

Introduction to superconducting qubits

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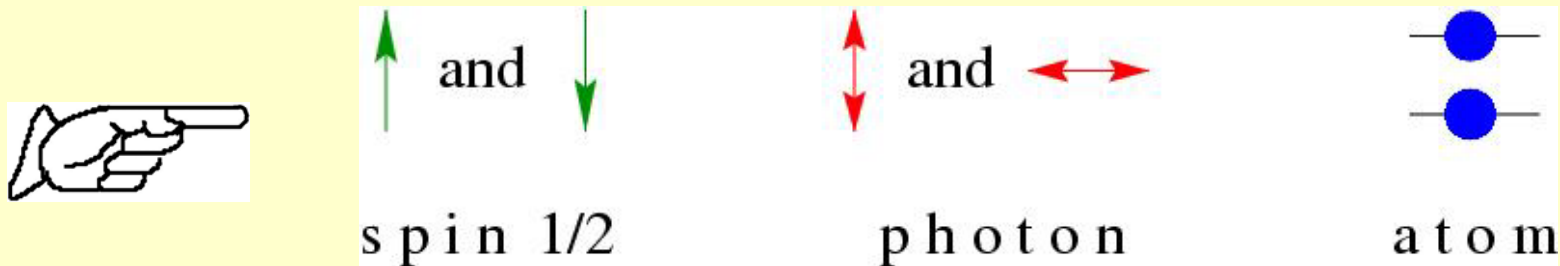
Landau Institute for Theoretical Physics

- **Quantum computation: concept + requirements to realizations**
- **Superconducting qubits: basic principles**
- **Coherence and fluctuations**
- **Read-out: Quantum measurement**
- **Quantum-coherent phenomena in sc-qubits**

Requirements to realizations of quantum computers

(D.DiVincenzo 1997)

- 2-level quantum systems — qubits



superpositions: $a |0\rangle + b |1\rangle$

- w/ well-defined physical parameters (Hamiltonian, ...)
- no leakage to higher states (3-rd, 4-th, ...)
- scalability (N qubits)

Requirements to realizations of quantum computers

$$\mathcal{H} = - \sum_{i=1}^N \mathbf{H}_i(t) \hat{\sigma}_i(t)$$

↑
controlled (on / off)

- 2-level quantum systems — N qubits
- controlled dynamics



- 1-qubit (logic) gates

- spin rotation: $\mathcal{H} = -H_x^i \sigma_x^i$ for time τ

$$U_x^i = \exp(-iH_x^i \hat{\sigma}_x^i \tau / \hbar) = \begin{pmatrix} \cos \alpha & i \sin \alpha \\ i \sin \alpha & \cos \alpha \end{pmatrix}, \quad \alpha = \frac{H_x^i \tau}{\hbar}$$

- phase shift: $\mathcal{H} = -H_z^i \sigma_z^i$ for time τ

$$U_z^i = \exp(iH_z^i \hat{\sigma}_z^i \tau / \hbar) = \begin{pmatrix} e^{i\beta} & 0 \\ 0 & e^{-i\beta} \end{pmatrix}, \quad \beta = \frac{H_z^i \tau}{\hbar}$$

Requirements to realizations of quantum computers

$$\mathcal{H} = - \sum_{i=1}^N \mathbf{H}_i(t) \hat{\sigma}_i(t) + \sum_{i < j} J^{ij}(t) (\hat{\sigma}_x^i \hat{\sigma}_x^j + \hat{\sigma}_y^i \hat{\sigma}_y^j)$$

↑
controlled

↑
controlled

- 2-level quantum systems — N qubits
- controlled dynamics

• 1-qubit gates

} universal set

universal set

• 2-qubit gates



$$\mathcal{H} = J_{ij} (\hat{\sigma}_x^i \hat{\sigma}_x^j + \hat{\sigma}_y^i \hat{\sigma}_y^j) \quad \text{for time } \tau$$

$$U_{2\text{-bit}}^{ij} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos \gamma & -i \sin \gamma & 0 \\ 0 & -i \sin \gamma & \cos \gamma & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{matrix} \left| \uparrow_i \uparrow_j \right\rangle \\ \left| \uparrow_i \downarrow_j \right\rangle \\ \left| \downarrow_i \uparrow_j \right\rangle \\ \left| \downarrow_i \downarrow_j \right\rangle \end{matrix}, \quad \gamma = \frac{2J^{ij}\tau}{\hbar}$$

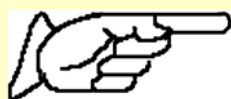
$$U_{2\text{-bit}}(\gamma = \pi/4) : \left| \uparrow \downarrow \right\rangle \rightarrow \frac{1}{\sqrt{2}} (\left| \uparrow \downarrow \right\rangle - i \left| \downarrow \uparrow \right\rangle) \quad - \text{ entangled state}$$

Requirements to realizations of quantum computers

$$\mathcal{H} = - \sum_{i=1}^N \mathbf{H}_i(t) \hat{\sigma}_i(t) + \sum_{i < j} J^{ij}(t) (\hat{\sigma}_x^i \hat{\sigma}_x^j + \hat{\sigma}_y^i \hat{\sigma}_y^j) + \mathcal{H}_{\text{meas}}(t) + \mathcal{H}_{\text{diss}}$$

↑ ↑ ↑ ↑
controlled controlled controlled weak

- 2-level quantum systems — N qubits
- controlled dynamics



- 1-qubit gates

- 2-qubit gates



- initialization - cooling or read-out



- coherence:

$$\tau_{\varphi} > 10^4 \tau_{\text{op}}$$

enough for
error correction

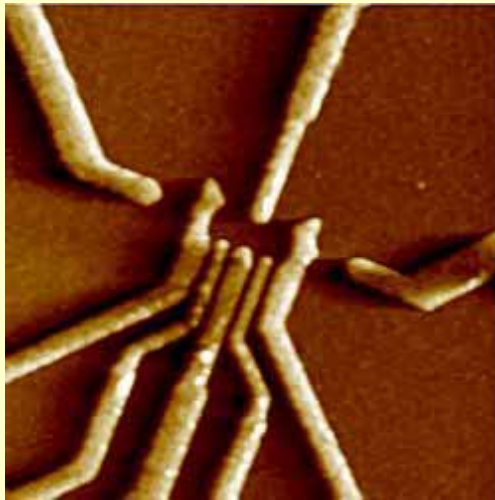
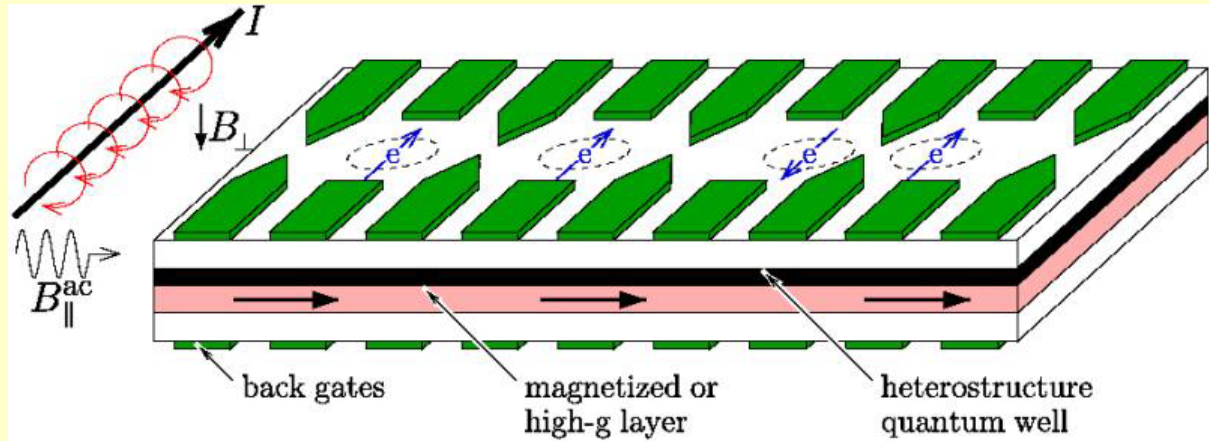


- read-out: quantum measurement

Electron spins in semiconductors Loss & DiVincenzo, 1998

+ τ_{ϕ} spins $>$ τ_{ϕ} charges

+ exactly 2 states



Kouwenhoven et al. (Delft)

Josephson quantum bits

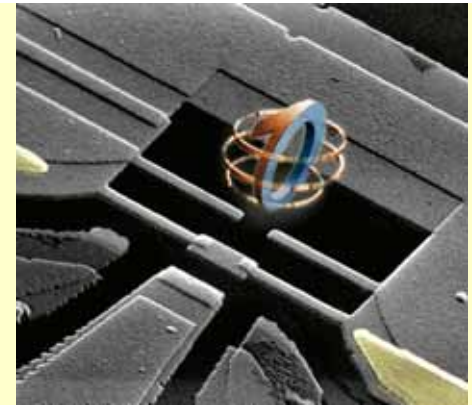
combine:

- coherence of superconducting state
- advanced control techniques for single-charge and SQUID systems

no excitations at low T

quantum degree of freedom: **charge** or **phase** (magn. flux)

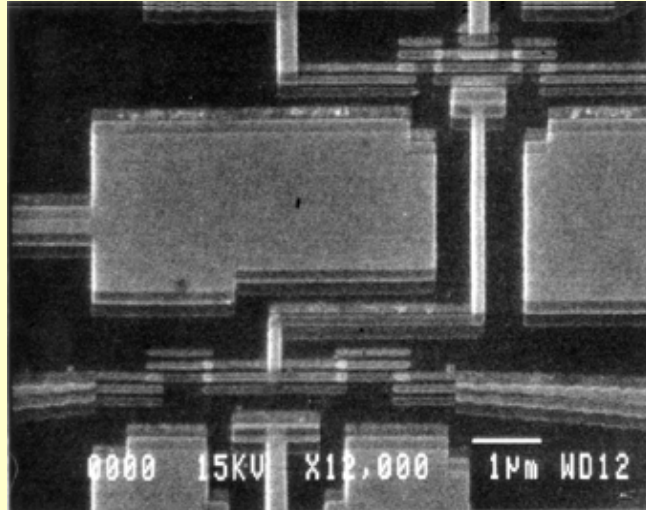
- **macroscopic quantum physics**
(unlike in ion traps, NMR, optical resonators)
- artificial atoms:
 - flexibility in fabrication
- **scalability** (many qubits)
- easy to **integrate** in el. structures



Quantonium

(Saclay)

Single-electron effects

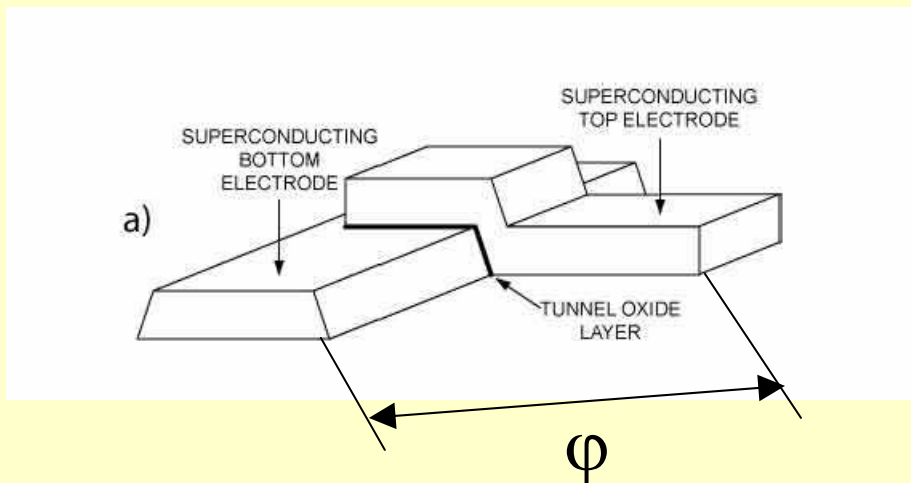


V. Bouchiat et al. (1996)

junctions w/ small area 10nm x 10nm

typical capacitance $C \approx 10^{-15}$ F

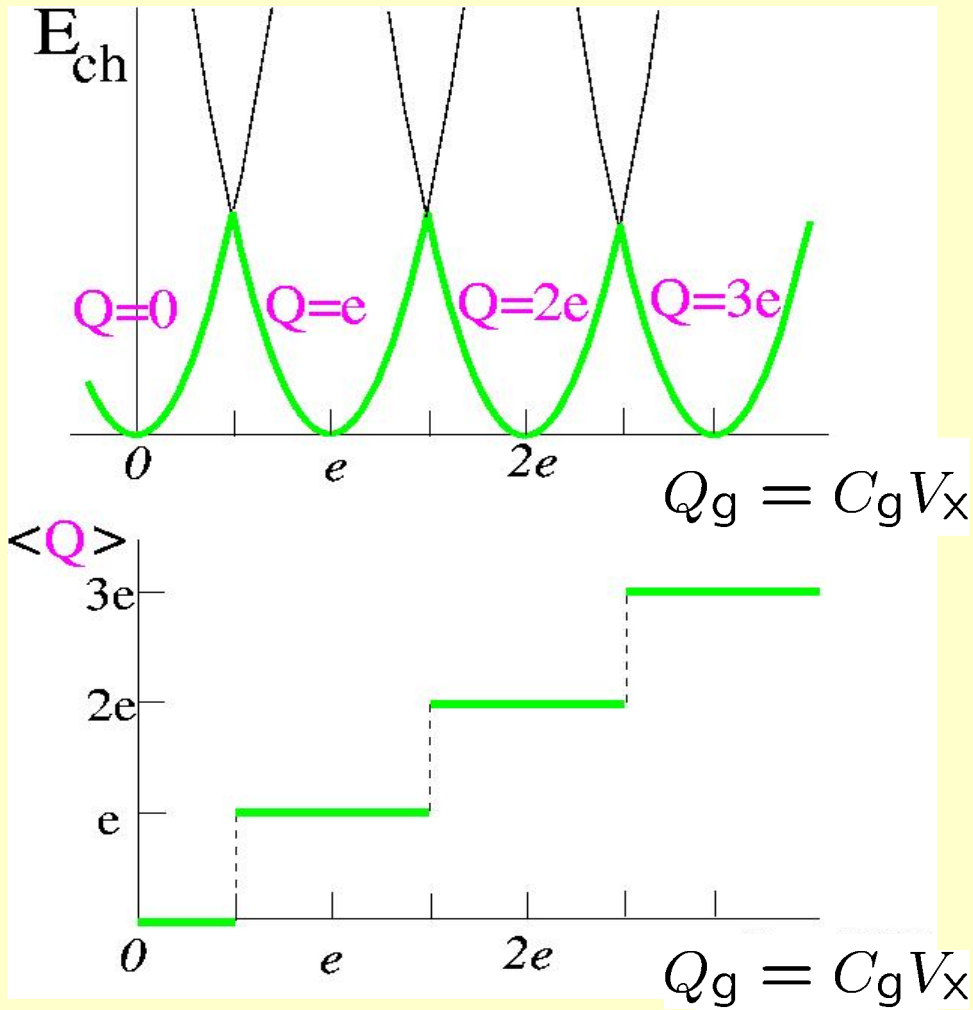
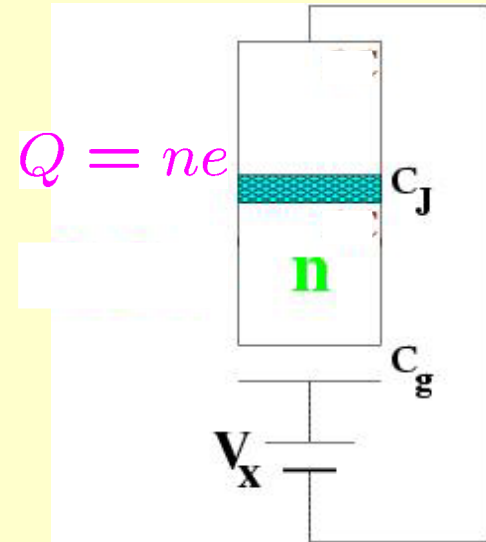
typical energies $E_C = e^2/2C \approx 1$ K



Charging effects in a single-electron box

charging energy:

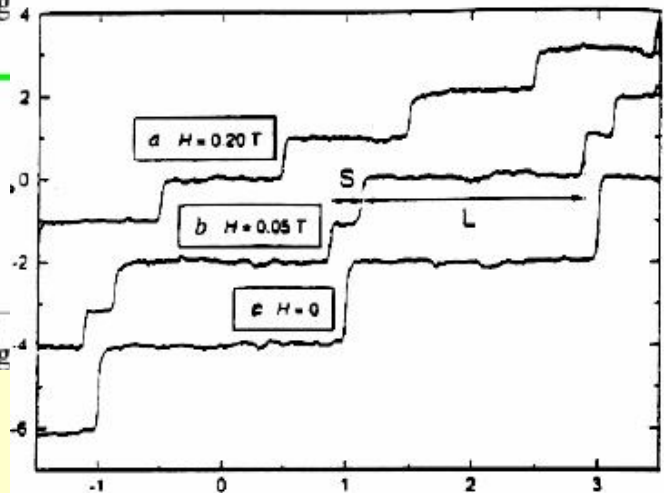
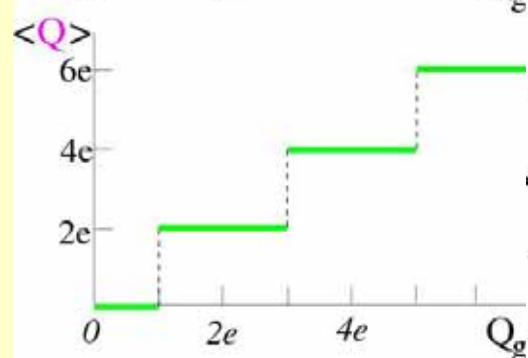
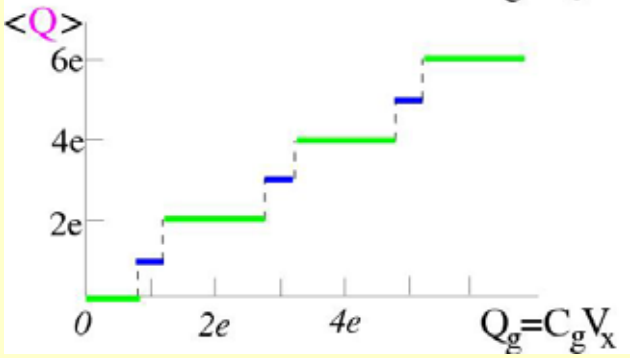
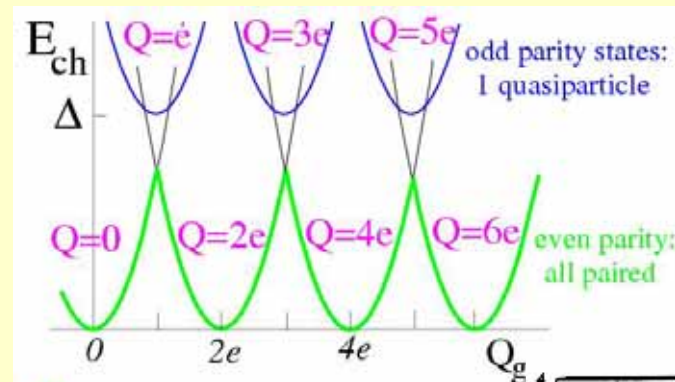
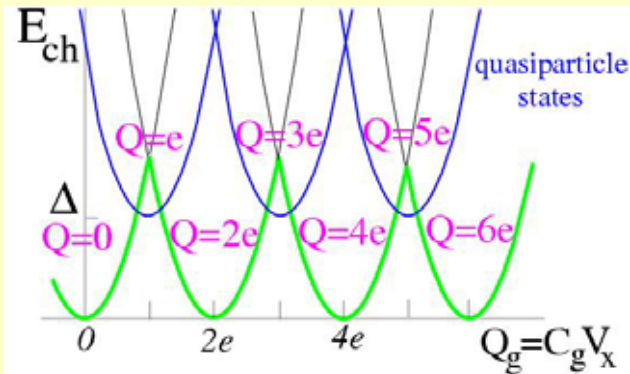
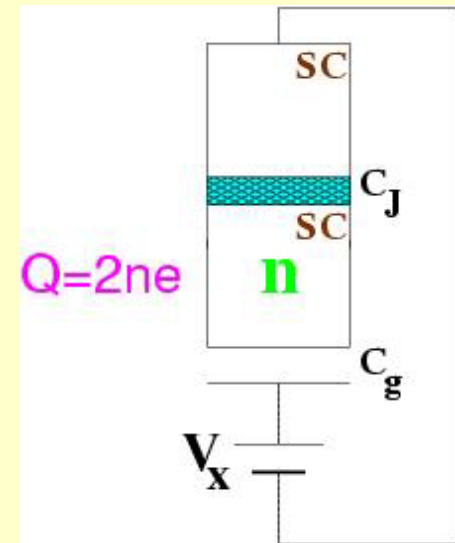
$$E_{\text{ch}}(n, V_x) = \frac{(ne - C_g V_x)^2}{2(C_g + C_J)}$$



Charging effects in a single-Cooper-pair box

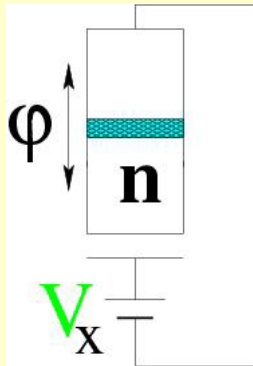
charging energy:

$$E_{\text{ch}}(n, V_x) = \frac{(2ne - C_g V_x)^2}{2(C_g + C_J)}$$

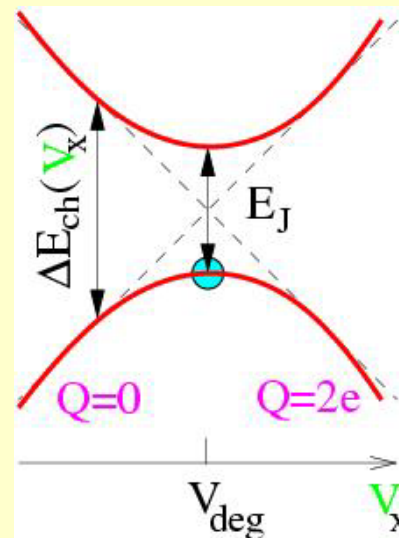
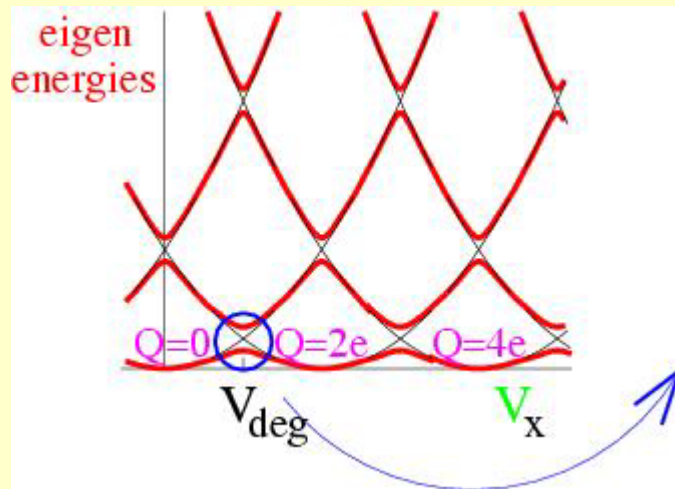


Coherent charge dynamics

$$E_C = e^2 / (2C_\Sigma) \gg E_J$$



charging energy + Josephson energy

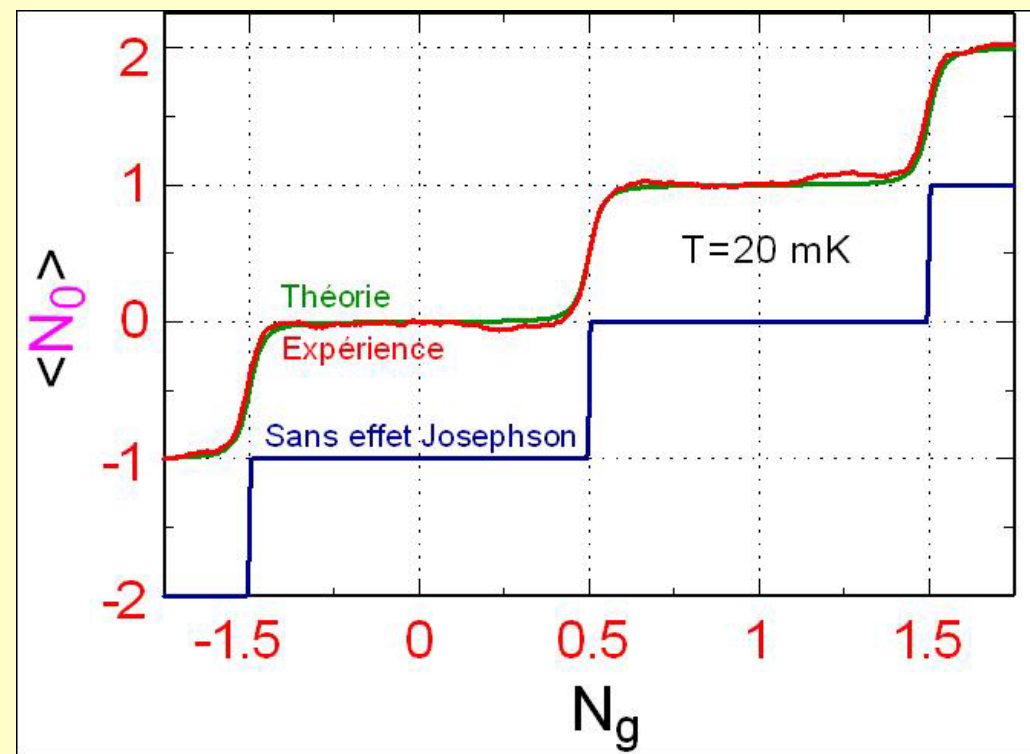
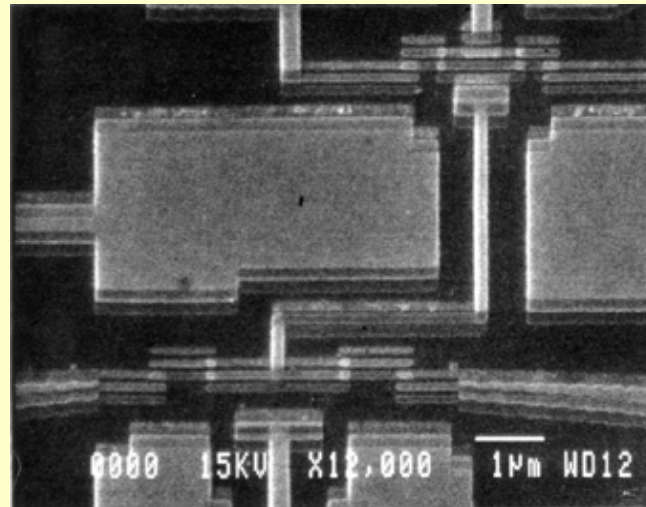


2-level system = qubit

Shnirman et al. 1997
Makhlin et al. 1999

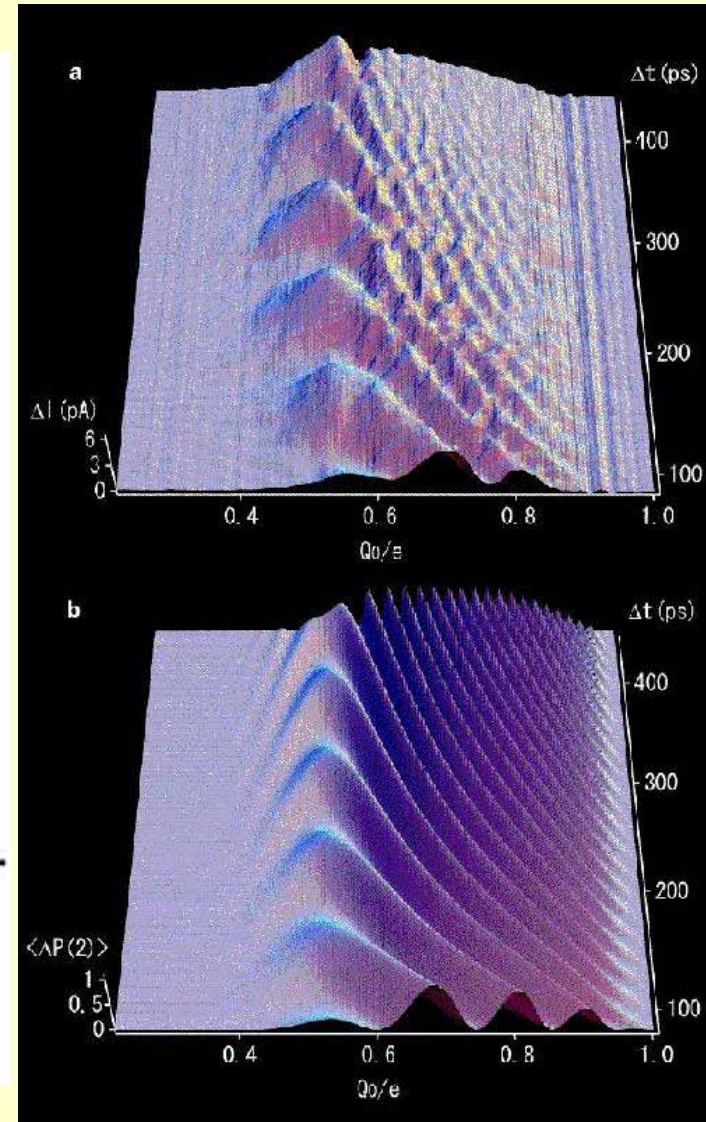
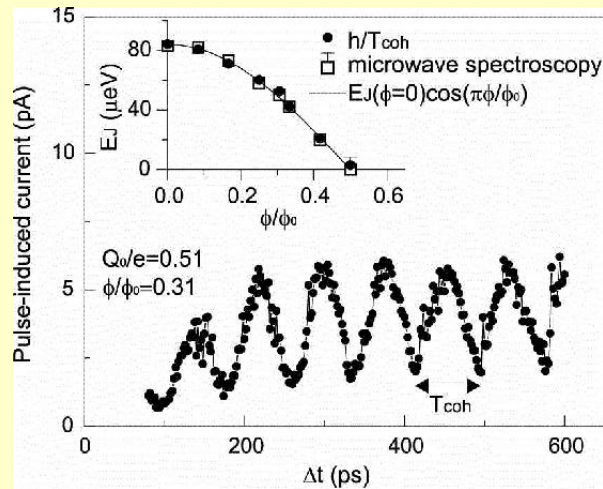
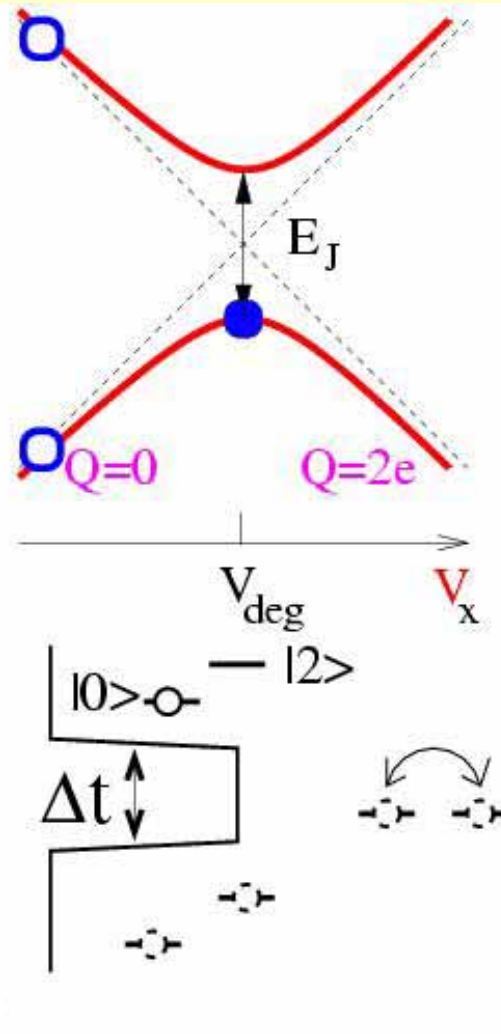
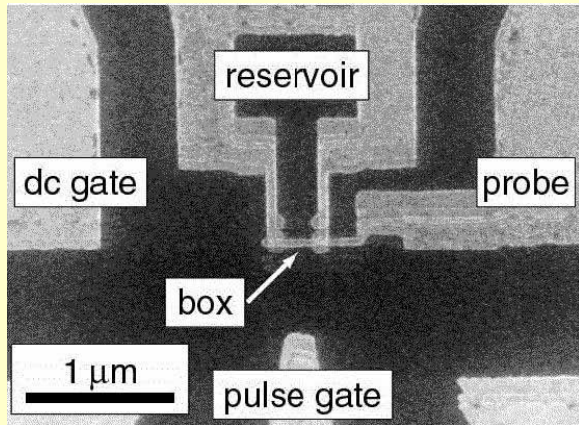
$$\mathcal{H} = -\frac{1}{2} \begin{pmatrix} \Delta E_{ch}(V_x) & E_J \\ E_J & -\Delta E_{ch}(V_x) \end{pmatrix} \begin{matrix} |Q=0\rangle \\ |Q=2e\rangle \end{matrix}$$

Bouchiat et al., Saclay 1997



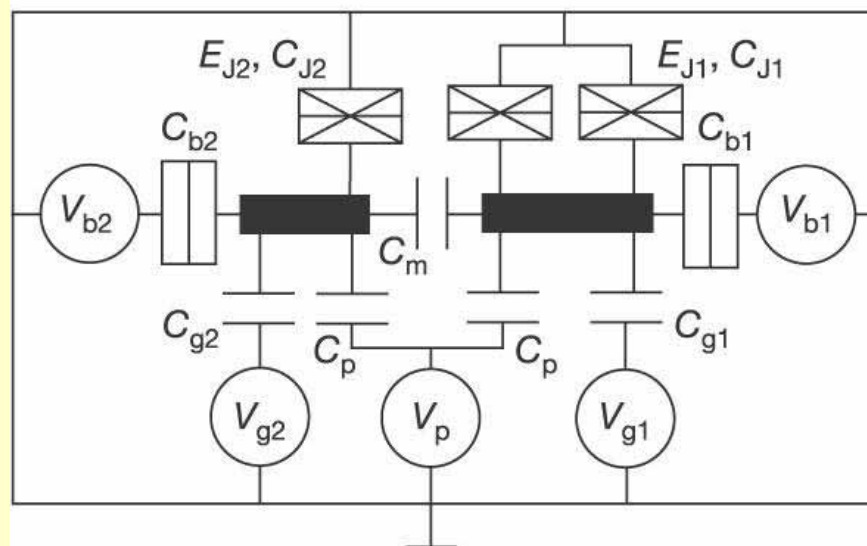
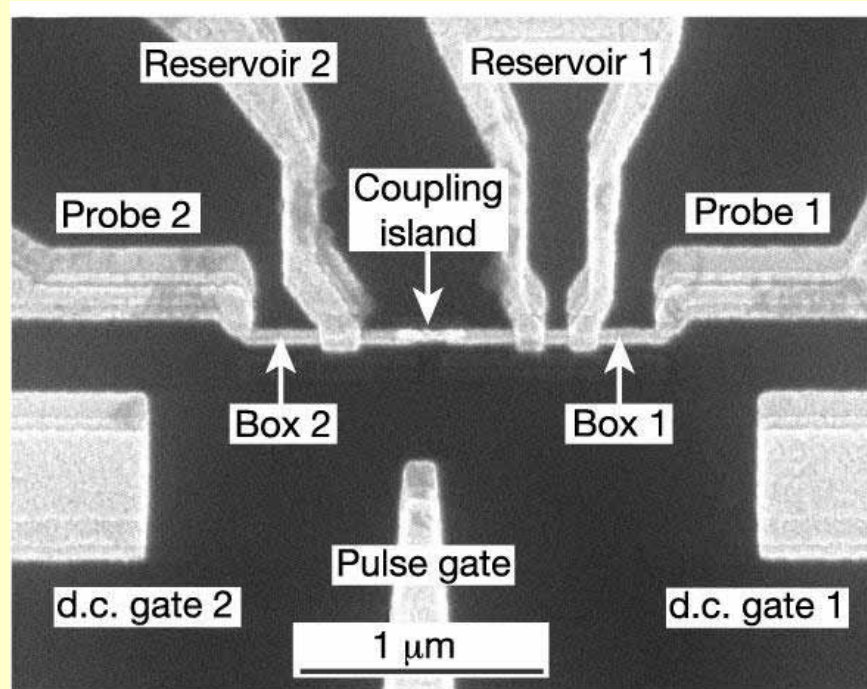
Experiment: coherent oscillations

Y.Nakamura et al., 1999



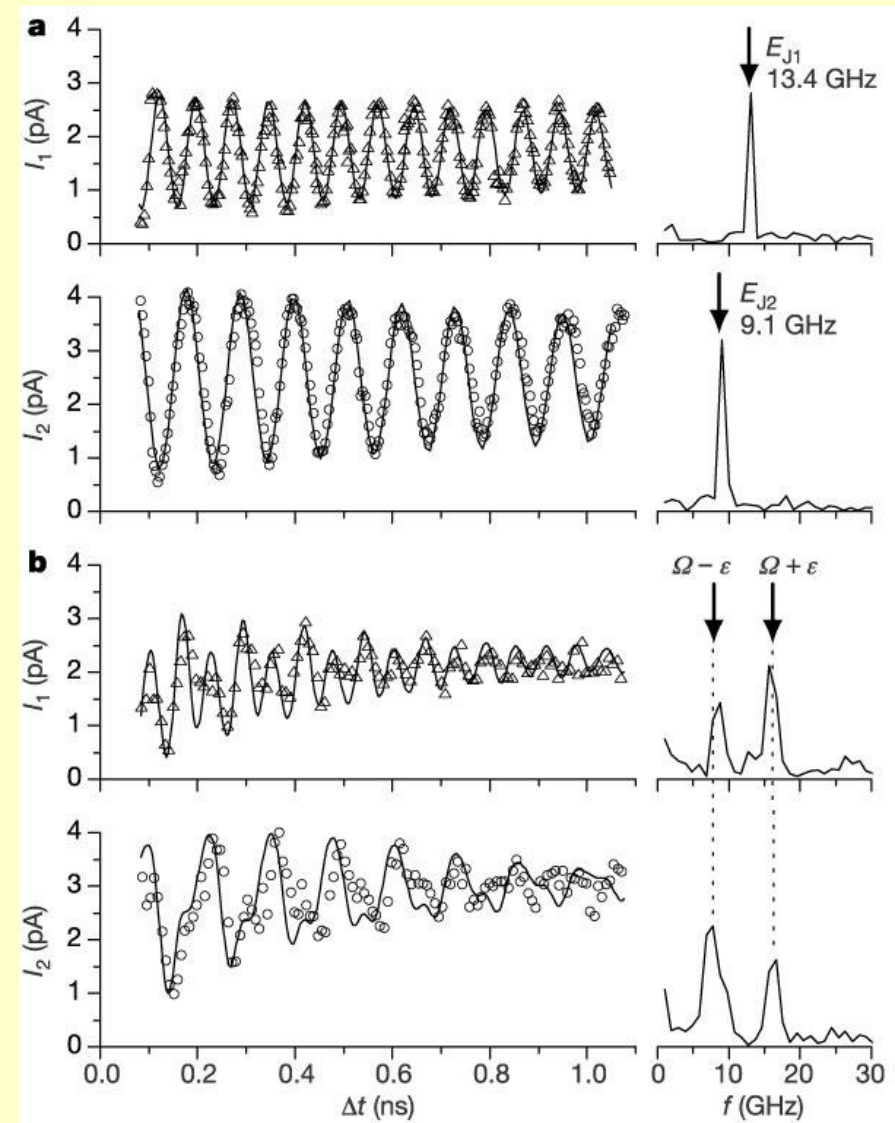
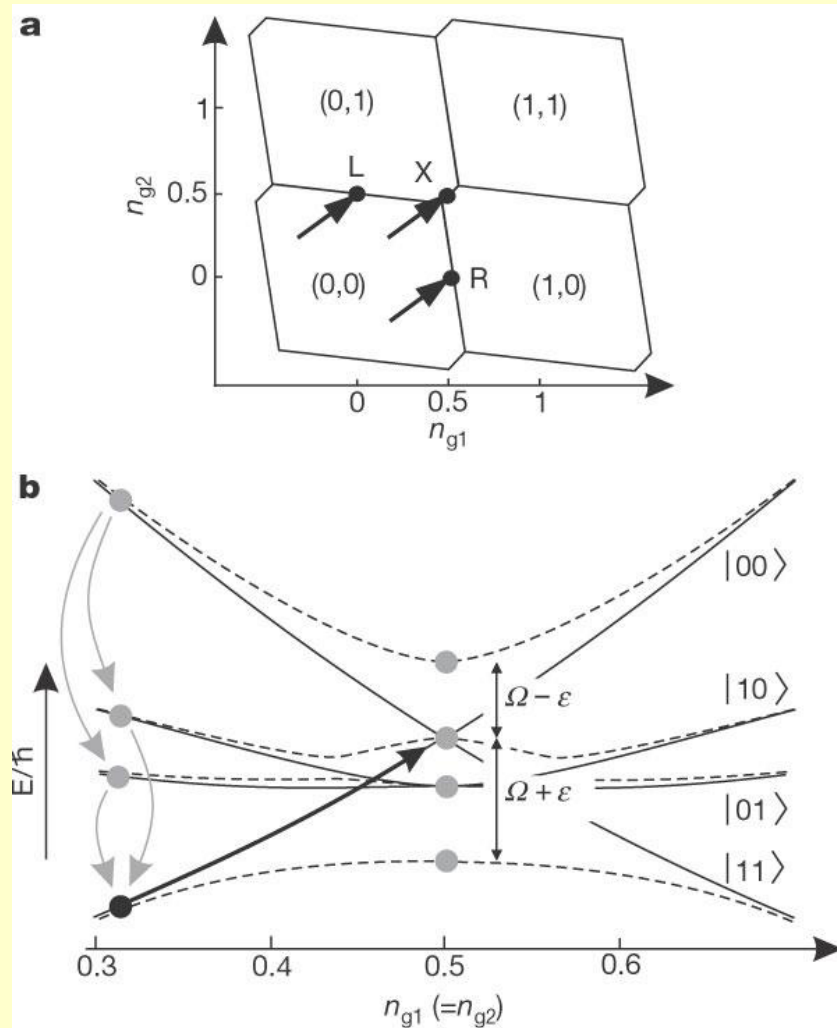
2 qubits: coherent oscillations

Yu.Pashkin et al., 2002

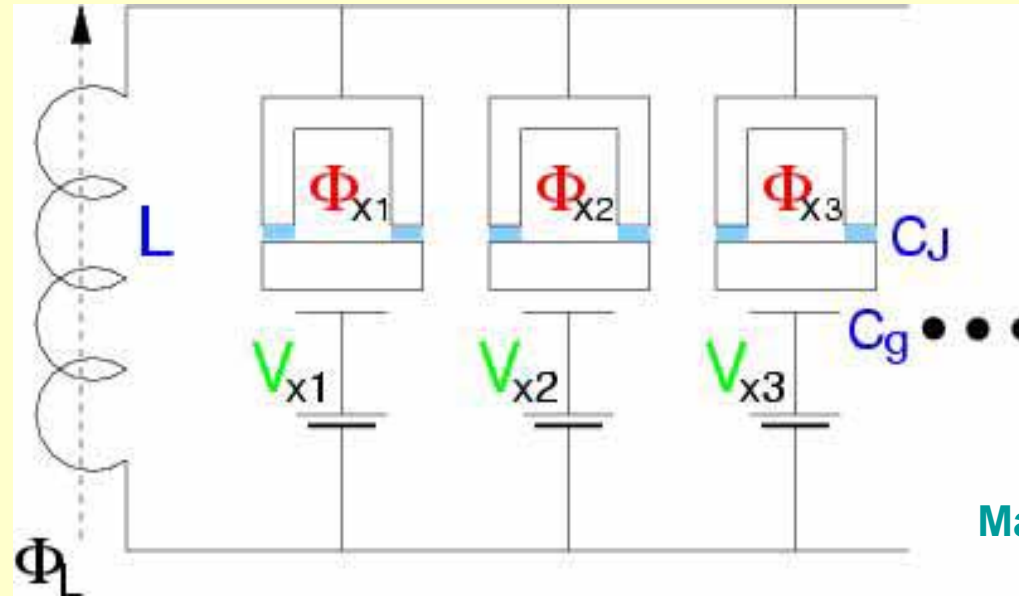


2 qubits: coherent oscillations

Yu.Pashkin et al., 2002



Interaction via an LC-circuit



Makhlin et al. 1999

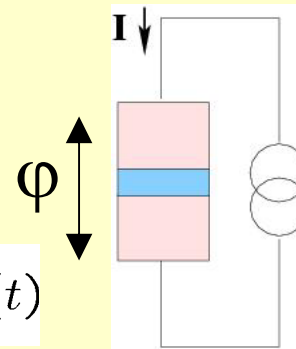
$$\mathcal{H} = -\frac{1}{2} \sum_{i=1}^N \left[\Delta E_{\text{ch}}(V_{xi}) \hat{\sigma}_z^i + E_J(\Phi_{xi}) \hat{\sigma}_x^i \right] + \sum_{i < j} \pi^2 \left(\frac{C_g}{C_J} \right)^2 \frac{E_J(\Phi_{xi}) E_J(\Phi_{xj})}{\Phi_0^2 / L} \hat{\sigma}_y^i \hat{\sigma}_y^j$$

controlled

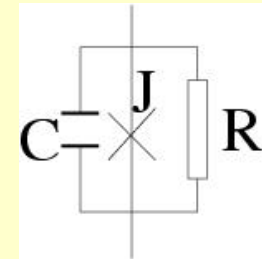
Phase dynamics, SQUIDs

- classical evolution, RCSJ model

$$C \frac{\hbar \ddot{\varphi}}{2e} + \frac{1}{R} \frac{\hbar \dot{\varphi}}{2e} + I_c \sin \varphi = I_{\text{ext}} + \delta I(t)$$



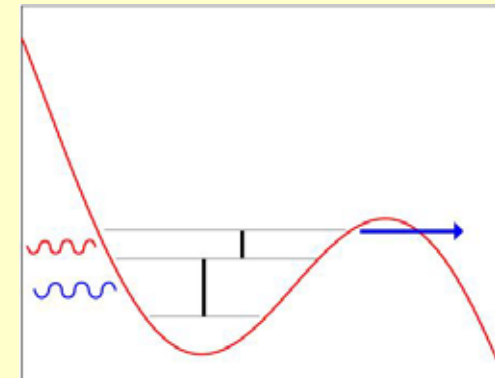
$$E_C = e^2 / (2C) \ll E_J$$



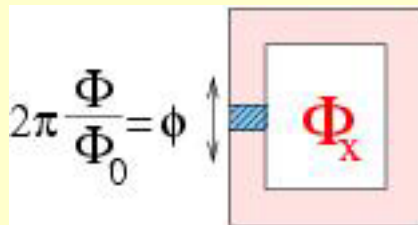
- macroscopic quantum phase evolution

$$\mathcal{H} = \frac{Q^2}{2C} - E_J \cos \varphi - I_{\text{ext}} \varphi + \mathcal{H}_{\text{diss}}; \quad Q = -i\hbar \frac{\partial}{\partial (\hbar \varphi / 2e)}$$

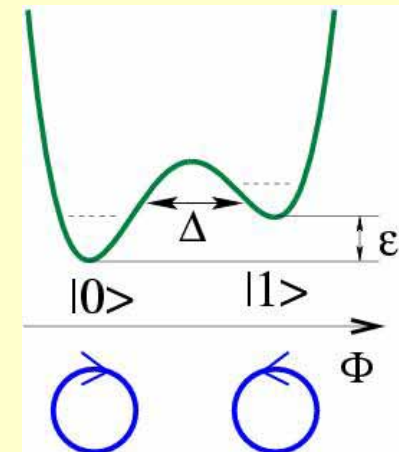
macroscopic & resonant
quantum tunneling



- rf-SQUID

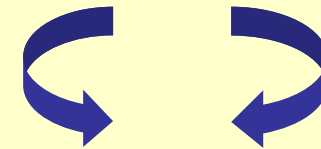
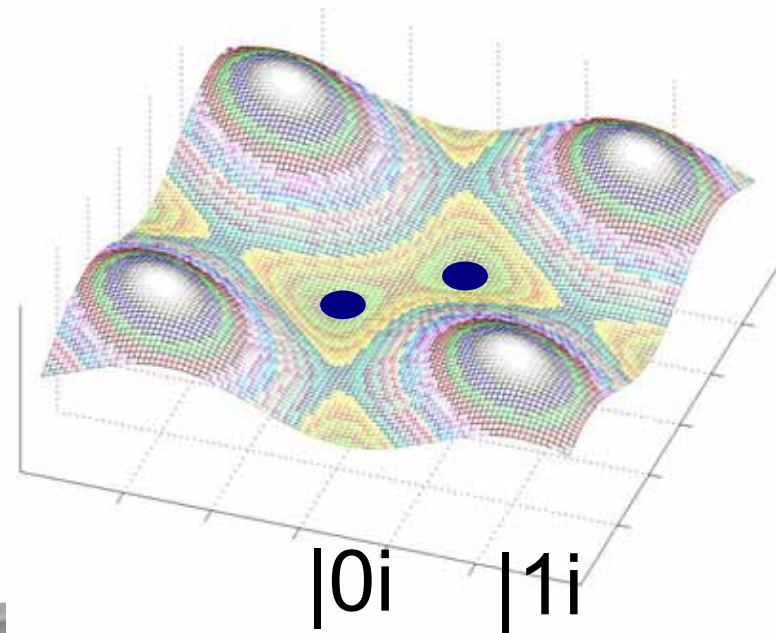
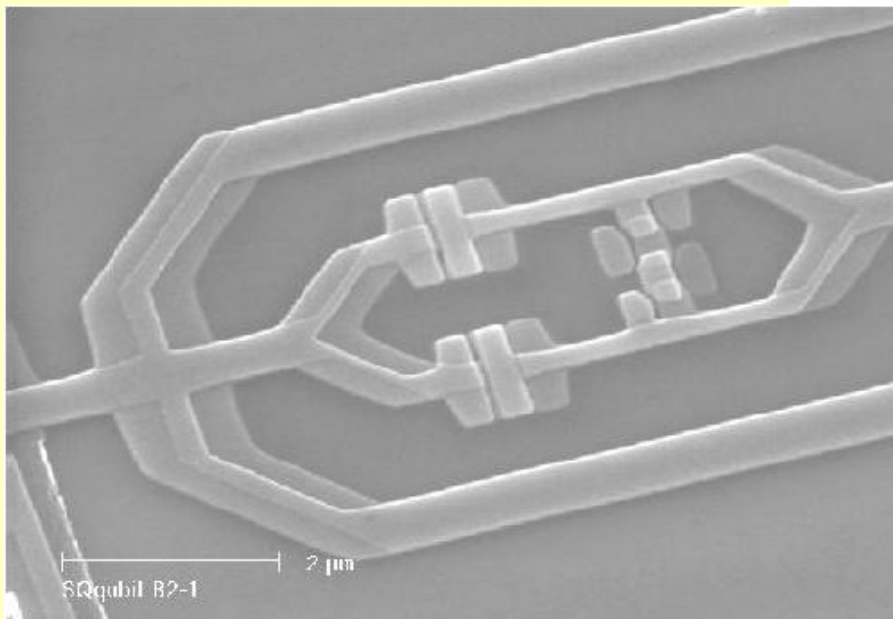
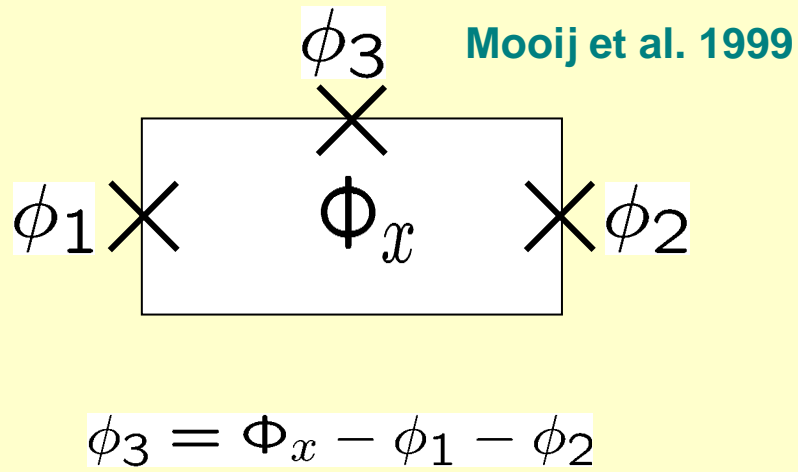


$$\mathcal{H} = \frac{Q^2}{2C} - E_J \cos \left(2\pi \frac{\Phi}{\Phi_0} \right) + \frac{(\Phi - \Phi_x)^2}{2L}$$



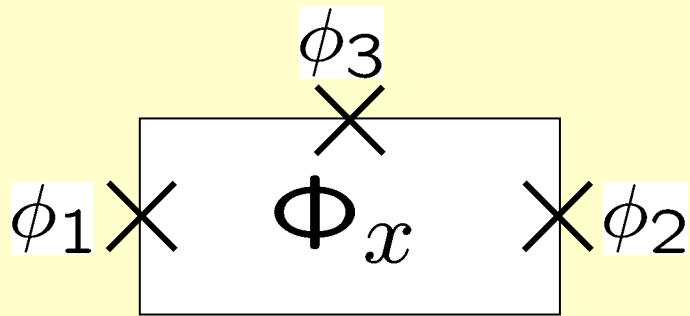
Macroscopic quantum coherence (Leggett)

Superconducting flux qubits



$$H \approx -\frac{1}{2}\epsilon(\Phi_x) \sigma_z - \frac{1}{2}\Delta \sigma_x$$

Chiorescu et al. 2002



$$\phi_3 = \frac{2\pi}{\Phi_0} \Phi_x - \phi_1 - \phi_2$$

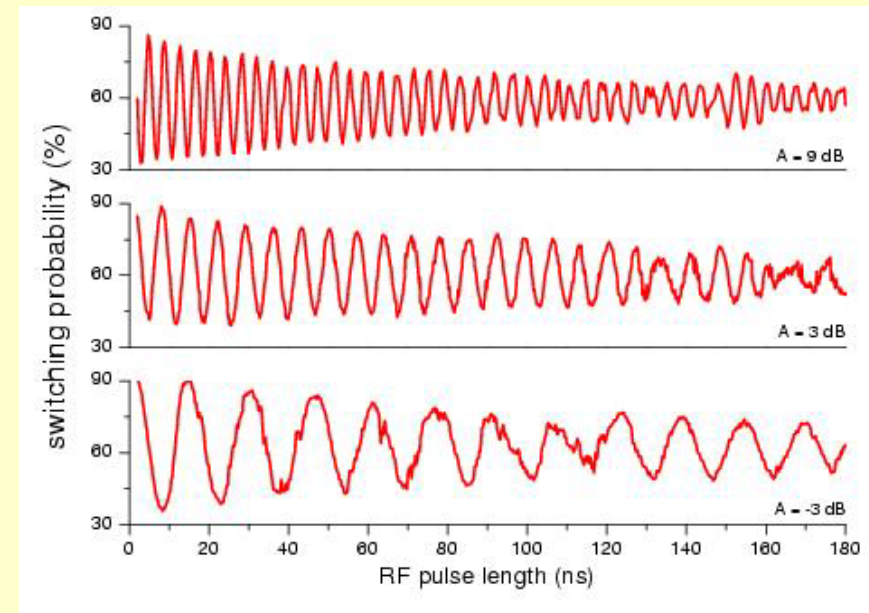
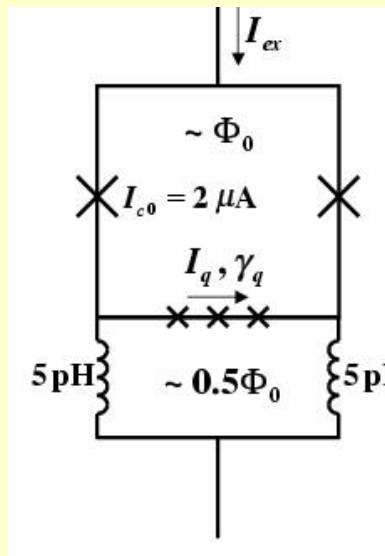
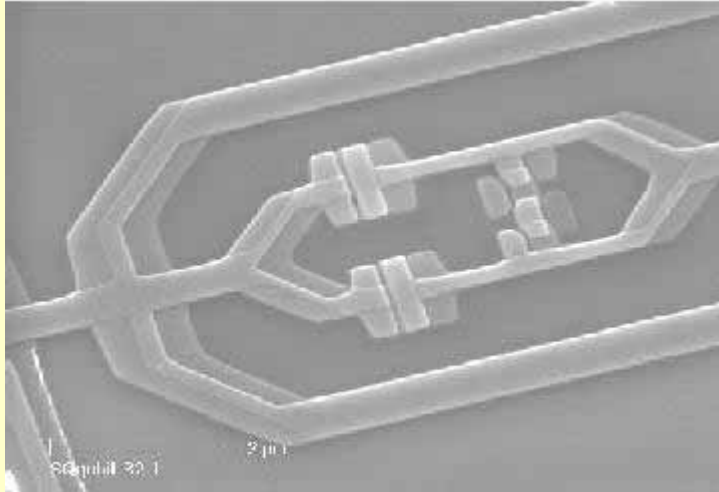
$$U(\phi_1, \phi_2) = -E_J \cos \phi_1 - E_J \cos \phi_2 - \tilde{E}_J \cos\left(\frac{2\pi}{\Phi_0} \Phi_x - \phi_1 - \phi_2\right)$$

for $\phi_1 = \phi_2, \quad \Phi_x = \Phi_0$

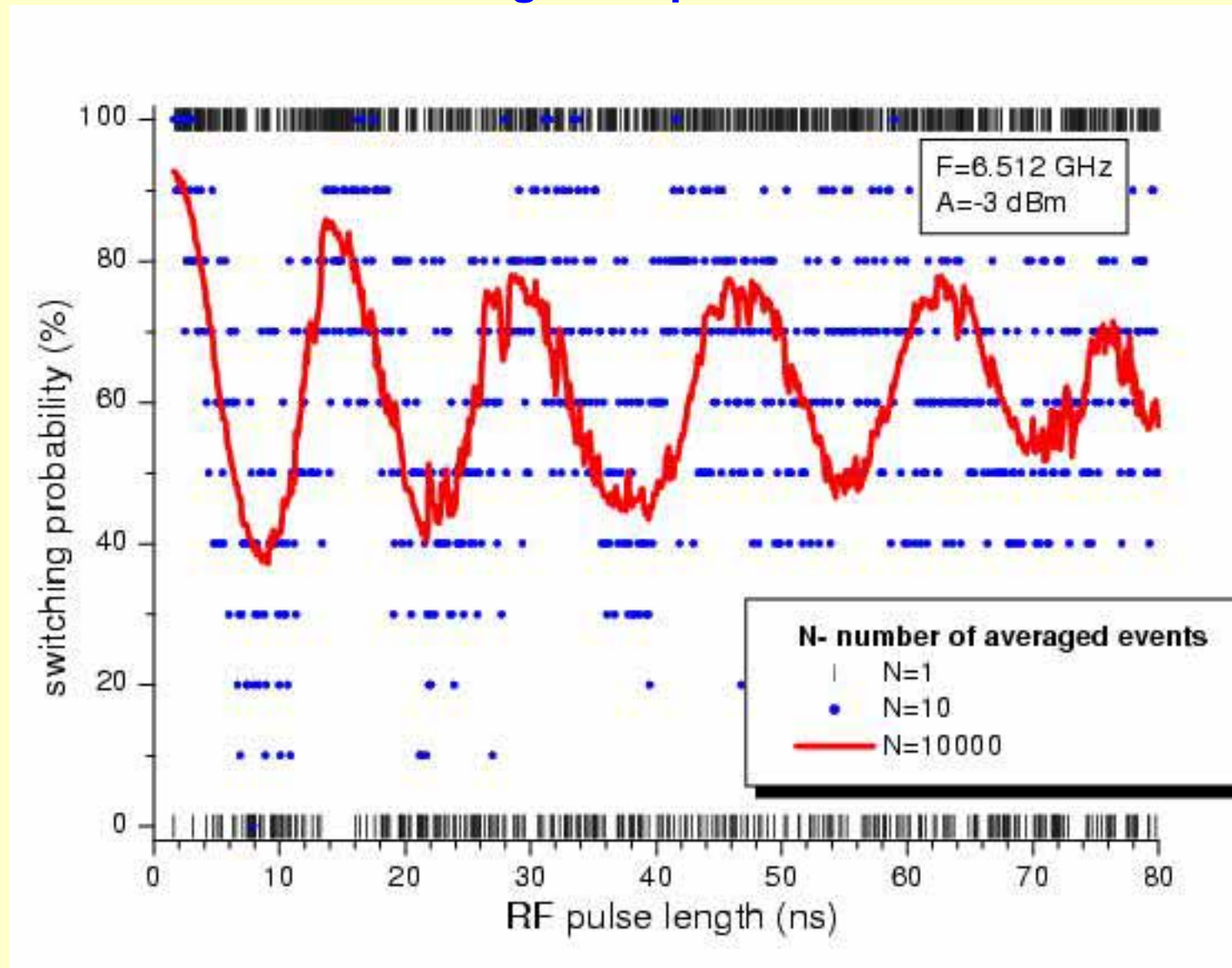
$$U = -2E_J \cos \phi_1 + \tilde{E}_J \cos(2\phi_1)$$

$$\tilde{E}_J \geq 0.5E_J$$

Coherent oscillations: Delft (2002)



Quantum measurement: magnetic qubit + dc-SQUID



strong threshold measurement