

# Direct Detection and Identification of WIMP Dark Matter

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in collaboration with M. Drees and M. Kakizaki

based on arXiv:0903.3300 [hep-ph]

## Direct Dark Matter detection

### Identification of WIMPs

- Reconstruction of the velocity distribution of halo WIMPs

- Determination of the WIMP mass

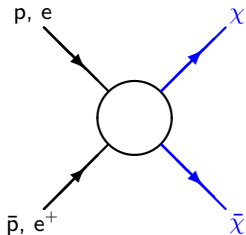
- Determinations of ratios of WIMP-nucleon cross sections

- Estimation of the SI WIMP-nucleon coupling

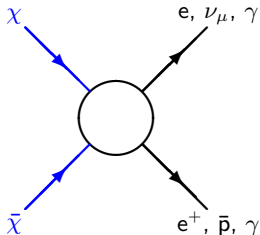
### Summary

## Dark Matter searches

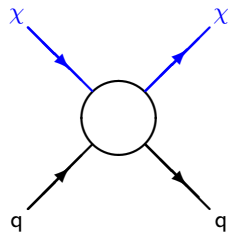
WIMPs should have **small, but non-zero couplings to ordinary matter.**



Colliders



Indirect detection



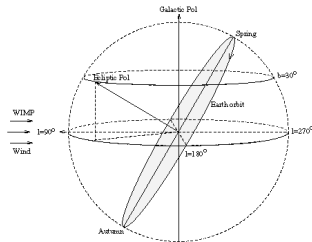
Direct detection

## Direct detection: elastic WIMP-nucleus scattering

- WIMPs could scatter elastically off target nuclei and produce nuclear recoils which deposit energy in the detector.
  - ◆ The event rate depends on the **WIMP density near the Earth**, the **WIMP-nucleus cross section**, the **WIMP mass** and the **velocity distribution of the incident WIMPs**.
  - ◆ In typical SUSY models with neutralino WIMPs, the WIMP-nucleus cross section is about  $10^{-1} \sim 10^{-6}$  pb, the optimistic expected event rate is then  $\sim 10^{-3}$  events/kg-day, but could be  $< 1$  event/ton-yr.
  - ◆ The recoil energy spectrum is **approximately exponential** and most events should be with energies **less than 50 keV**.
  - ◆ Typical background events due to cosmic rays and ambient radioactivity is much larger.

## Direct detection: elastic WIMP-nucleus scattering

- Annual modulation of the WIMP event rate



[Y. Ramachers, Nucl. Phys. Proc. Suppl. 118, 341 (2003)]

- ◆ Due to the orbital motion of the Earth around the Sun.
- ◆ A **cosinusoidal function** with a **one-year period**, a peak around **June 2nd**, and a **modulation amplitude  $\sim 5\%$** .

- ◆ The peak **could be shifted by  $\sim 3$  weeks!**

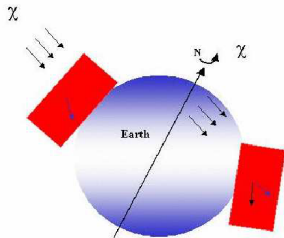
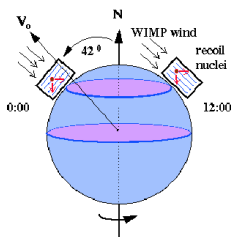
[T. Bruch, J. Read, L. Baudis, and G. Lake, Astrophys. J. 696, 920 (2009)]

- ◆ The **signal identification should also be performed!**

[M. Drees and G. Gerbier, Review of Particle Physics 2008]

## Direct detection: elastic WIMP-nucleus scattering

- Diurnal modulation of the event rate



[Y. Ramachers (2003); M. de Jesus, Int. J. Mod. Phys. A19, 1142 (2004)]

- ◆ Due to the rotation of the Earth.
- ◆ Directionality of the WIMP wind  
A daily forward/backward asymmetry of the nuclear recoil direction.
- ◆ Shielding of the detector by the Earth of the incident WIMP flux.
- ◆ requires a large WIMP-nucleus cross section.

## Direct detection: elastic WIMP-nucleus scattering

- Target material dependence
  - ◆ Spin-independent (SI) coupling  
a scalar (and/or vector) interaction  
the cross section for scalar interaction is approximately proportional to the square of the mass of the nucleus.
  - ◆ Spin-dependent (SD) coupling  
an axial-vector (spin-spin) interaction
  - ◆ For nuclei with  $A \geq 30$ , the SI interaction almost always dominates the SD interaction.
  - ◆ The scattering event rate and the detector sensitivity depend on the mass of the target material directly.

## Basic expression

- Differential event rate for elastic WIMP-nucleus scattering

$$\frac{dR}{dQ} = \mathcal{A} F^2(Q) \int_{v_{\min}}^{v_{\max}} \left[ \frac{f_1(v)}{v} \right] dv$$

Here

$$v_{\min} = \alpha \sqrt{Q}$$

is the minimal incoming velocity of incident WIMPs that can deposit the recoil energy  $Q$  in the detector.

$$\mathcal{A} \equiv \frac{\rho_0 \sigma_0}{2 m_\chi m_{r,N}^2}$$

$$\alpha \equiv \sqrt{\frac{m_N}{2 m_{r,N}^2}}$$

$$m_{r,N} = \frac{m_\chi m_N}{m_\chi + m_N}$$

Particle Physics

$\rho_0$ : WIMP density near the Earth

$\sigma_0$ : total cross section ignoring the form factor suppression

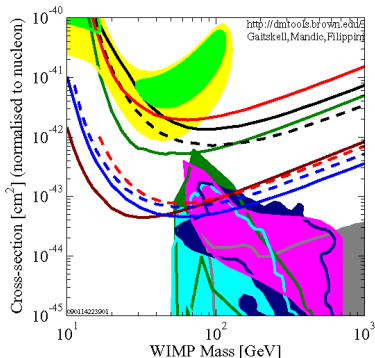
$F(Q)$ : elastic nuclear form factor

$f_1(v)$ : one-dimensional velocity distribution of halo WIMPs



## Basic expression

- Exclusion limits on the (predicted) SI WIMP-nucleon cross section



- DATA listed top to bottom on plot
  - DAMA/LIBRA: 2008 3sigma, no ion channeling
  - KIMS 2007 - 3409 kg-days CsI
  - Edelweiss I final limit, 62 kg-days Ge 2000+2002+2003 limit
  - DAMA/LIBRA: 2008 5sigma, no ion channeling
  - WARP 2.3L, 96.5 kg-days 40 keV threshold
  - CREST 2007 60 kg-day CaWO<sub>4</sub>
  - ZEPLIN III (Dec 2008) result
  - CDMS 2008 Ge
  - CDMS: 2004+2005 (reanalysis) +2008 Ge
  - XENON10 2007 (Net 136 kg-d)
  - Baer et. al well-tempered neutralinos (mixed higgsino,  $\Omega_{\tilde{\chi}^0_1} h^2 \sim 0.1$ )
  - Trotta et al 2008, CMSSM Bayesian: 95% contour
  - Baltz and Gondolo, 2004, Markov Chain Monte Carlo (1 sigma)
  - Ellis et al 2005 CMSSM ( $\mu = 0$ , pion Sigma = 64 MeV)
  - Mastero, Profumo and Ullio: general Split SUSY
- 090114223903  
090114223904

[<http://dmtools.berkeley.edu/limitplots/>]

## Basic expression

- Differential event rate for elastic WIMP-nucleus scattering

$$\left( \frac{dR}{dQ} \right) = \mathcal{A} F^2(Q) \int_{v_{\min}}^{v_{\max}} \left[ \frac{f_1(v)}{v} \right] dv$$

Here

$$v_{\min} = \alpha \sqrt{Q}$$

is the minimal incoming velocity of incident WIMPs that can deposit the recoil energy  $Q$  in the detector.

$$\mathcal{A} \equiv \frac{\rho_0 \sigma_0}{2 m_\chi m_{r,N}^2}$$

$$\alpha \equiv \sqrt{\frac{m_N}{2 m_{r,N}^2}}$$

$$m_{r,N} = \frac{m_\chi m_N}{m_\chi + m_N}$$

$\rho_0$ : WIMP density near the Earth

$\sigma_0$ : total cross section ignoring the form factor suppression

$F(Q)$ : elastic nuclear form factor

$f_1(v)$ : one-dimensional velocity distribution of halo WIMPs

## Reconstruction of the velocity distribution of halo WIMPs

- Normalized one-dimensional velocity distribution function

$$f_1(v) = \mathcal{N} \left\{ -2Q \cdot \frac{d}{dQ} \left[ \frac{1}{F^2(Q)} \left( \frac{dR}{dQ} \right) \right] \right\}_{Q=v^2/\alpha^2}$$

$$\mathcal{N} = \frac{2}{\alpha} \left\{ \int_0^\infty \frac{1}{\sqrt{Q}} \left[ \frac{1}{F^2(Q)} \left( \frac{dR}{dQ} \right) \right] dQ \right\}^{-1}$$

- Moments of the velocity distribution function

$$\langle v^n \rangle = \mathcal{N}(Q_{\text{thre}}) \left( \frac{\alpha^{n+1}}{2} \right) \left[ \frac{2Q_{\text{thre}}^{(n+1)/2}}{F^2(Q_{\text{thre}})} \left( \frac{dR}{dQ} \right)_{Q=Q_{\text{thre}}} + (n+1)I_n(Q_{\text{thre}}) \right]$$

$$\mathcal{N}(Q_{\text{thre}}) = \frac{2}{\alpha} \left[ \frac{2Q_{\text{thre}}^{1/2}}{F^2(Q_{\text{thre}})} \left( \frac{dR}{dQ} \right)_{Q=Q_{\text{thre}}} + I_0(Q_{\text{thre}}) \right]^{-1}$$

$$I_n(Q_{\text{thre}}) = \int_{Q_{\text{thre}}}^\infty Q^{(n-1)/2} \left[ \frac{1}{F^2(Q)} \left( \frac{dR}{dQ} \right) \right] dQ$$

## Reconstruction of the velocity distribution of halo WIMPs

- Ansatz for the reconstructed recoil spectrum in the  $n$ th  $Q$ -bin

$$\left( \frac{dR}{dQ} \right)_{Q \simeq Q_n} \equiv r_n e^{k_n(Q - Q_{s,n})} \quad r_n \equiv \frac{N_n}{b_n}$$

- Logarithmic slope and shifted point in the  $n$ th  $Q$ -bin

$$\overline{Q - Q_n}|_n \equiv \frac{1}{N_n} \sum_{i=1}^{N_n} (Q_{n,i} - Q_n) = \left( \frac{b_n}{2} \right) \coth \left( \frac{k_n b_n}{2} \right) - \frac{1}{k_n}$$

$$Q_{s,n} = Q_n + \frac{1}{k_n} \ln \left[ \frac{\sinh(k_n b_n / 2)}{k_n b_n / 2} \right]$$

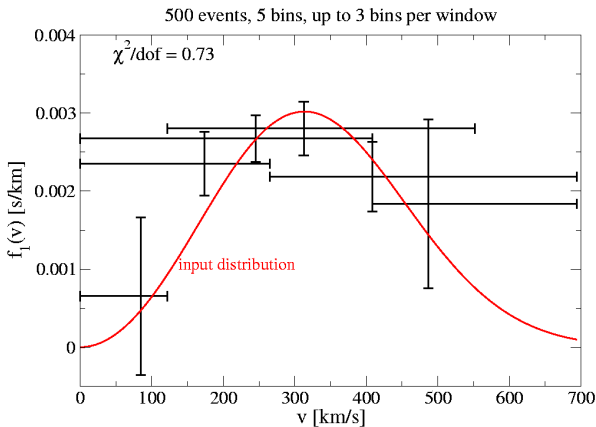
- Reconstructing the one-dimensional velocity distribution

$$f_{1,\text{rec}}(v_{s,n}) = \mathcal{N} \left[ \frac{2Q_{s,n} r_n}{F^2(Q_{s,n})} \right] \left[ \left. \frac{d}{dQ} \ln F^2(Q) \right|_{Q=Q_{s,n}} - k_n \right]$$

$$\mathcal{N} = \frac{2}{\alpha} \left[ \sum_a \frac{1}{\sqrt{Q_a} F^2(Q_a)} \right]^{-1} \quad v_{s,n} = \alpha \sqrt{Q_{s,n}}$$

## Reconstruction of the velocity distribution of halo WIMPs

- Reconstructed  $f_{1,\text{rec}}(v_{s,n})$   
(500 events, 5 bins, up to 3 bins per window)



[M. Drees and CLS, JCAP 0706, 011]

## Determination of the WIMP mass

- Estimating the moments of the WIMP velocity distribution

$$\langle v^n \rangle = \alpha^n \left[ \frac{2Q_{\min}^{1/2} r_{\min}}{F^2(Q_{\min})} + I_0 \right]^{-1} \left[ \frac{2Q_{\min}^{(n+1)/2} r_{\min}}{F^2(Q_{\min})} + (n+1)I_n \right]$$

$$I_n = \sum_a \frac{Q_a^{(n-1)/2}}{F^2(Q_a)} \quad r_{\min} = \left( \frac{dR}{dQ} \right)_{\text{expt}, Q=Q_{\min}}$$

[M. Drees and CLS, JCAP 0706, 011]

- Determining the WIMP mass

$$m_X |_{\langle v^n \rangle} = \frac{\sqrt{m_X m_Y} - m_X \mathcal{R}_n}{\mathcal{R}_n - \sqrt{m_X / m_Y}}$$

$$\mathcal{R}_n = \left[ \frac{2Q_{\min,X}^{(n+1)/2} r_{\min,X} / F_X^2(Q_{\min,X}) + (n+1)I_{n,X}}{2Q_{\min,X}^{1/2} r_{\min,X} / F_X^2(Q_{\min,X}) + I_{0,X}} \right]^{1/n} (X \rightarrow Y)^{-1} \quad (n \neq 0)$$

[CLS and M. Drees, arXiv:0710.4296]

## Determination of the WIMP mass

- Spin-independent (SI) WIMP-nucleus cross section

$$\sigma_0^{\text{SI}} = \left(\frac{4}{\pi}\right) m_{r,N}^2 \left[ Z f_p + (A - Z) f_n \right]^2 \simeq \left(\frac{4}{\pi}\right) m_{r,N}^2 A^2 |f_p|^2 = A^2 \left(\frac{m_{r,N}}{m_{r,p}}\right)^2 \sigma_{\chi p}^{\text{SI}}$$

$$\sigma_{\chi p}^{\text{SI}} = \left(\frac{4}{\pi}\right) m_{r,p}^2 |f_p|^2$$

$f_p, f_n$ : effective SI WIMP-proton/neutron coupling

- Determining the WIMP mass

$$m_X|_{\sigma} = \frac{(m_X/m_Y)^{5/2} m_Y - m_X \mathcal{R}_{\sigma}}{\mathcal{R}_{\sigma} - (m_X/m_Y)^{5/2}}$$

$$\mathcal{R}_{\sigma} = \frac{\mathcal{E}_Y}{\mathcal{E}_X} \left[ \frac{2Q_{\min,X}^{1/2} r_{\min,X} / F_X^2(Q_{\min,X}) + I_{0,X}}{2Q_{\min,Y}^{1/2} r_{\min,X} / F_Y^2(Q_{\min,Y}) + I_{0,Y}} \right]$$

## Determination of the WIMP mass

### ○ $\chi^2$ -fit

$$\chi^2(m_X) = \sum_{i,j} (f_{i,X} - f_{i,Y}) C_{ij}^{-1} (f_{j,X} - f_{j,Y})$$

where

$$f_{i,X} = \alpha_X^i \left[ \frac{2Q_{\min,X}^{(i+1)/2} r_{\min,X} / F_X^2(Q_{\min,X}) + (i+1)l_{i,X}}{2Q_{\min,X}^{1/2} r_{\min,X} / F_X^2(Q_{\min,X}) + l_{0,X}} \right] \left( \frac{1}{300 \text{ km/s}} \right)^i$$

$$f_{n_{\max}+1,X} = \mathcal{E}_X \left[ \frac{A_X^2}{2Q_{\min,X}^{1/2} r_{\min,X} / F_X^2(Q_{\min,X}) + l_{0,X}} \right] \left( \frac{\sqrt{m_X}}{m_X + m_X} \right)$$

$$C_{ij} = \text{cov}(f_{i,X}, f_{j,X}) + \text{cov}(f_{i,Y}, f_{j,Y})$$

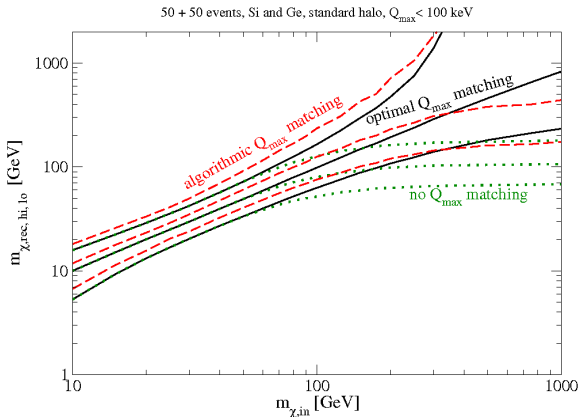
### ○ Algorithmic $Q_{\max}$ matching

$$Q_{\max,Y} = \left( \frac{\alpha_X}{\alpha_Y} \right)^2 Q_{\max,X} \quad \left( v_{\text{cut}} = \alpha \sqrt{Q_{\max}} \right)$$



## Determination of the WIMP mass

- Reconstructed  $m_{\chi, \text{rec}}$   
( $Q_{\text{max}} < 100 \text{ keV}$ ,  $^{28}\text{Si} + ^{76}\text{Ge}$ ,  $2 \times 50 \text{ events}$ )



[M. Drees and CLS, JCAP 0806, 012]

## Determinations of ratios of WIMP-nucleon cross sections

- **-1-st moment** of the WIMP velocity distribution

$$\begin{aligned} \left( \frac{dR}{dQ} \right)_{\text{expt, } Q=Q_{\min}} &= \mathcal{E} A F^2(Q_{\min}) \int_{v(Q_{\min})}^{v(Q_{\max})} \left[ \frac{f_1(v)}{v} \right] dv \\ &= \mathcal{E} \left( \frac{\rho_0 \sigma_0}{2 m_\chi m_{r,N}^2} \right) F^2(Q_{\min}) \cdot \frac{1}{\alpha} \left[ \frac{2 r_{\min} / F^2(Q_{\min})}{2 Q_{\min}^{1/2} r_{\min} / F^2(Q_{\min}) + l_0} \right] \end{aligned}$$

- Product of the local density times the WIMP-nucleus cross section

$$\rho_0 \sigma_0 = \left( \frac{1}{\mathcal{E}} \right) m_\chi m_{r,N} \sqrt{\frac{m_N}{2}} \left[ \frac{2 Q_{\min}^{1/2} r_{\min}}{F^2(Q_{\min})} + l_0 \right]$$

[M. Drees, M. Kakizaki, and CLS, UCLA Dark Matter 2008]

- Ratio of two WIMP-nucleus cross sections

$$\frac{\sigma_{0,X}}{\sigma_{0,Y}} = \left( \frac{\mathcal{E}_Y}{\mathcal{E}_X} \right) \frac{m_{r,X} \sqrt{m_X}}{m_{r,Y} \sqrt{m_Y}} \left[ \frac{2 Q_{\min,X}^{1/2} r_{\min,X} + l_{0,X} F_X^2(Q_{\min,X})}{2 Q_{\min,Y}^{1/2} r_{\min,Y} + l_{0,Y} F_Y^2(Q_{\min,Y})} \right] \left[ \frac{F_Y^2(Q_{\min,Y})}{F_X^2(Q_{\min,X})} \right]$$

## Determination of the ratio of two SD WIMP-nucleon couplings

- Spin-dependent (SD) WIMP-nucleus cross section

$$\sigma_0^{\text{SD}} = \left(\frac{32}{\pi}\right) G_F^2 m_{r,N}^2 \left(\frac{J+1}{J}\right) [a_p \langle S_p \rangle + a_n \langle S_n \rangle]^2$$

$$\sigma_{\text{xp/n}}^{\text{SD}} = \left(\frac{32}{\pi}\right) G_F^2 m_{r,p/n}^2 \cdot \left(\frac{3}{4}\right) a_{p/n}^2$$

$J$ : total nuclear spin

$\langle S_p \rangle, \langle S_n \rangle$ : expectation value of the proton/neutron group spin

$a_p, a_n$ : effective SD WIMP-proton/neutron coupling

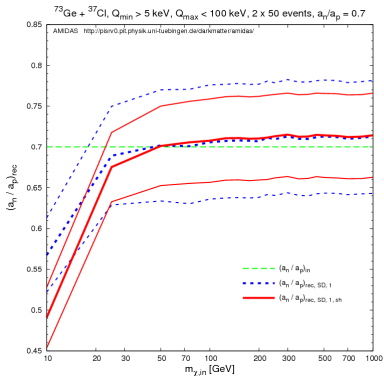
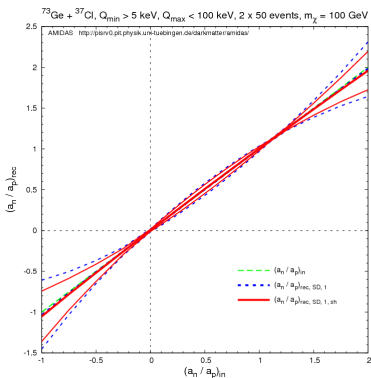
- Determining the ratio of two SD WIMP-nucleon couplings

$$\left(\frac{a_n}{a_p}\right)_{\pm,n}^{\text{SD}} = -\frac{\langle S_p \rangle_X \pm \langle S_p \rangle_Y \mathcal{R}_{J,n}}{\langle S_n \rangle_X \pm \langle S_n \rangle_Y \mathcal{R}_{J,n}}$$

$$\mathcal{R}_{J,n} \equiv \left[ \left(\frac{J_X}{J_X+1}\right) \left(\frac{J_Y+1}{J_Y}\right) \frac{\mathcal{R}_\sigma}{\mathcal{R}_n} \right]^{1/2} \quad (n \neq 0)$$

## Determination of the ratio of two SD WIMP-nucleon couplings

- Reconstructed  $(a_n/a_p)_{\text{rec},1}^{\text{SD}}$   
 $(Q_{\text{min}} > 5 \text{ keV}, Q_{\text{max}} < 100 \text{ keV}, {}^{73}\text{Ge} + {}^{37}\text{Cl}, 2 \times 50 \text{ events},$   
 $m_\chi = 100 \text{ GeV or } a_n/a_p = 0.7)$



[M. Drees, M. Kakizaki, and CLS, in progress]

## Determination of the ratio of two WIMP-proton cross sections

- Differential rate for the combination of the SI and SD cross sections

$$\left(\frac{dR}{dQ}\right)_{\text{expt, } Q=Q_{\min}} = \mathcal{E} \left( \frac{\rho_0 \sigma_0^{\text{SI}}}{2m_\chi m_r^2 N} \right) F_{\text{SI}}^{\prime 2}(Q_{\min}) \cdot \frac{1}{\alpha} \left[ \frac{2r_{\min}/F_{\text{SI}}^{\prime 2}(Q_{\min})}{2Q_{\min}^{1/2} r_{\min}/F_{\text{SI}}^{\prime 2}(Q_{\min}) + I_0} \right]$$

$$F_{\text{SI}}^{\prime 2}(Q) \equiv F_{\text{SI}}^2(Q) + \left( \frac{\sigma_{\text{XP}}^{\text{SD}}}{\sigma_{\text{XP}}^{\text{SI}}} \right) C_p F_{\text{SD}}^2(Q) \quad C_p \equiv \frac{4}{3} \left( \frac{J+1}{J} \right) \left[ \frac{\langle S_p \rangle + (a_n/a_p) \langle S_n \rangle}{A} \right]^2$$

- Determining the ratio of two WIMP-proton cross sections

$$\frac{\sigma_{\text{XP}}^{\text{SD}}}{\sigma_{\text{XP}}^{\text{SI}}} = \frac{F_{\text{SI},Y}^2(Q_{\min},Y) \mathcal{R}_{m,XY} - F_{\text{SI},X}^2(Q_{\min},X)}{C_{p,X} F_{\text{SD},X}^2(Q_{\min},X) - C_{p,Y} F_{\text{SD},Y}^2(Q_{\min},Y) \mathcal{R}_{m,XY}}$$

$$\mathcal{R}_{m,XY} \equiv \left( \frac{r_{\min,X}}{\mathcal{E}_X} \right) \left( \frac{\mathcal{E}_Y}{r_{\min,Y}} \right) \left( \frac{m_Y}{m_X} \right)^2$$

- Determining the ratio of two SD WIMP-nucleon couplings

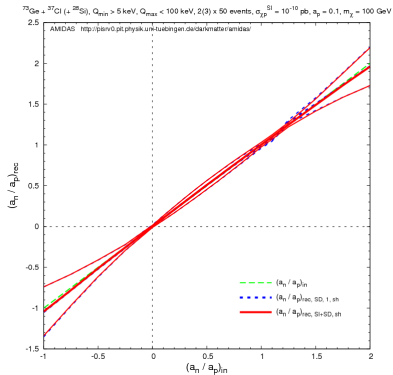
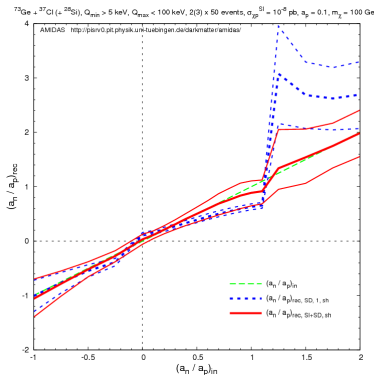
$$\left( \frac{a_n}{a_p} \right)_{\pm}^{\text{SI+SD}} = - \frac{\sqrt{c_{p,X}} \mp \sqrt{c_{p,Y}}}{\sqrt{c_{p,X} s_{n/p,X}} \mp \sqrt{c_{p,Y} s_{n/p,Y}}} \quad (s_{n/p,X} > s_{n/p,Y}, s_{n/p} \equiv \langle S_p \rangle / \langle S_p \rangle)$$

$$c_{p,X} \equiv \frac{4}{3} \left( \frac{J_X + 1}{J_X} \right) \left[ \frac{\langle S_p \rangle_X}{A_X} \right]^2 \left[ F_{\text{SI},Z}^2(Q_{\min},Z) \mathcal{R}_{m,YZ} - F_{\text{SI},Y}^2(Q_{\min},Y) \right] F_{\text{SD},X}^2(Q_{\min},X)$$

[M. Drees and CLS, arXiv:0903.3300]

## Determination of the ratio of two WIMP-proton cross sections

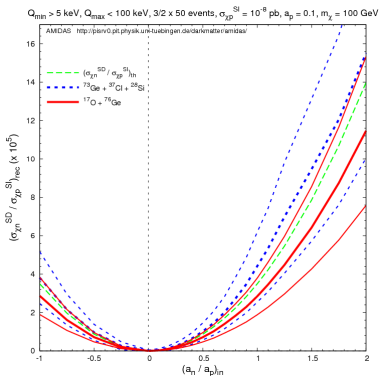
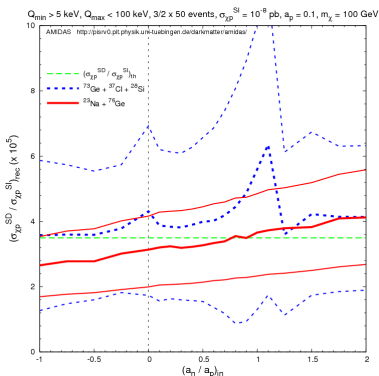
- Reconstructed  $(a_n/a_p)_{rec}^{SI+SD}$  vs  $(a_n/a_p)_{rec,1}^{SD}$   
 $(Q_{min} > 5 \text{ keV}, Q_{max} < 100 \text{ keV}, {}^{73}\text{Ge} + {}^{37}\text{Cl} + {}^{28}\text{Si}, 3 \times 50 \text{ events},$   
 $\sigma_{\chi p}^{SI} = 10^{-8} / 10^{-10} \text{ pb}, a_p = 0.1, m_\chi = 100 \text{ GeV})$



[M. Drees, M. Kakizaki, and CLS, in progress]

## Determination of the ratio of two WIMP-proton cross sections

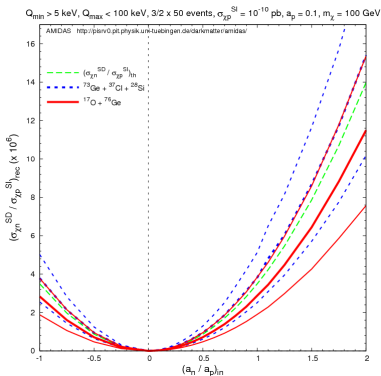
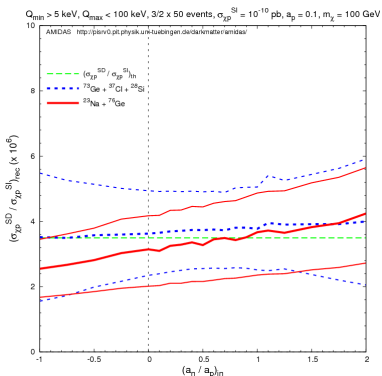
- Reconstructed  $(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{rec}$  and  $(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{rec}$   
 ( $Q_{min} > 5$  keV,  $Q_{max} < 100$  keV,  $^{73}\text{Ge} + ^{37}\text{Cl} + ^{28}\text{Si}$  vs  $^{76}\text{Ge} + ^{23}\text{Na}/^{17}\text{O}$ ,  
 $\sigma_{\chi p}^{SI} = 10^{-8}$  pb,  $a_p = 0.1$ ,  $m_\chi = 100$  GeV,  $3/2 \times 50$  events)



[M. Drees, M. Kakizaki, and CLS, in progress]

## Determination of the ratio of two WIMP-proton cross sections

- Reconstructed  $(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{rec}$  and  $(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{rec}$   
 $(Q_{min} > 5 \text{ keV}, Q_{max} < 100 \text{ keV}, {}^{73}\text{Ge} + {}^{37}\text{Cl} + {}^{28}\text{Si} \text{ vs } {}^{76}\text{Ge} + {}^{23}\text{Na}/{}^{17}\text{O},$   
 $\sigma_{\chi p}^{SI} = 10^{-10} \text{ pb}, a_p = 0.1, m_{\chi} = 100 \text{ GeV}, 3/2 \times 50 \text{ events})$



[M. Drees, M. Kakizaki, and CLS, in progress]



## Estimation of the SI WIMP-nucleon coupling

- We can estimate ratios of each two of the three WIMP-nucleon cross sections model-independently.
  - ◆ Can we estimate **any one** of them further?
  - ◆ Unfortunately, **no!**
  - ◆ Expression for the product of the local density times the WIMP-nucleus cross section

$$\rho_0 \sigma_0^{\text{SI}} = \left( \frac{1}{\mathcal{E}} \right) m_\chi m_{r,N} \sqrt{\frac{m_N}{2}} \left[ \frac{2Q_{\min}^{1/2} r_{\min}}{F_{\text{SI}}^2(Q_{\min})} + I_0 \right]$$

➔ Making an **assumption** for the local WIMP density

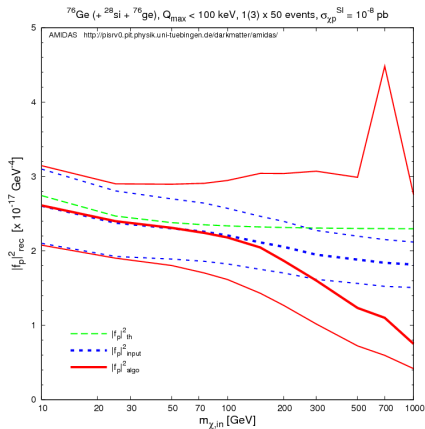
- Estimating the SI WIMP-nucleon coupling

$$|f_p|^2 = \frac{1}{\rho_0} \left[ \frac{\pi}{4\sqrt{2}} \left( \frac{1}{\mathcal{E} A^2 \sqrt{m_N}} \right) \right] \left[ \frac{2Q_{\min}^{1/2} r_{\min}}{F_{\text{SI}}^2(Q_{\min})} + I_0 \right] (m_\chi + m_N)$$

[M. Drees and CLS, arXiv:0809.2441]

## Estimation of the SI WIMP-nucleon coupling

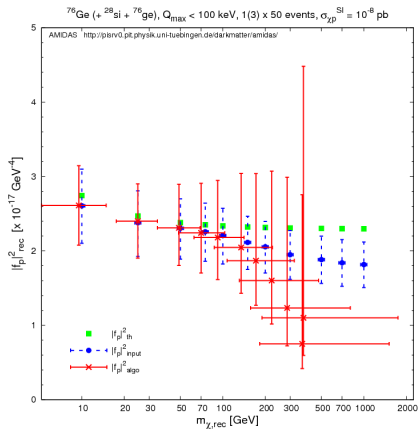
- Reconstructed  $|f_p|_{\text{rec}}^2$   
 $(Q_{\text{max}} < 100 \text{ keV}, {}^{76}\text{Ge} (+ {}^{28}\text{Si} + {}^{76}\text{Ge}), \sigma_{\chi p}^{\text{SI}} = 10^{-8} \text{ pb}, 1(3) \times 50 \text{ events})$



[M. Drees and CLS, in progress]

## Estimation of the SI WIMP-nucleon coupling

- Reconstructed  $|f_p|_{\text{rec}}^2$  vs. reconstructed  $m_{\chi, \text{rec}}$   
 ( $Q_{\text{max}} < 100 \text{ keV}$ ,  $^{76}\text{Ge} (+^{28}\text{Si} + ^{76}\text{Ge})$ ,  $\sigma_{\chi p}^{\text{SI}} = 10^{-8} \text{ pb}$ ,  $1(3) \times 50 \text{ events}$ )



[M. Drees and CLS, in progress]

## Summary

- Once two or more experiments **with different target nuclei** observe positive WIMP signals, we could estimate
  - ◆ WIMP mass  $m_\chi$
  - ◆ SI WIMP-proton coupling  $|f_p|^2$
  - ◆ ratio between the SD WIMP-nucleon couplings,  $a_n/a_p$
  - ◆ ratios between the SD and SI WIMP-nucleon cross sections,  $\sigma_{\chi p/n}^{SD}/\sigma_{\chi p}^{SI}$
- These analyses are **independent** of the **velocity distribution**, the **local density**, and the **mass/couplings on nucleons** of halo WIMPs (none of them is yet known).
- For a WIMP mass of **100 GeV**, these quantities could be estimated with statistical errors of **10 – 40%** with only  **$\mathcal{O}(50)$  events** from one experiment.

## Summary

- These information will help us to
  - ◆ constrain the parameter space
  - ◆ distinguish the (neutralino) LSP from the (first KK hypercharge) LKP
    - G. Bertone *et al.*, PRL 99, 151301 (2007); V. Barger *et al.*, PRD 78, 056007 (2008);  
G. Belanger *et al.*, PRD 79, 015008 (2009); R. C. Cotta *et al.*, arXiv:0903.4409 (2009)
  - ◆ identify the particle produced at colliders to be indeed halo WIMP
  - ◆ predict the WIMP annihilation cross section  $\langle \sigma_{\text{anni}} v \rangle$
  - ◆ .....
  
- Furthermore, we could
  - ◆ determine the local WIMP density  $\rho_0$
  - ◆ predict the indirect detection event rate  $d\Phi/dE$
  - ◆ test our understanding of the early Universe
  - ◆ .....

## Current projects and related research interests

- With **direct DM detection** experiments
  - ◆ Identifying the **annual modulation** of WIMP events
  - ◆ Extracting **directional information** of WIMP signals
  - ◆ Taking **background events** into account
  - ◆ (Online interactive) simulation/data analysis system: **AMIDAS**  
<http://pisrv0.pit.physik.uni-tuebingen.de/darkmatter/amidas/>
  
- With **indirect DM detection** experiments
  - ◆ Predicting the **WIMP annihilation cross section** and the **event rate**
  - ◆ Information on the **spin-dependent WIMP-proton coupling**
  - ◆ Information on (the **anisotropy** of) the **halo structure**

Thank you very much for your attention

[<http://dmrc.snu.ac.kr/~cshan/>]