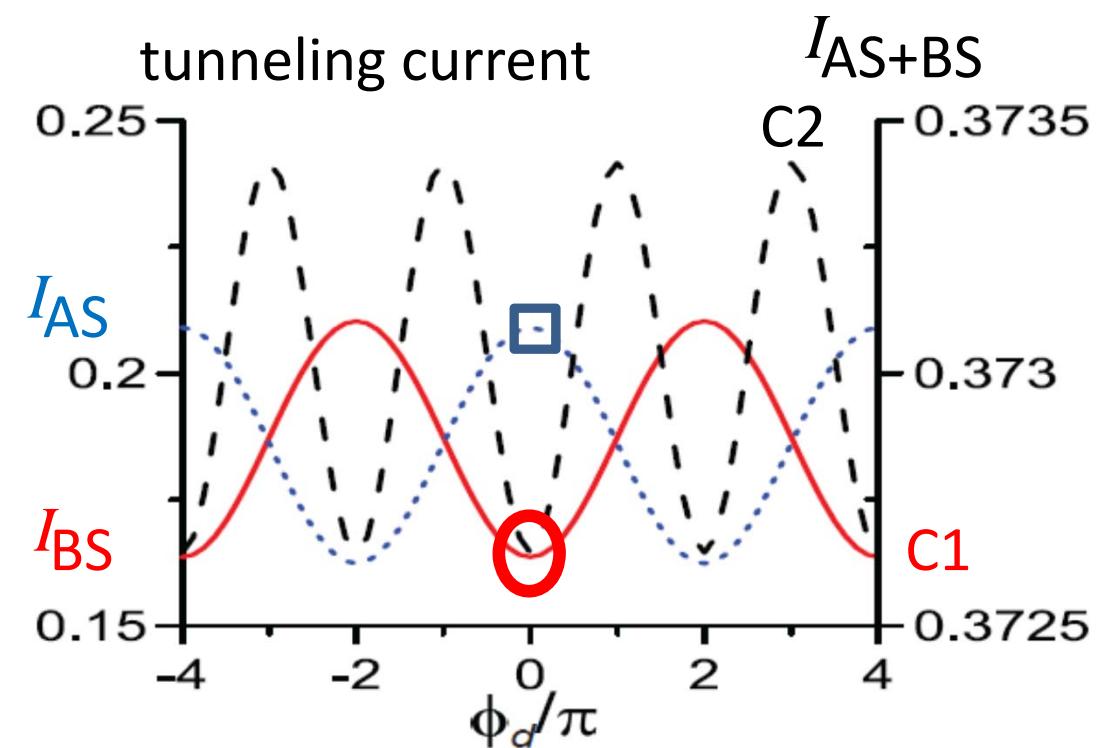
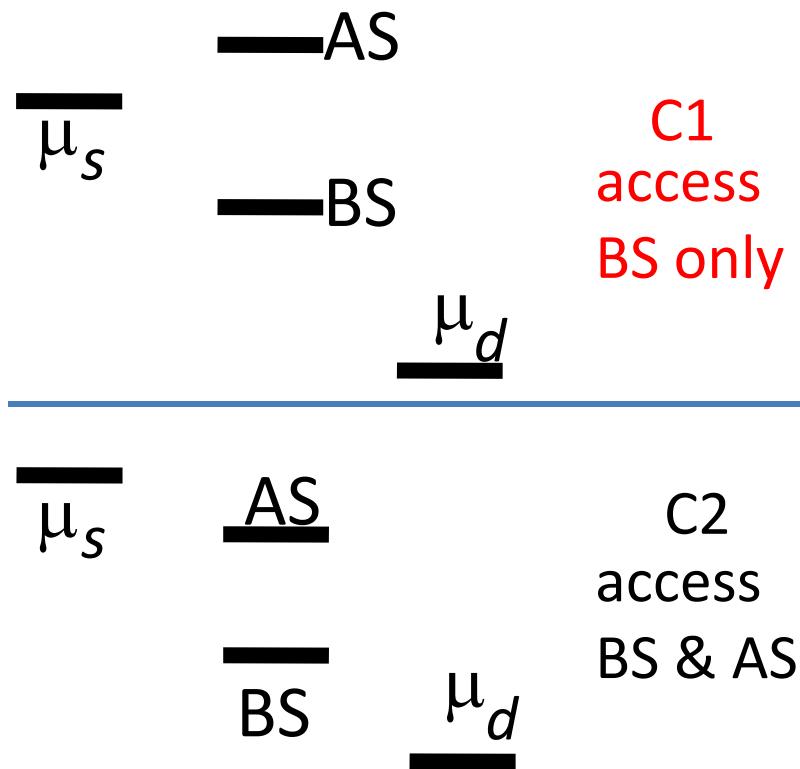
current through BS



current through AS



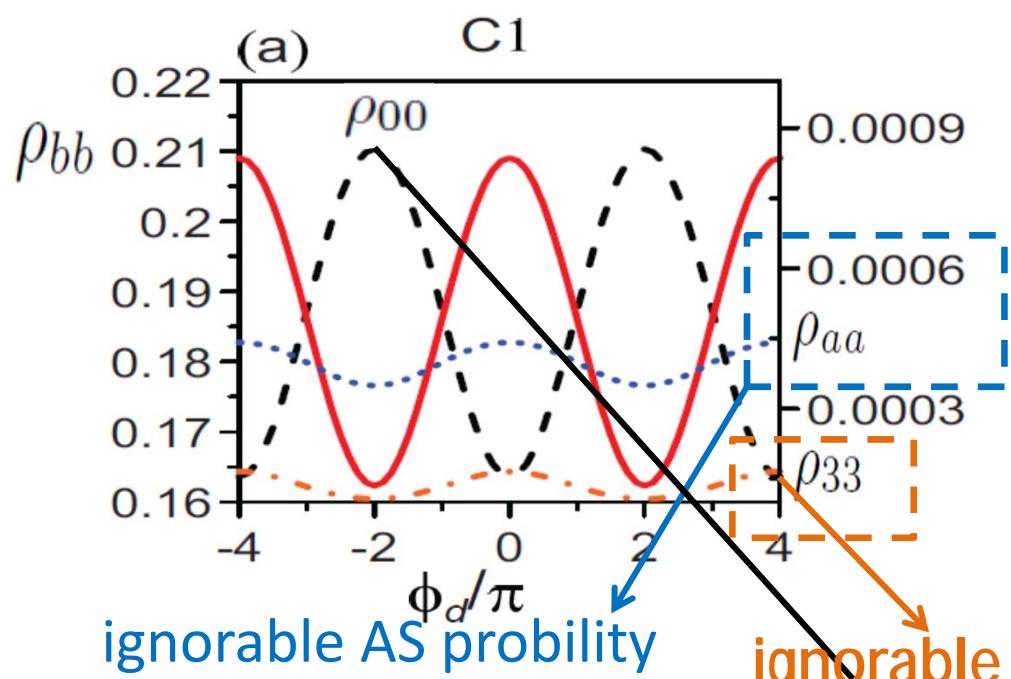
T. Hatano, T. Kubo, Y. Tokura, S. Amaha, S. Teraoka, and S. Tarucha, PRL (2011).



III. Example of application

quantum states under nonequilibrium transports

flux response at steady-state limit

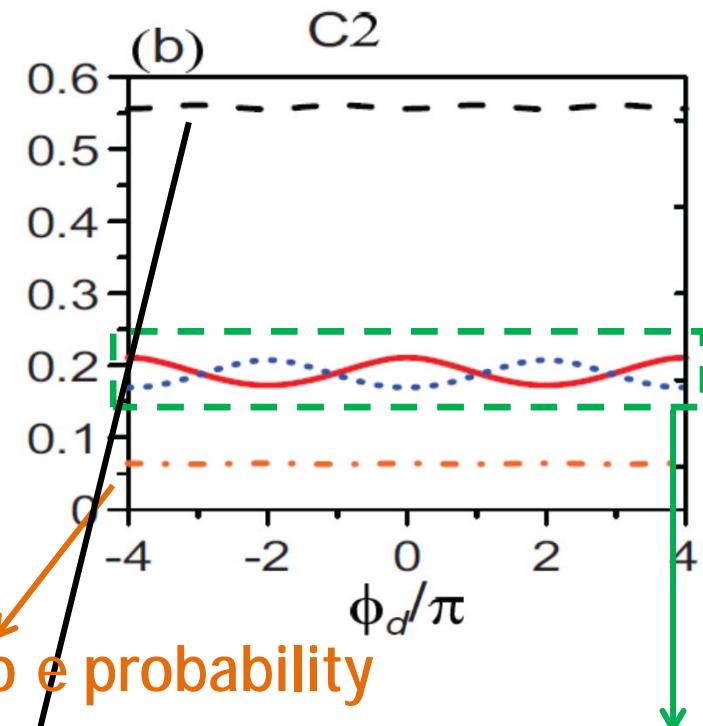


ρ_{00} : empty state probability

ρ_{bb} : one-electron BS probability

ρ_{aa} : one-electron AS probability

ρ_{33} : two-electron state probability



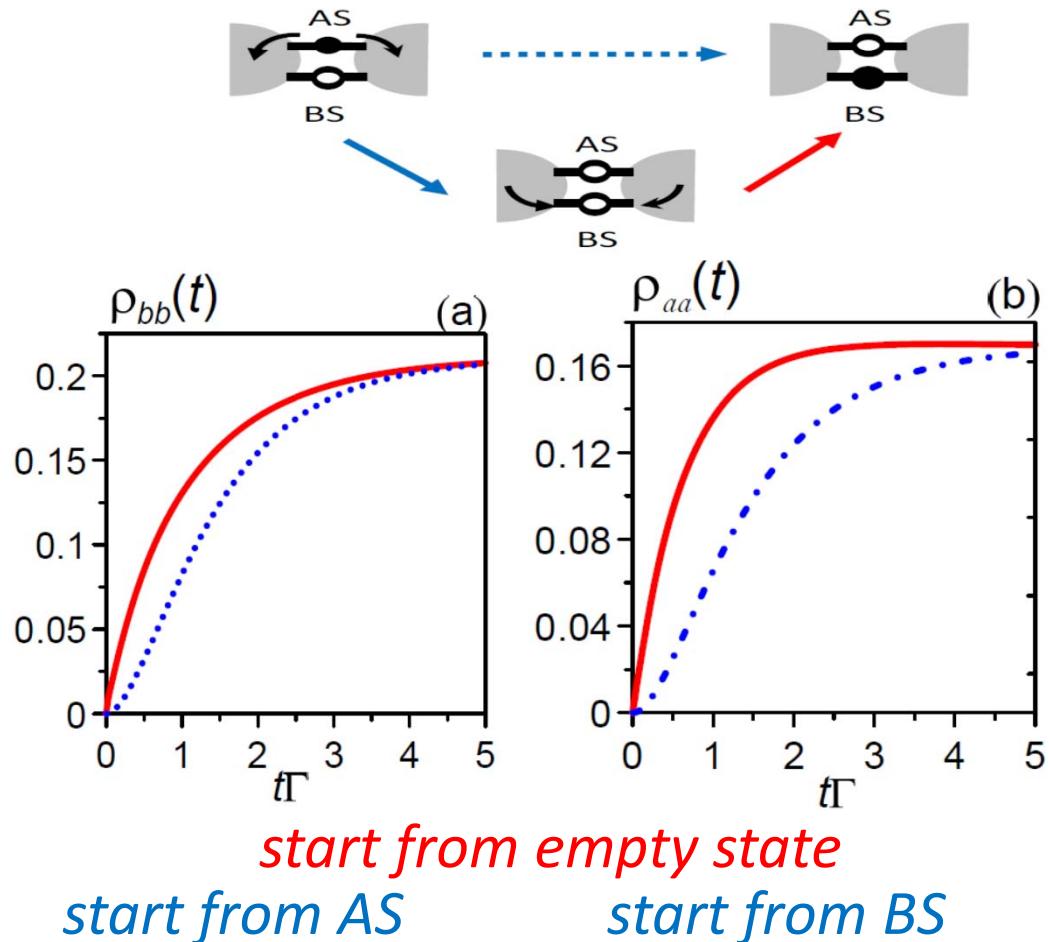
comparable BS and AS probabilities

considerable empty probability

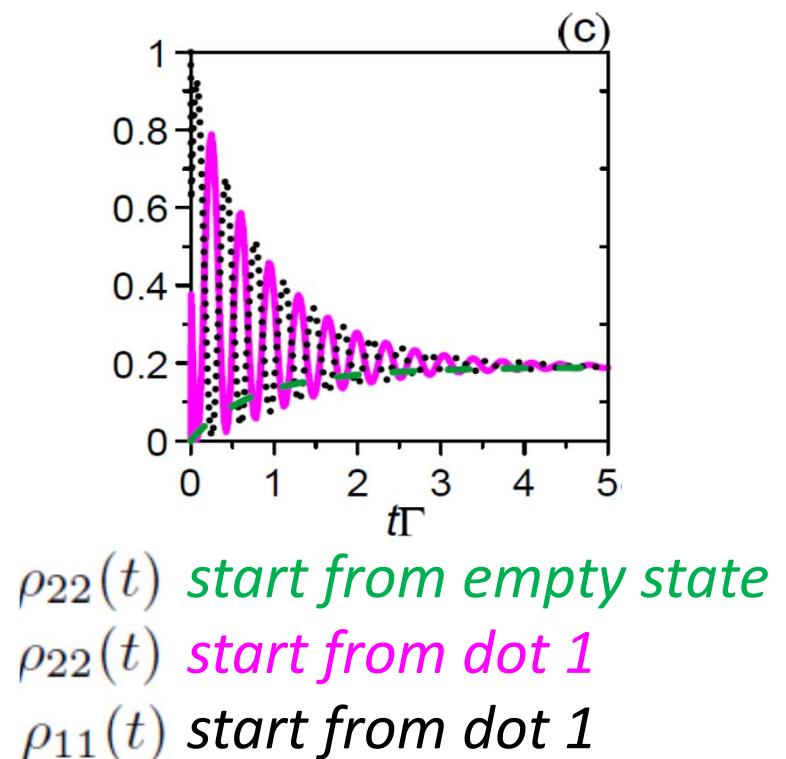
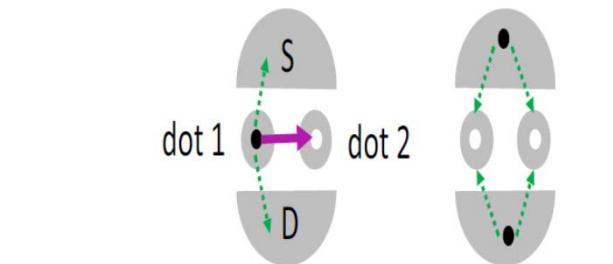
III. Example of application

*quantum states under nonequilibrium transports
real-time quantum state transitions*

external pathways



internal pathway



VI. Summary

Methodological side:

- General analytical solution, applicable to arbitrary number of orbitals and initial preparations for a class of open nanoelectronic systems

Application example :

- Connection between orbital occupation and transport through certain orbitals in DQD AB interferometer
- Distinguishing pathways for quantum state transitions from intermediate evolution processes

Thank you for your attentions