

PPP8, May 21 2009

Determining the neutrino flavor ratio at the astrophysical source

By G.-L. Lin

NCTU and LeCosPA

Determination of the Neutrino Flavor Ratio at the Astrophysical Source

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Abstract

We discuss the reconstruction of neutrino flavor ratios at astrophysical sources from future neutrino-telescope measurements, given the knowledge of neutrino mixing angles obtained from terrestrial experiments. With a statistical analysis, we demonstrate that the pion source and the muon damped source can be distinguished at the 3σ level provided the accuracies on measuring $R \equiv \phi(\nu_\mu)/(\phi(\nu_e) + \phi(\nu_\tau))$ and $S \equiv \phi(\nu_e)/\phi(\nu_\tau)$ can both reach about 10%. On the other hand, the above two sources are very difficult to distinguish by merely measuring R alone. We also discuss the effect of leptonic CP phase on such a flavor-ratio reconstruction.

PACS numbers: 95.85.Ry, 14.60.Pq, 95.55.Vj

β source

$(1,0,0)$

**Pion
source**

$(1/3, 2/3, 0)$

$(0,1,0)$

$(0,0,1)$

$\downarrow \frac{1}{\sqrt{6}}(-2,1,1)$

**Damped muon
source**

$\longrightarrow \frac{1}{\sqrt{2}}(0,-1,1)$

$\Phi_0 = (\phi_0(\nu_e), \phi_0(\nu_\mu), \phi_0(\nu_\tau))$

$\phi_0(\nu_e) + \phi_0(\nu_\mu) + \phi_0(\nu_\tau) = 1$

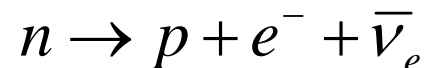
The β source (1,0,0)

Motivated by the correlation of the arrival direction of the cosmic rays to the Galactic Plane (GP) near EeV (10^{18} eV) energies

AGASA 1998; Fly's Eye 1998

Directional signal requires relatively-stable neutral primaries.

Neutron decay length is about 10 kpc for $E_n=1$ EeV. Smaller energy neutrons can decay



L. A. Anchordoqui, H. Goldberg, F. Halzen and T. J. Weiler, 2004

Pion source (1/3,2/3,0)

$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$

$$\mu^+ \rightarrow \bar{\nu}_\mu + e^+ + \nu_e$$

Energies of various neutrinos are comparable, i.e., muon decays before losing its energy by interactions.

Cosmogenic (GZK) neutrinos produced by $p + \gamma_{CMB} \rightarrow \Delta^+ \rightarrow n + \pi^+$ and the subsequent pion decay fit into this category.

Damped muon source (0,1,0)

$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$

$$\mu^+ \rightarrow \bar{\nu}_\mu + e^+ + \nu_e$$

Muon loses significant amount of energy before it decays:

(1) muon interacts with matter

J. P. Rachen and P. Meszaros, 1998

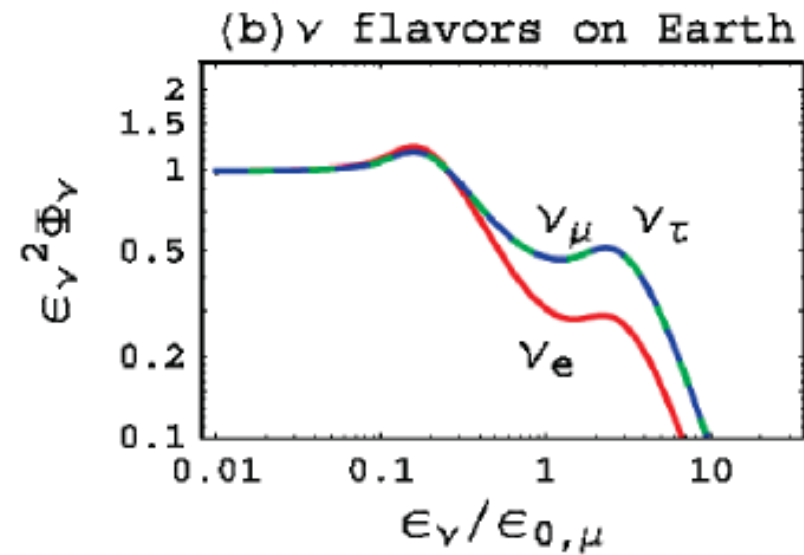
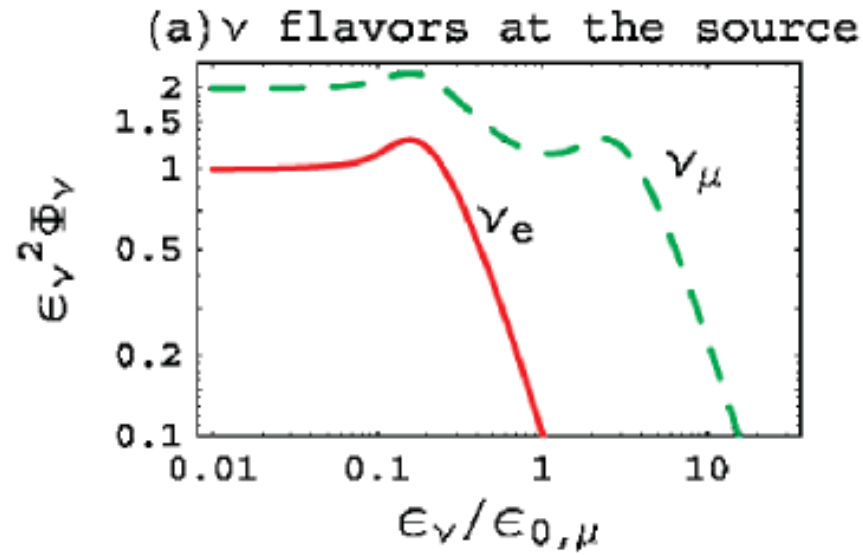
(2) Muon interacts with background photon field

**M. Kacherliess, O. Ostapchenko and R. Tomas,
arXiv: 0708.3007**

See also

T. Kashti and E. Waxman *Phy. Rev. Lett.* 2005

Neutrinos from muon decays are out of the spectrum



T. Kashti and E. Waxman *Phys. Rev. Lett.* 2005

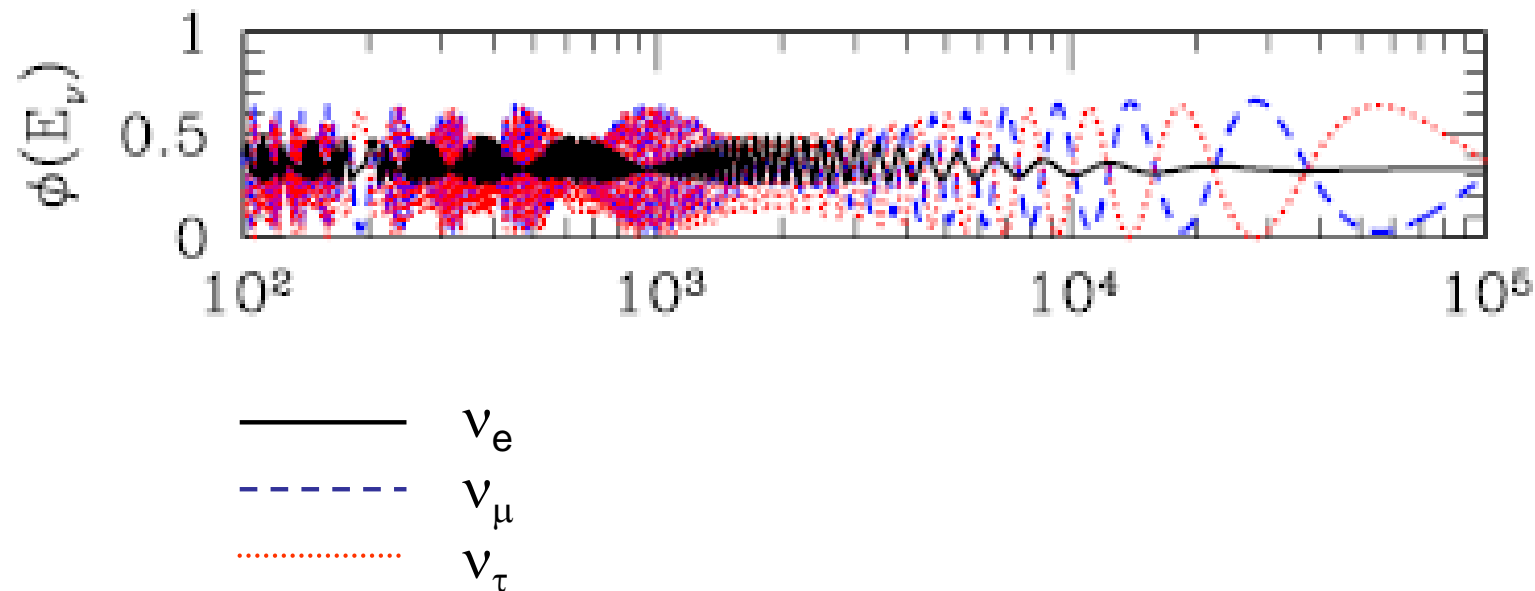
Transition from pion source to muon-damped source

Source with a significant tau neutrino flux

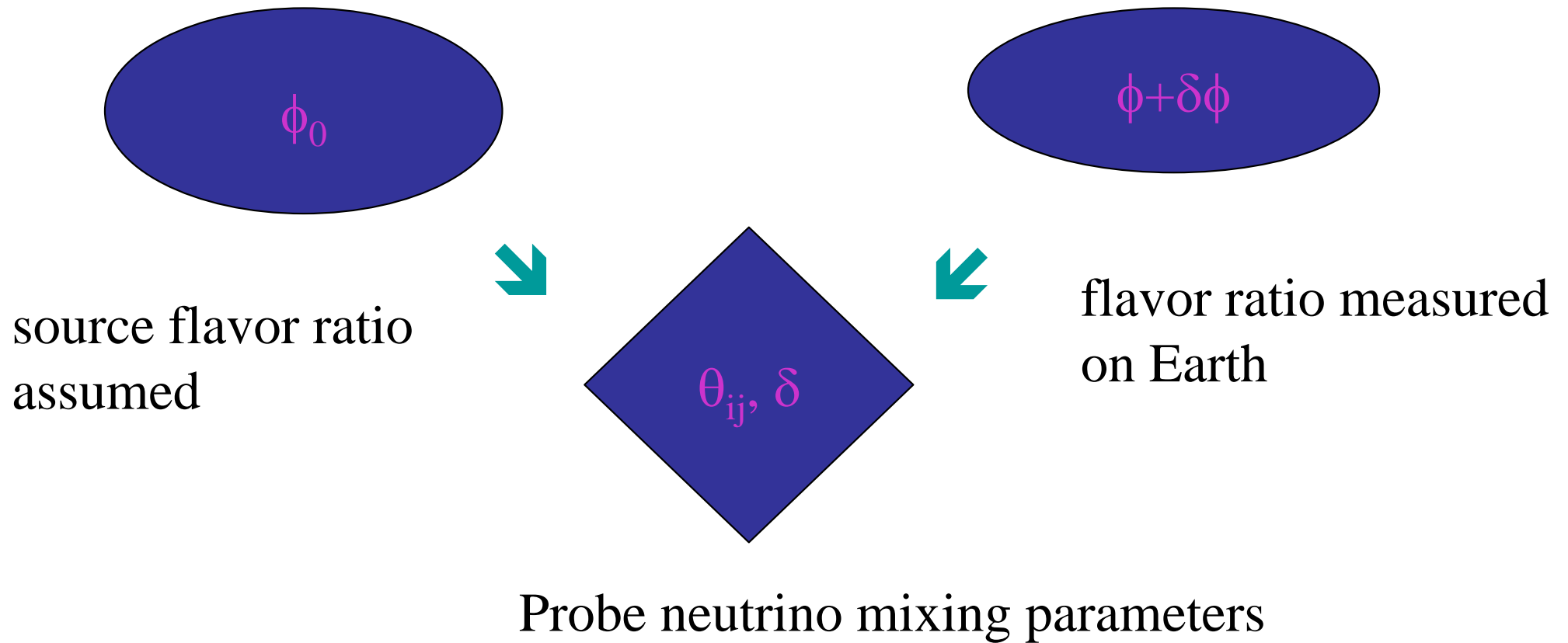
Optically thick sources: highly relativistic GRB jets

Neutrinos already oscillate inside the object.

O. Mena, I. Mocioiu and S. Razzaque, 2006



The usual approach



S. Pakvasa, Mod. Phys. Lett. A **19**, 1163 (2004) [Yad. Fiz. **67**, 1179 (2004)].

M. L. Costantini and F. Vissani, Astropart. Phys. **23**, 477 (2005); F. Vissani, Astropart. Phys. **26**, 310 (2006); astro-ph/0609575.

P. Bhattacharjee and N. Gupta, arXiv:hep-ph/0501191.

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W. Winter, Phys. Rev. D **74**, 033015 (2006).

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D. Majumdar and A. Ghosal, Phys. Rev. D **75**, 113004 (2007).

W. Rodejohann, JCAP **0701**, 029 (2007).

D. Meloni and T. Ohlsson, Phys. Rev. D **75**, 125017 (2007).

K. Blum, Y. Nir and E. Waxman, arXiv:0706.2070 [hep-ph].

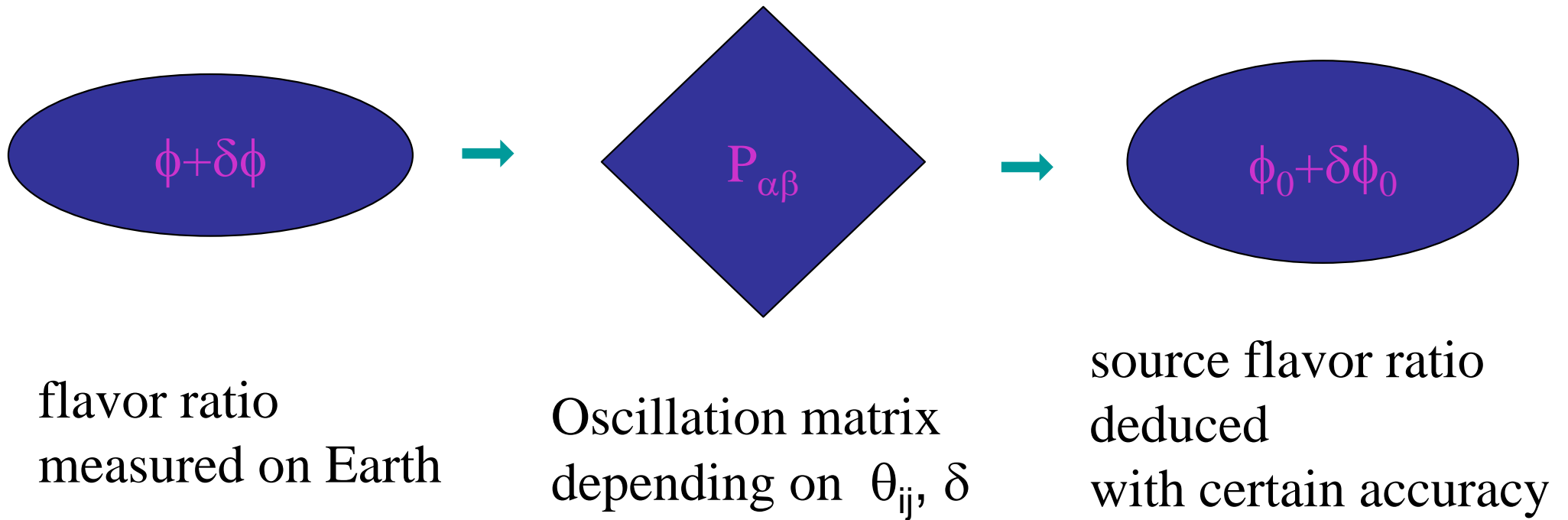
G. R. Hwang and S. Kim, Phys. Rev. D **78**, 093008 (2008).

S. Pakvasa, W. Rodejohann and T. J. Weiler, JHEP **0802**, 005 (2008).

S. Choubey, V. Niro and W. Rodejohann, Phys. Rev. D **77**, 113006 (2008).

M. Maltoni and W. Winter, JHEP **0807**, 064 (2008).

Our approach



Reconstructing the source flavor ratio

$$(\phi(\nu_e), \phi(\nu_\mu), \phi(\nu_\tau)) = (\phi_0(\nu_e), \phi_0(\nu_\mu), \phi_0(\nu_\tau)) \begin{pmatrix} P_{ee} & P_{e\mu} & P_{e\tau} \\ P_{\mu e} & P_{\mu\mu} & P_{\mu\tau} \\ P_{\tau e} & P_{\tau\mu} & P_{\tau\tau} \end{pmatrix}$$

Measured flux

$$P_{\alpha\beta} \equiv P(\nu_\alpha \rightarrow \nu_\beta) = \sum_{i=1}^3 |U_{\alpha i}|^2 |U_{\beta i}|^2, \text{ where } \nu_\alpha = U_{\alpha i} \nu_i$$

P

Flavor Eigenstate

Mass Eigenstate

$U_{\alpha i}$ contains 3 mixing angles-- θ_{12} , θ_{23} , and θ_{13}
one CP phase δ_{CP}

$$\begin{aligned}
P_{ee} &= \left(1 - \frac{1}{2}\omega\right) (1 - D^2)^2 + D^4, \\
P_{e\mu} &= \frac{1}{4}(1 - D^2) \left[\omega(1 + \Delta) + (4 - \omega)(1 - \Delta)D^2 + 2\sqrt{\omega(1 - \omega)(1 - \Delta^2)}D \cos \delta \right], \\
P_{e\tau} &= \frac{1}{4}(1 - D^2) \left[\omega(1 - \Delta) + (4 - \omega)(1 + \Delta)D^2 - 2\sqrt{\omega(1 - \omega)(1 - \Delta^2)}D \cos \delta \right], \\
P_{\mu\mu} &= \frac{1}{2} \left[(1 + \Delta^2) - (1 - \Delta)^2 D^2 (1 - D^2) \right] \\
&\quad - \frac{1}{8}\omega \left[(1 + \Delta)^2 + (1 - \Delta)^2 D^4 - (1 - \Delta^2)D^2(2 + 4\cos^2 \delta) \right] \\
&\quad - \frac{1}{2}\sqrt{\omega(1 - \omega)(1 - \Delta^2)} \left[(1 + \Delta) - (1 - \Delta)D^2 \right] D \cos \delta, \\
P_{\mu\tau} &= \frac{1}{2}(1 - \Delta^2)(1 - D^2 + D^4) \\
&\quad - \frac{1}{8}\omega \left[(1 - \Delta^2)(1 + 4D^2 \cos^2 \delta + D^4) - 2(1 + \Delta^2)D^2 \right] \\
&\quad + \frac{1}{2}\sqrt{\omega(1 - \omega)(1 - \Delta^2)}\Delta(1 + D^2)D \cos \delta, \\
P_{\tau\tau} &= \frac{1}{2} \left[(1 + \Delta^2) - (1 + \Delta)^2 D^2 (1 - D^2) \right] \\
&\quad - \frac{1}{8}\omega \left[(1 - \Delta)^2 + (1 + \Delta)^2 D^4 - (1 - \Delta^2)D^2(2 + 4\cos^2 \delta) \right] \\
&\quad + \frac{1}{2}\sqrt{\omega(1 - \omega)(1 - \Delta^2)} \left[(1 - \Delta) - (1 + \Delta)D^2 \right] D \cos \delta,
\end{aligned}
\tag{A1}$$

$\omega \equiv \sin^2 2\theta_{12}, \Delta \equiv \cos 2\theta_{23}, D \equiv \sin \theta_{13}$
 δ CP phase

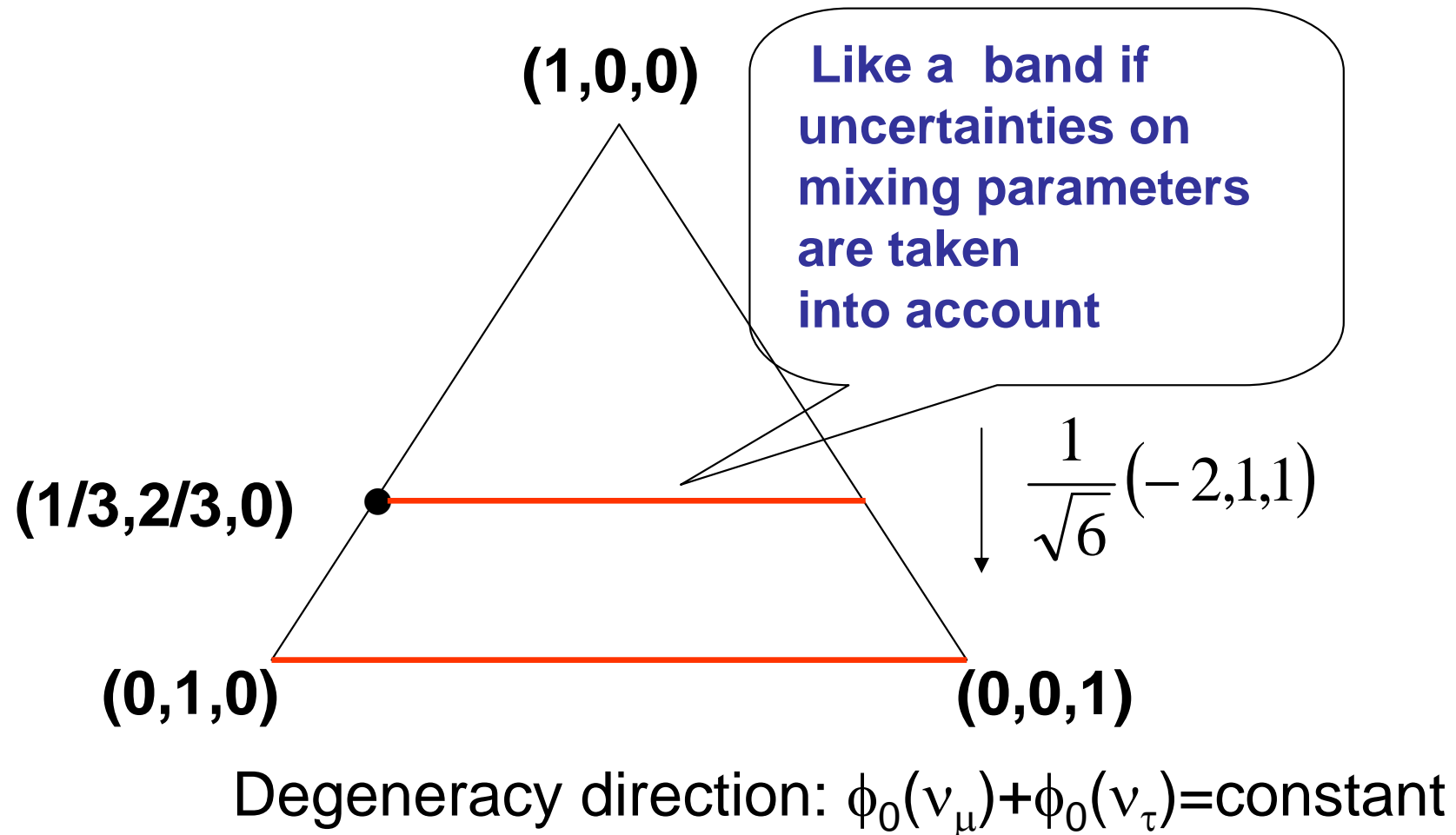
$$\theta_{23}=45^\circ, \theta_{13}=0^\circ$$

$$\mathbf{P} = \frac{1}{8} \begin{pmatrix} 8-4\omega & 2\omega & 2\omega \\ 2\omega & 4-\omega & 4-\omega \\ 2\omega & 4-\omega & 4-\omega \end{pmatrix},$$

where $\omega = \sin^2 2\theta_{12}$

This matrix is singular!

Uncertainties: mixing parameter and measurements

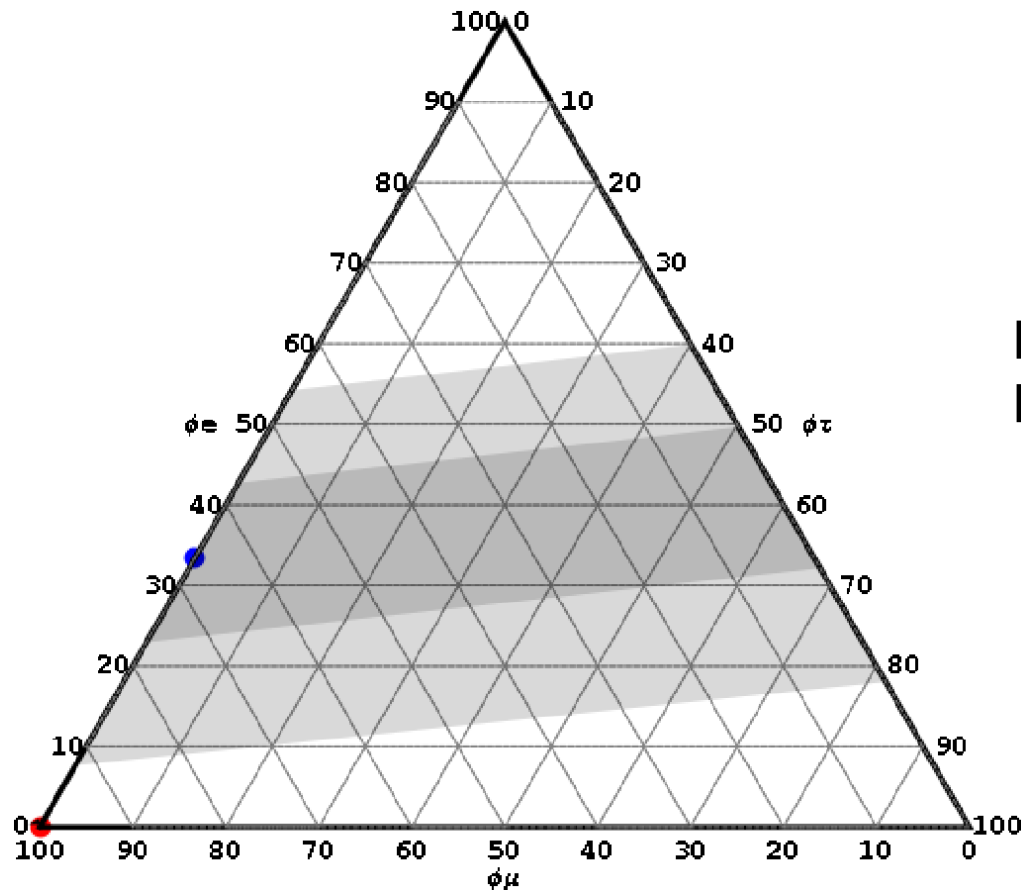


$$\longrightarrow \frac{1}{\sqrt{2}} (0, -1, 1)$$

ν flavor astronomy require
a large number of events

Pion Source

Require statistical error $\sim 10\%$
So ~ 100 events needed.



$$\sin^2 \theta_{12} = 0.32^{+0.02}_{-0.02},$$

$$\sin^2 \theta_{23} = 0.45^{+0.09}_{-0.06},$$

$$\sin^2 \theta_{13} < 0.019 \text{ (90\% C.L.)}$$

$$\delta_{\text{CP}} = 0$$

M.C. Gonzalez-Garcia and
M. Maltoni, Phys. Rept. 2008

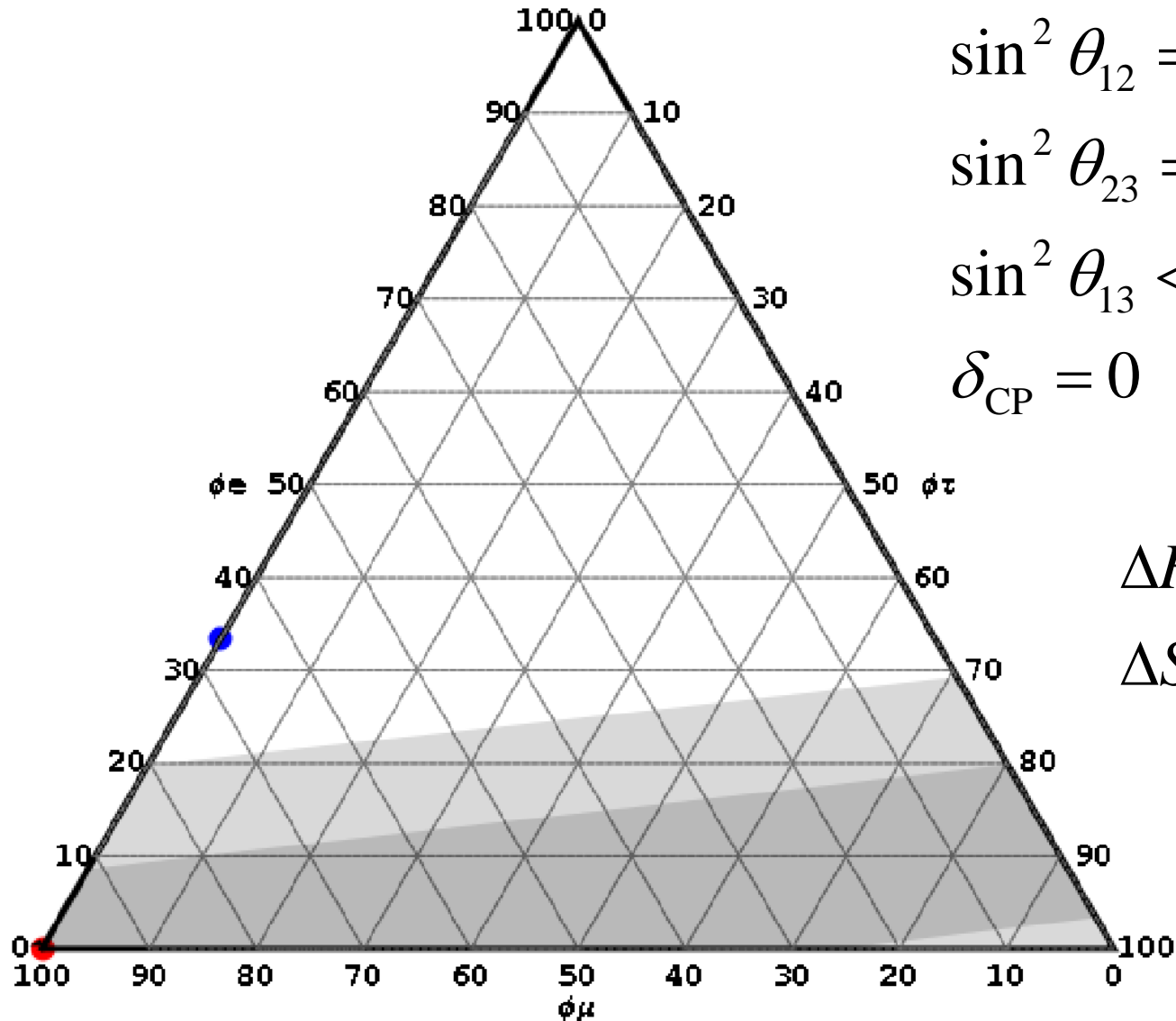
$$R \equiv \phi(\nu_\mu) / (\phi(\nu_e) + \phi(\nu_\tau))$$

$$S \equiv \phi(\nu_e) / \phi(\nu_\tau)$$

$$\Delta R / R = 10\%$$

$$\Delta S / S = 1.2 \Delta R / R$$

Damped Muon Source



$$\sin^2 \theta_{12} = 0.32^{+0.02}_{-0.02},$$

$$\sin^2 \theta_{23} = 0.45^{+0.09}_{-0.06},$$

$$\sin^2 \theta_{13} < 0.019 \text{ (90\% C.L.)}$$

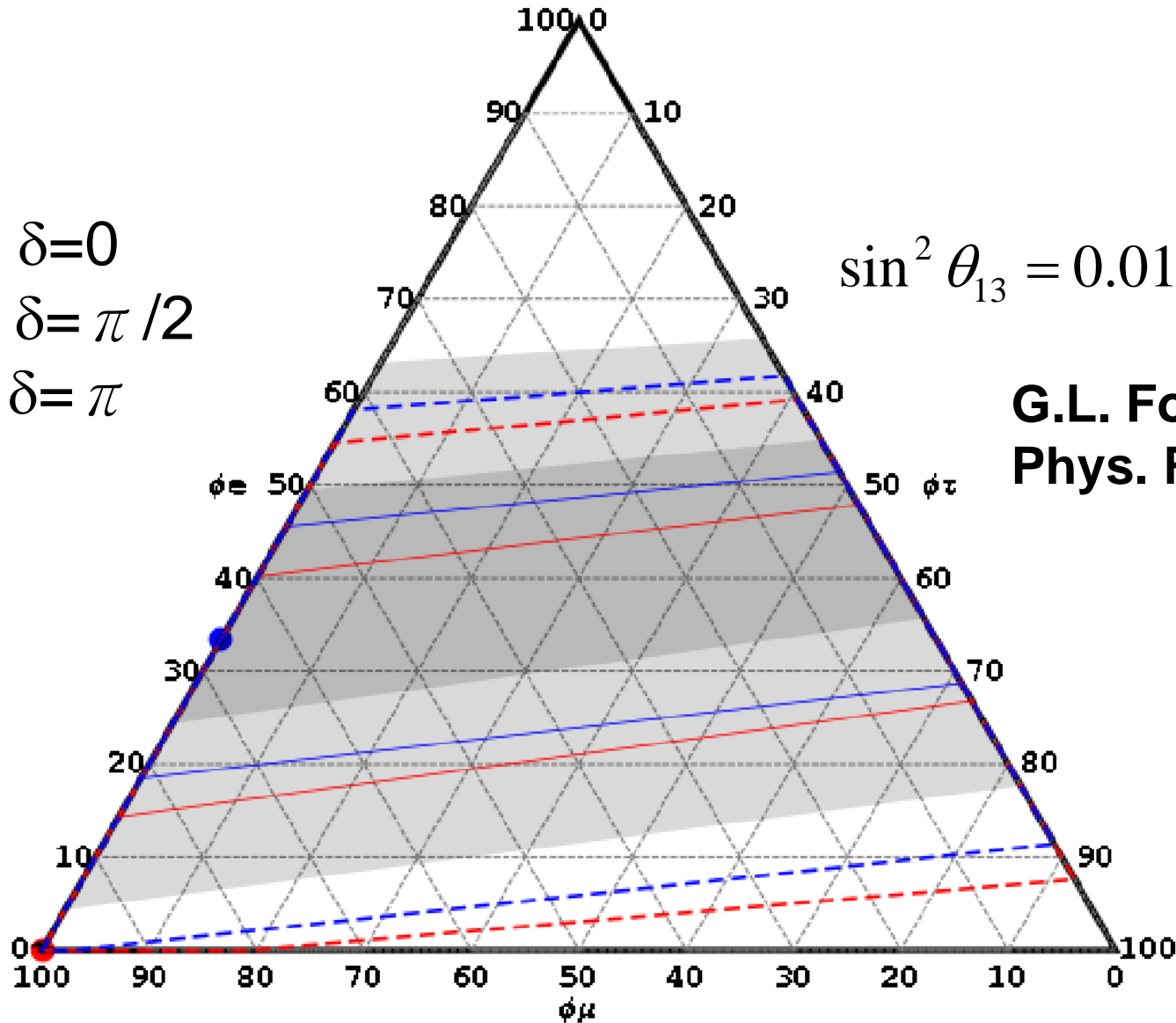
$$\delta_{\text{CP}} = 0$$

$$\Delta R / R = 10\%$$

$$\Delta S / S = 1.3 \Delta R / R$$

Pion Source

Gray: $\delta=0$
Red: $\delta=\pi/2$
Blue: $\delta=\pi$



$$\sin^2 \theta_{13} = 0.016 \pm 0.010(1\sigma)$$

**G.L. Fogli et al.,
Phys. Rev. Lett. 2008**

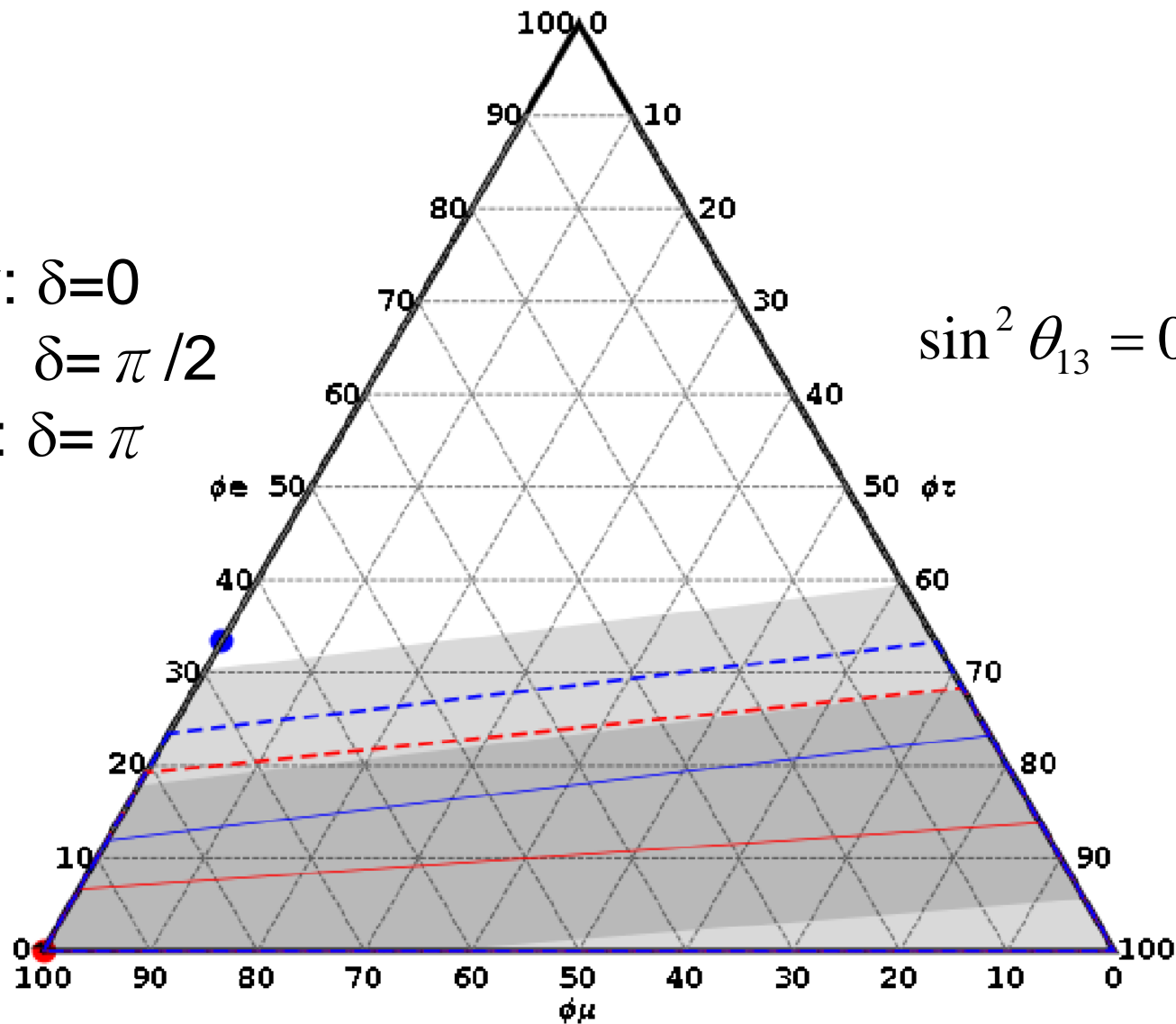
Damped muon source

Gray: $\delta=0$

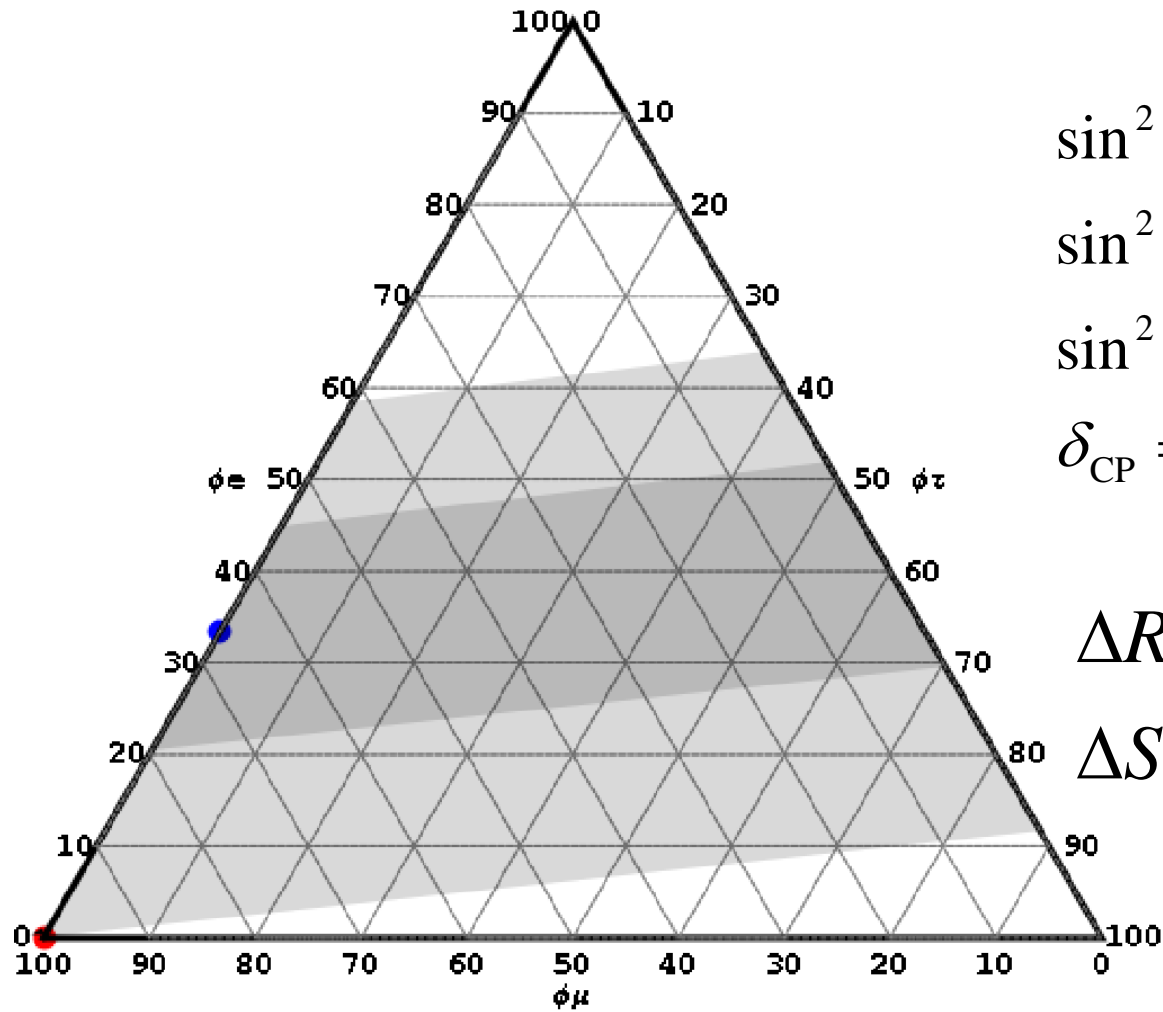
Red: $\delta=\pi/2$

Blue: $\delta=\pi$

$$\sin^2 \theta_{13} = 0.016 \pm 0.010(1\sigma)$$



Critical measurement accuracy for ruling out damped muon source



$$\sin^2 \theta_{12} = 0.32^{+0.02}_{-0.02},$$

$$\sin^2 \theta_{23} = 0.45^{+0.09}_{-0.06},$$

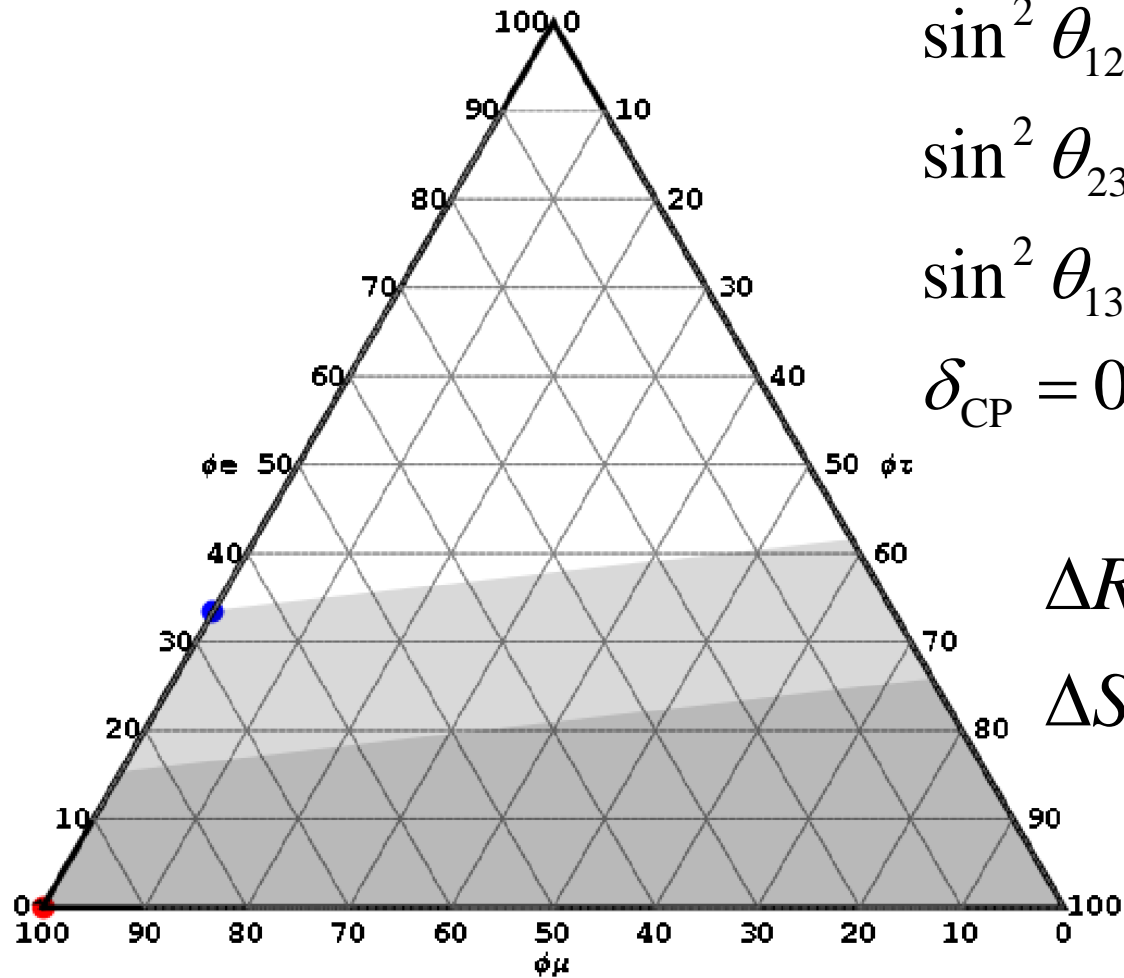
$$\sin^2 \theta_{13} < 0.019 \text{ (90\% C.L.)}$$

$$\delta_{\text{CP}} = 0$$

$$\Delta R / R \sim 13\%$$

$$\Delta S / S = 1.2 \Delta R / R$$

Critical measurement accuracy for ruling out pion source



$$\sin^2 \theta_{12} = 0.32^{+0.02}_{-0.02},$$

$$\sin^2 \theta_{23} = 0.45^{+0.09}_{-0.06},$$

$$\sin^2 \theta_{13} < 0.019 \text{ (90\% C.L.)}$$

$$\delta_{\text{CP}} = 0$$

$$\Delta R / R \sim 18\%$$

$$\Delta S / S = 1.3 \Delta R / R$$

What if one only measures R ?

It is more challenging to measure S
Relies on double bang signature of ν_τ

Pion source

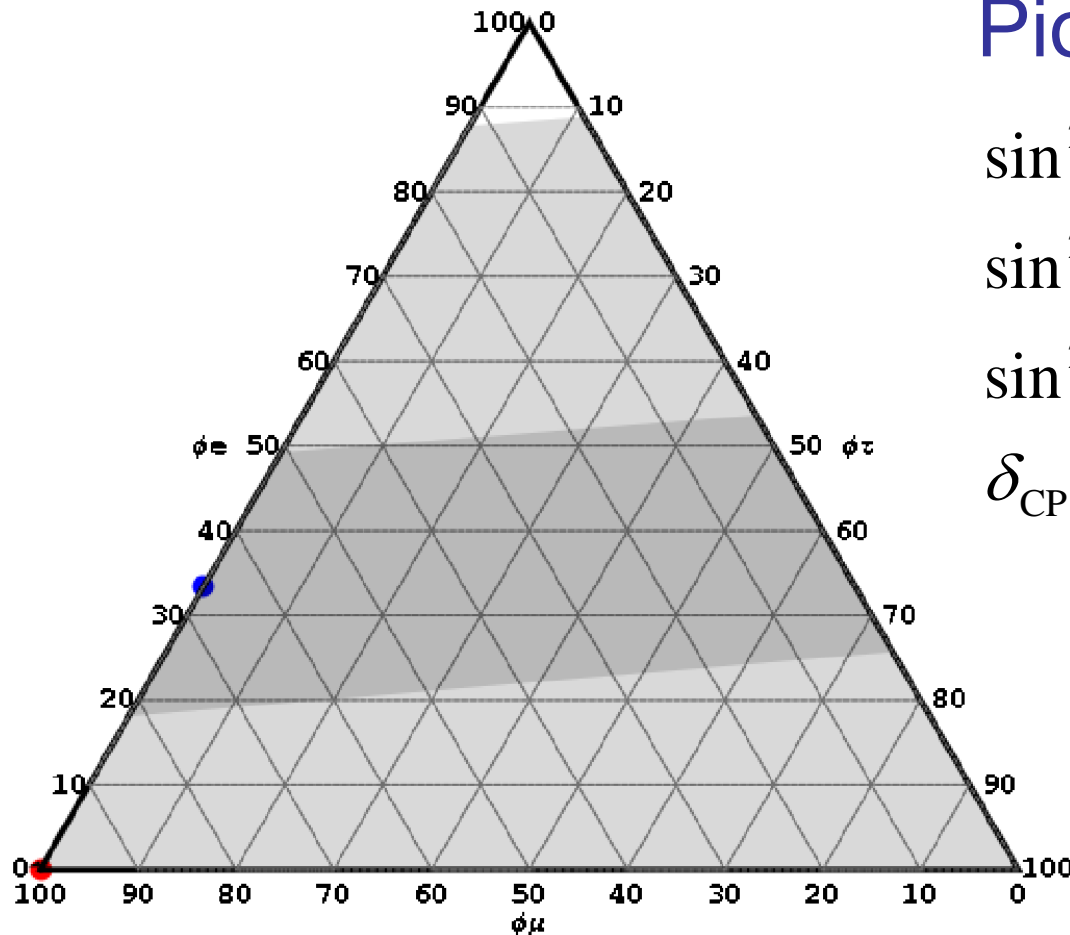
$$\sin^2 \theta_{12} = 0.32^{+0.02}_{-0.02},$$

$$\sin^2 \theta_{23} = 0.45^{+0.09}_{-0.06},$$

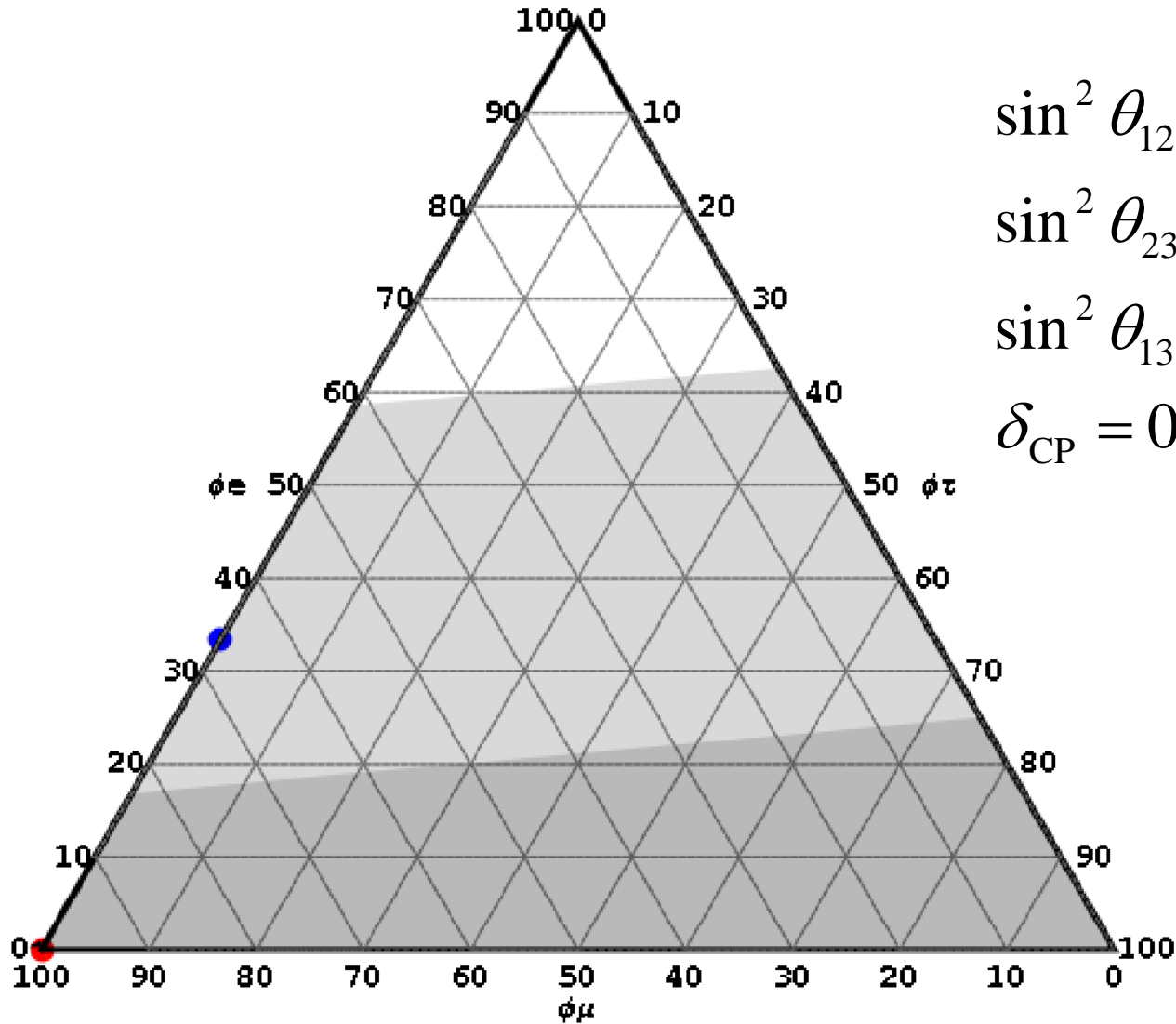
$$\sin^2 \theta_{13} < 0.019 \text{ (90\% C.L.)}$$

$$\delta_{\text{CP}} = 0$$

$$\Delta R / R = 10\%$$



Damped muon source



$$\sin^2 \theta_{12} = 0.32^{+0.02}_{-0.02},$$

$$\sin^2 \theta_{23} = 0.45^{+0.09}_{-0.06},$$

$$\sin^2 \theta_{13} < 0.019 \text{ (90\% C.L.)}$$

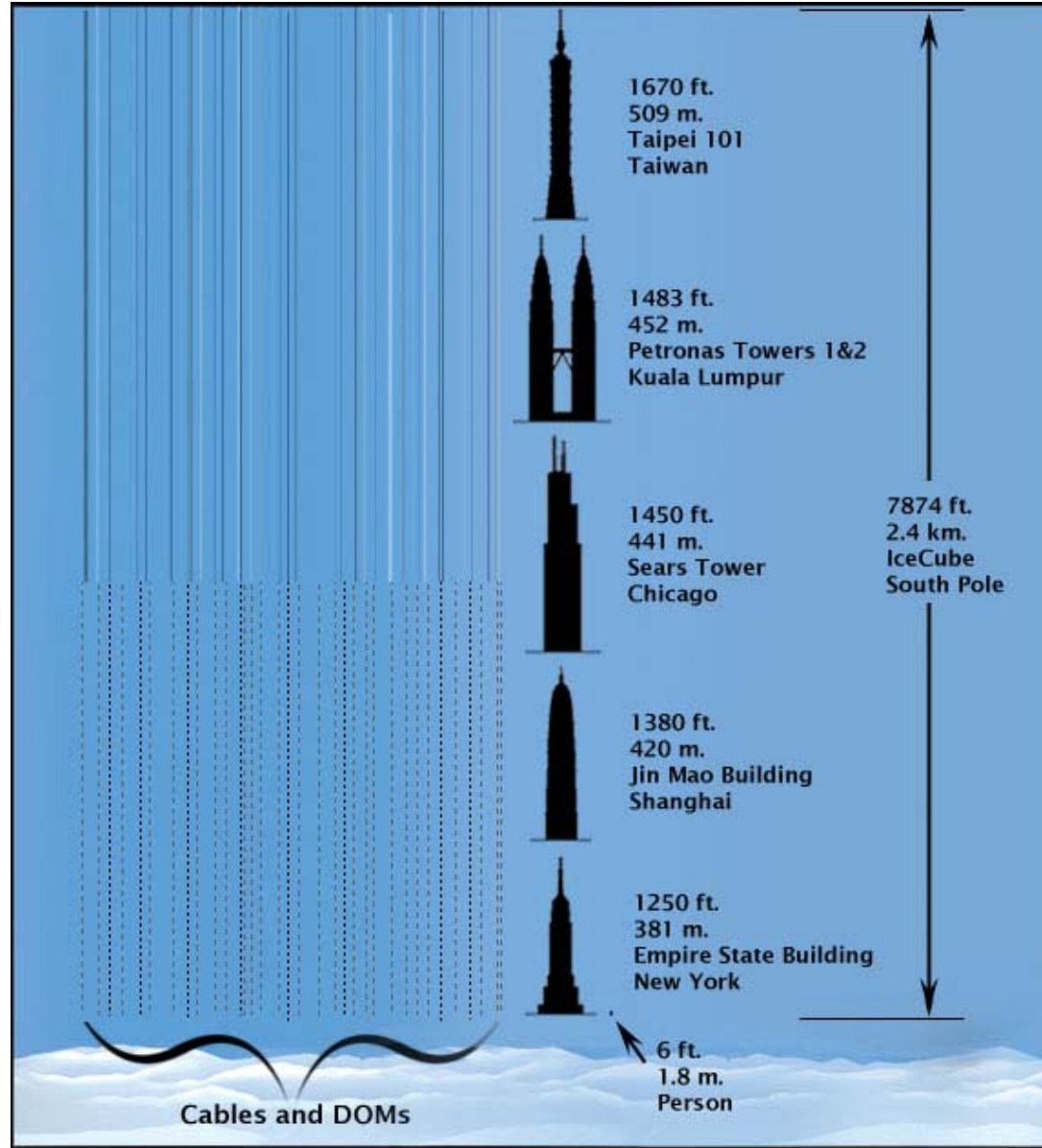
$$\delta_{\text{CP}} = 0$$

$$\Delta R / R = 10\%$$

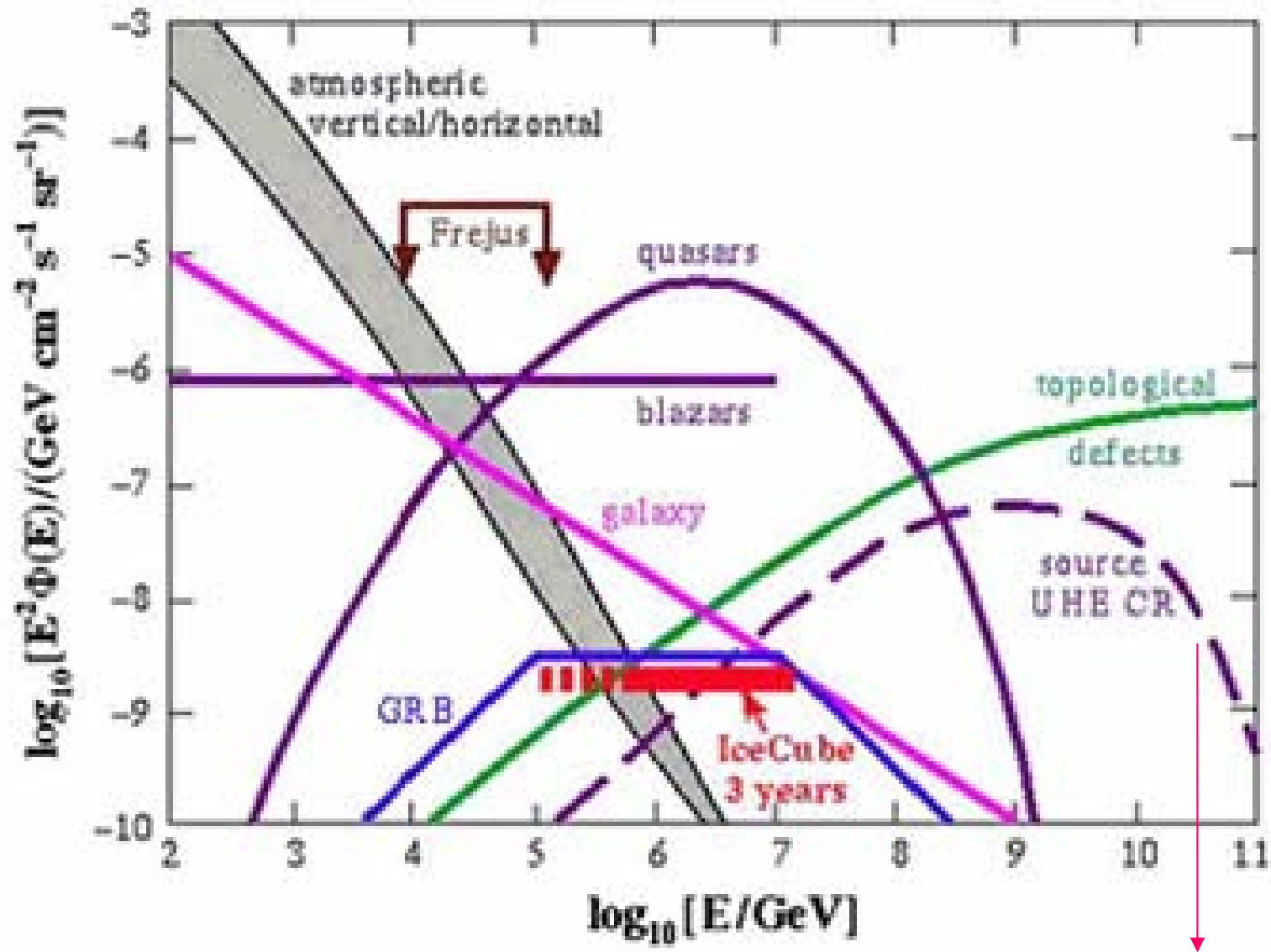
How to detect astrophysical neutrinos ?

- IceCube—PMT array in South Pole ice
- ANTARES→KM3Net—PMT array in the Mediterranean
- ANITA—Radio wave detector above South Pole
- Pierre Auger—Earth skimming tau neutrinos

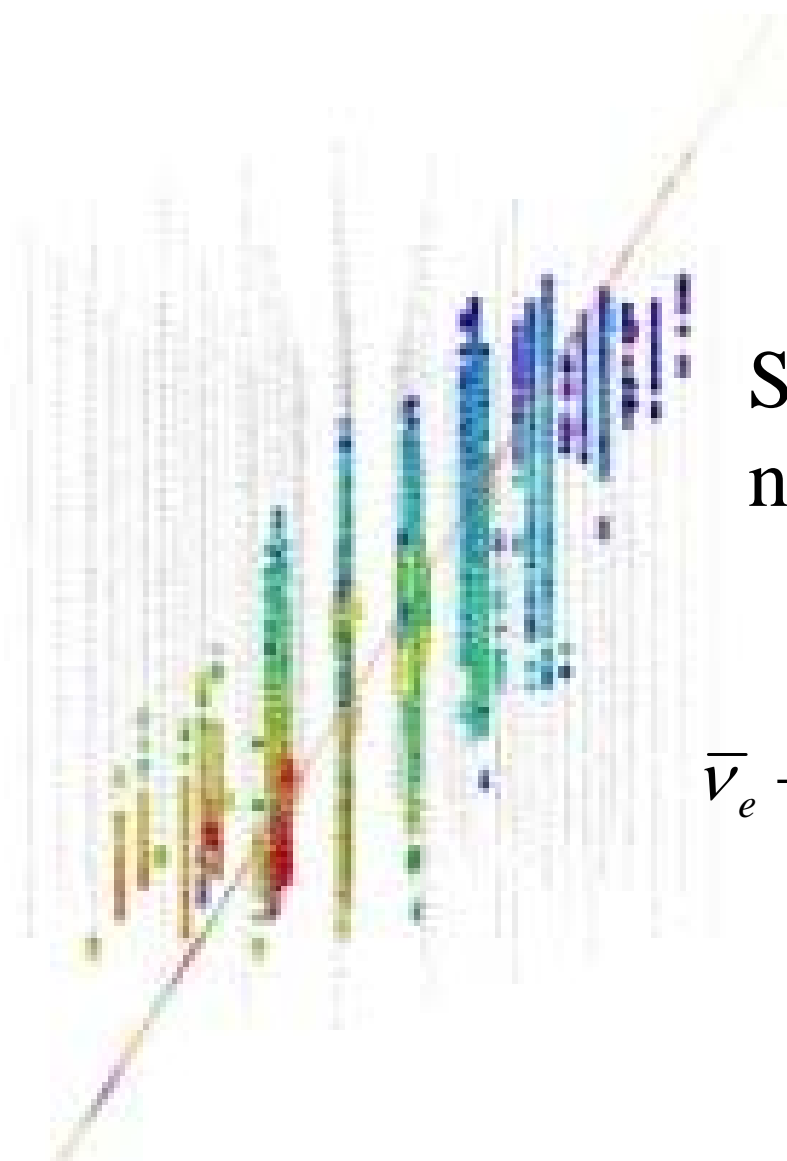
IceCube in Scale



IceCube—PMT array in South Pole ice



Cosmogenic neutrino flux



Simulated high energy
neutrino event

$$\bar{\nu}_e + e^- \rightarrow W^- \rightarrow \bar{\nu}_\mu + \mu^-$$

E = 6400 TeV

ANITA

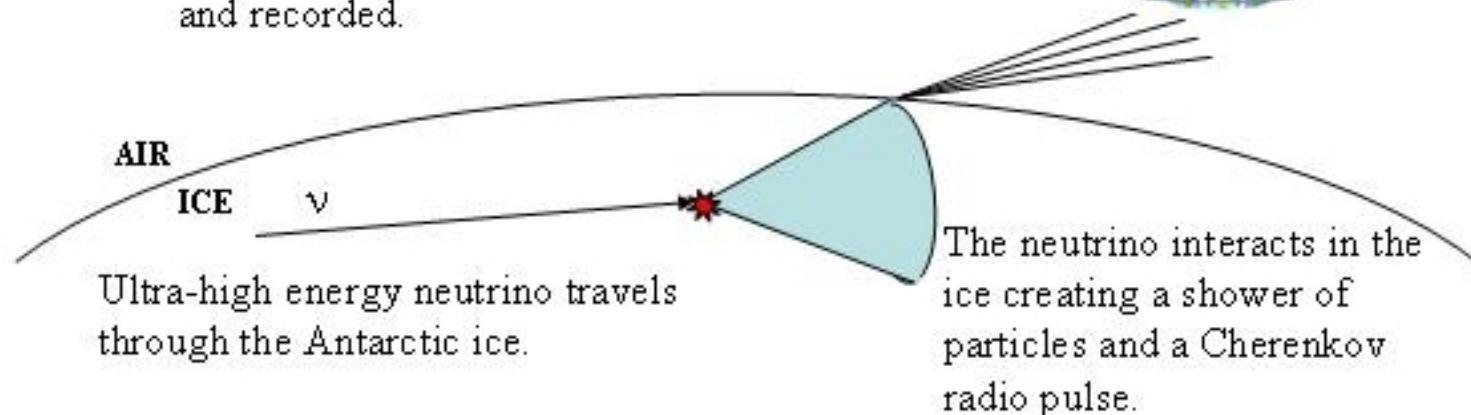
ANTARCTIC IMPULSIVE TRANSIENT ANTENNA

A LONG DURATION BALLOON MISSION TO CONSTRAIN THE ORIGIN OF THE HIGHEST ENERGY PARTICLES IN THE UNIVERSE

**ANITA
Experiment
at altitude of
~120,000 ft**



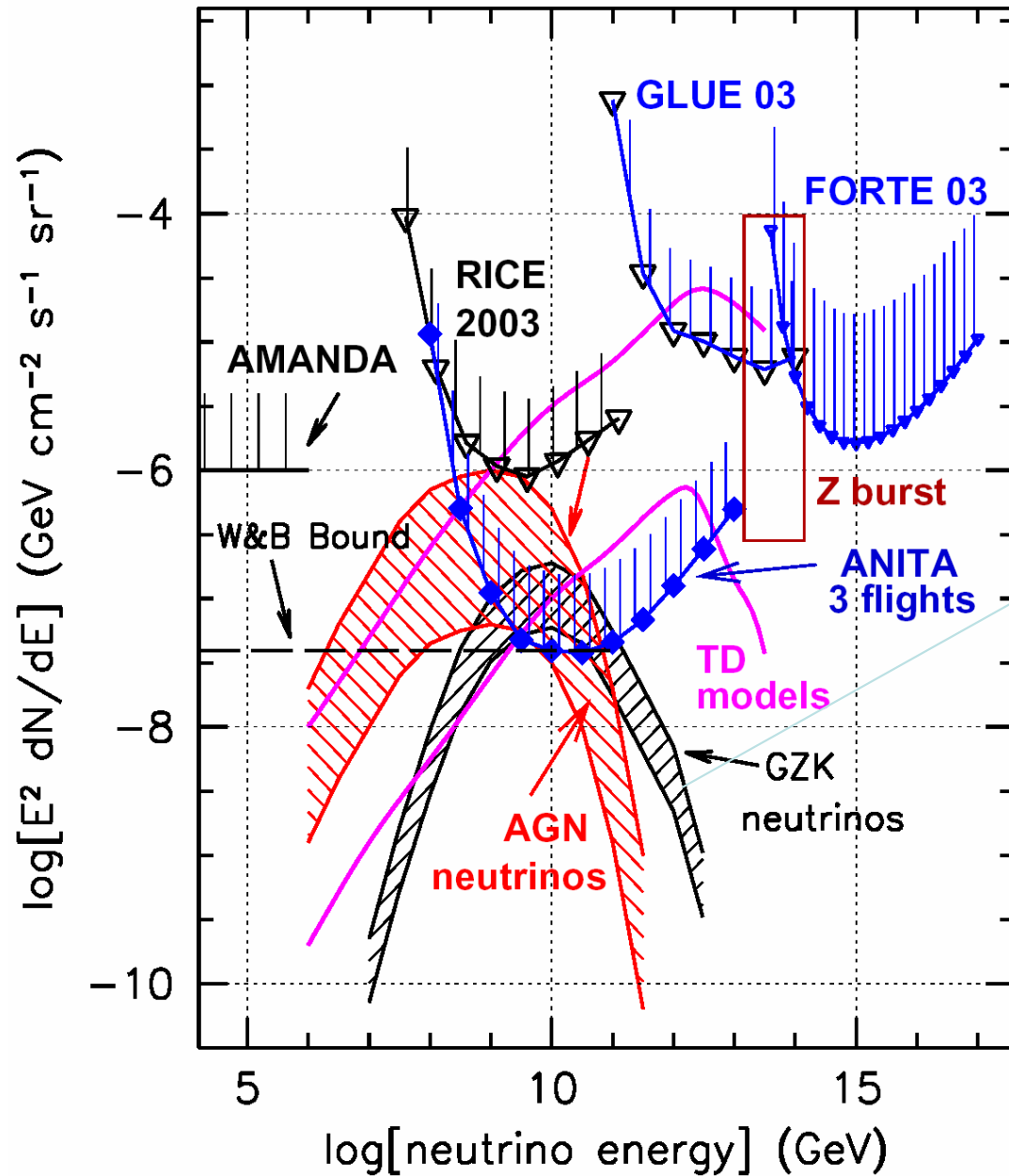
The pulse propagates through the ice, is refracted at the ice surface of the ice and travels to the ANITA instrument antennas where the waveform is digitized and recorded.



Ultra-high energy neutrino travels through the Antarctic ice.

The neutrino interacts in the ice creating a shower of particles and a Cherenkov radio pulse.

Neutrino Models & Limits, mid-2003



Cosmogenic neutrinos

ANITA—best sensitivity
near 10^{19} eV

Most recent launch Jan. 2009

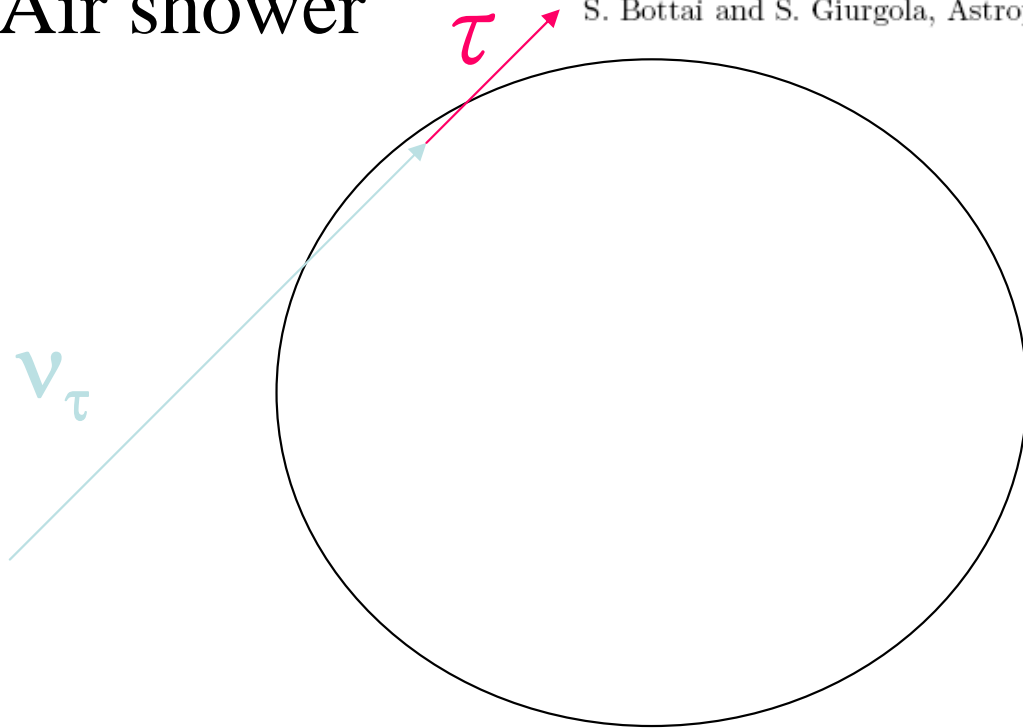
G. Domokos and S. Kovesi-Domokos, arXiv:hep-ph/9801362; arXiv:hep-ph/9805221. See also, D. Fargion, *Astrophys. J.* **570**, 909 (2002) [arXiv:astro-ph/0002453].

X. Bertou, P. Billoir, O. Deligny, C. Lachaud and A. Letessier-Selvon, *Astropart. Phys.* **17**, 183 (2002) [arXiv:astro-ph/0104452].

J. L. Feng, P. Fisher, F. Wilczek and T. M. Yu, *Phys. Rev. Lett.* **88**, 161102 (2002) [arXiv:hep-ph/0105067].

S. Bottai and S. Giurgola, *Astropart. Phys.* **18**, 539 (2003) [arXiv:astro-ph/0205325].

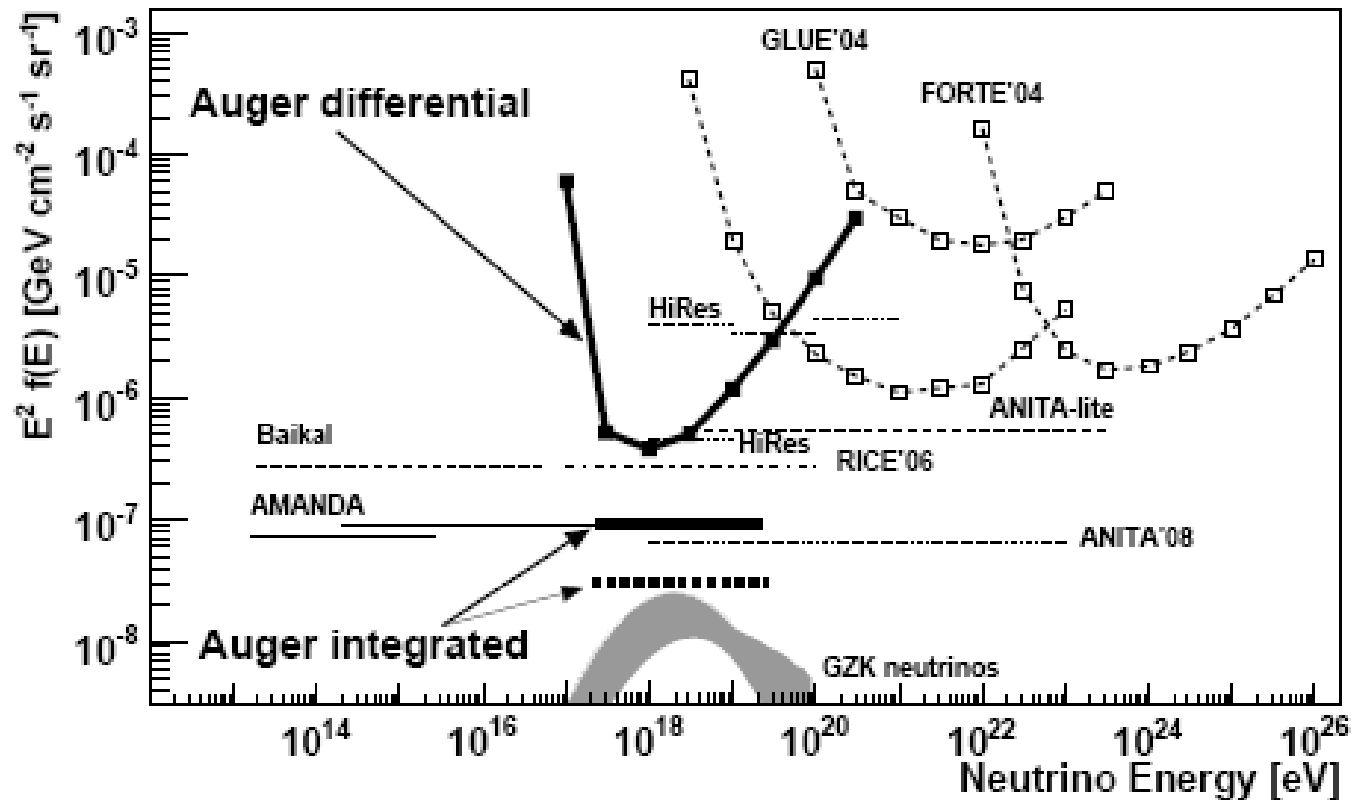
Air shower



- J.-J. Tseng, T.-W. Yeh, H. Athar, M. A. Huang, F.-F. Lee and GLL, *Phys. Rev. D* 2003
- GLL,

Earth-skimming tau neutrinos

Auger result on UHE tau neutrino flux



90% C.L.

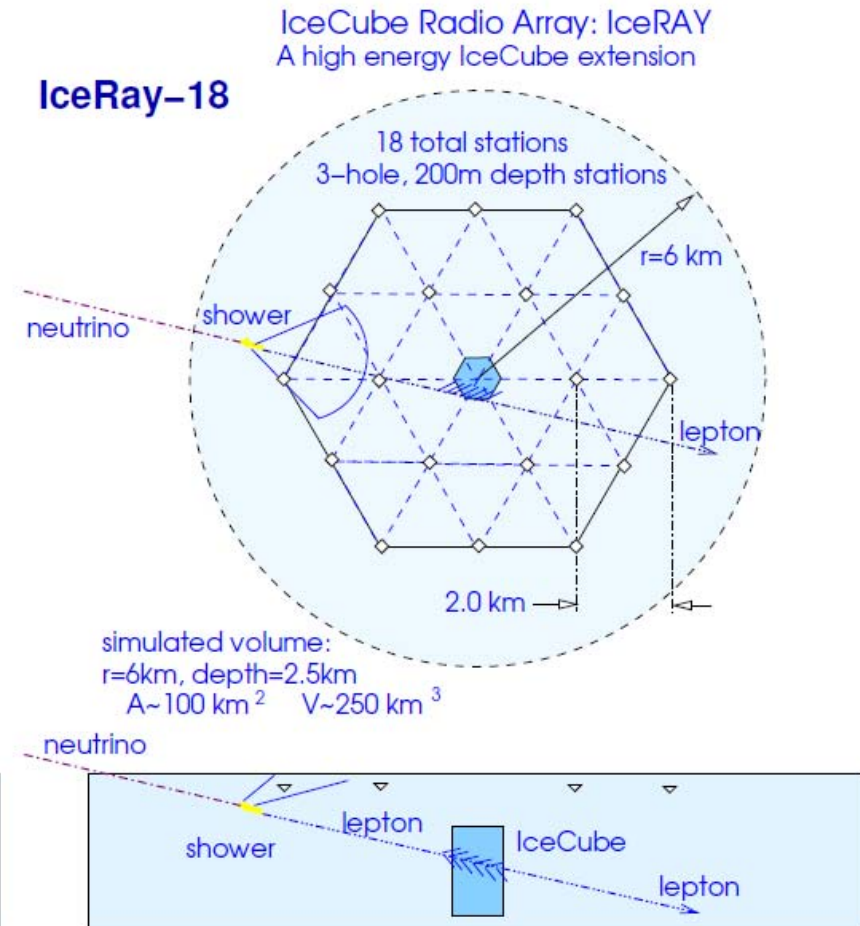
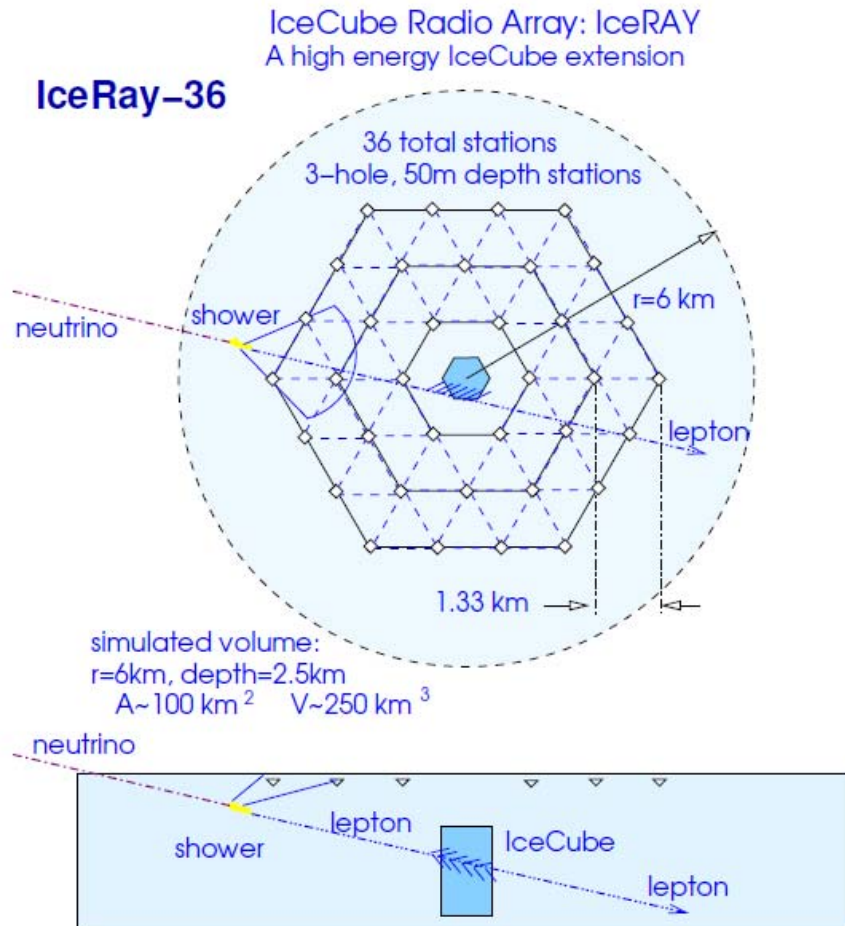
arXiv:0903.3385 [astro-ph.HE]

IceRay deployment-radio and water

Cherenkov technique put together

Higher density, shallow **50m**

sparse, deep **200m**



←→ **8km**

Summary—omitting details on how neutrino telescopes measure R and S

- The structure of the oscillation probability matrix P makes it difficult to constrain $\phi_0(\nu_\mu) - \phi_0(\nu_\tau)$
- 10% accuracy in terrestrial measurements on both R and S is required for distinguishing the pion source and muon-damped source—caused by the structure of P .
- Taking Waxman-Bahcall upper bound $E^2 \Phi_0 = 2 \times 10^{-8}$ GeV/cm² s sr, 10% accuracy ($O(100)$ events) takes IceCube detector more than a decade \rightarrow neutrino flavor astronomy challenging!
- Improved measurements on neutrino mixing parameters are not very helpful to the above situation. See P. 20
- To use the astrophysical source for probing the neutrino mixing parameter, it is unrealistic to assume a precise knowledge on the neutrino flavor ratio at the source.