

Electroweak phase transition in the light Higgs boson scenario

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in collaboration with Koichi Funakubo (Saga U, Japan)

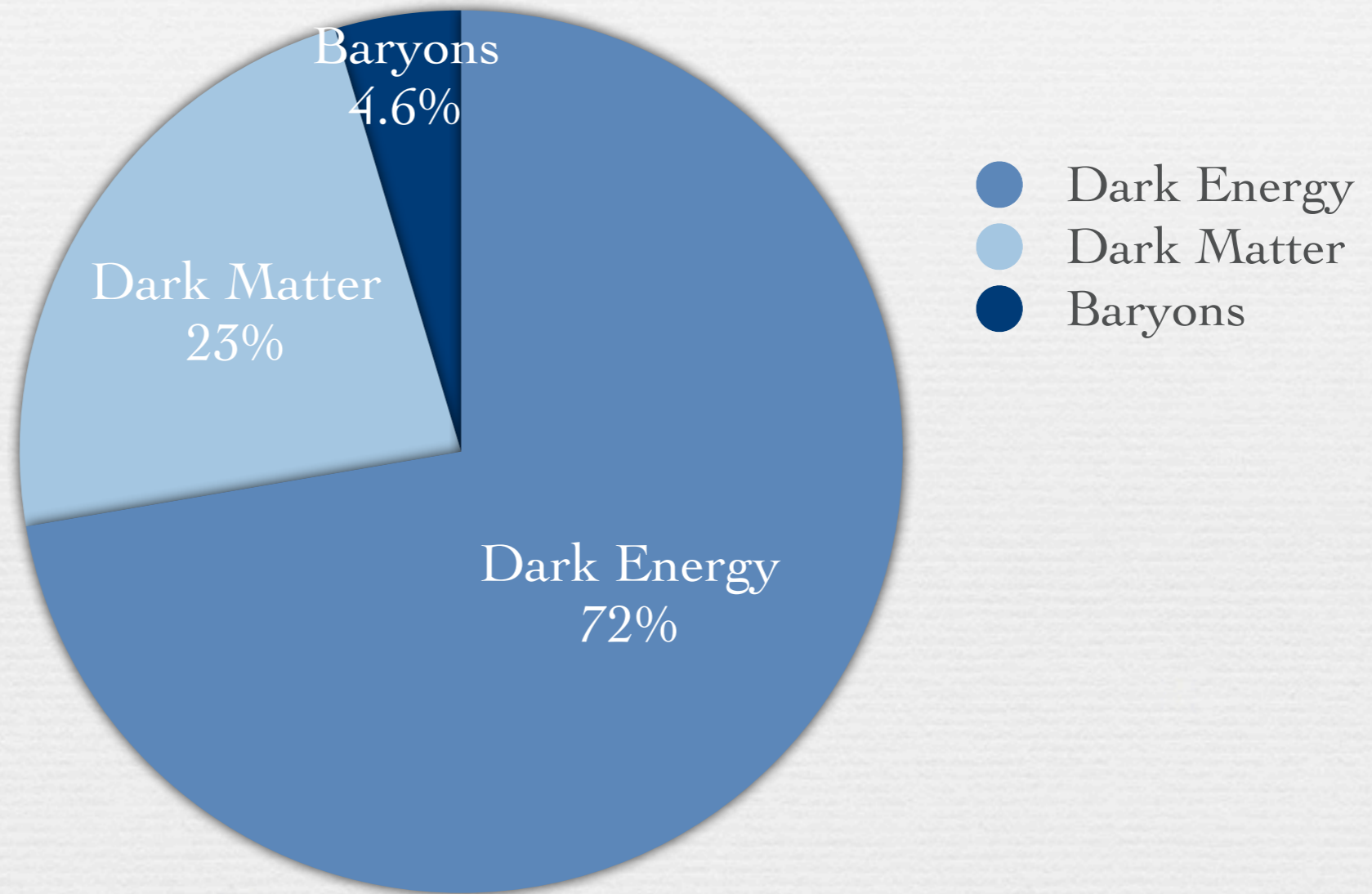
[Ref. arXiv:0905.2022](https://arxiv.org/abs/0905.2022)

Outline

- ❧ Motivation
Baryon Asymmetry of the Universe (BAU)
- ❧ Electroweak baryogenesis in the light Higgs boson scenario of the MSSM
 - Sphaleron decoupling condition
 - Electroweak phase transition (EWPT)
 - Critical bubble
- ❧ Summary

Motivation

- Energy budget of the Universe



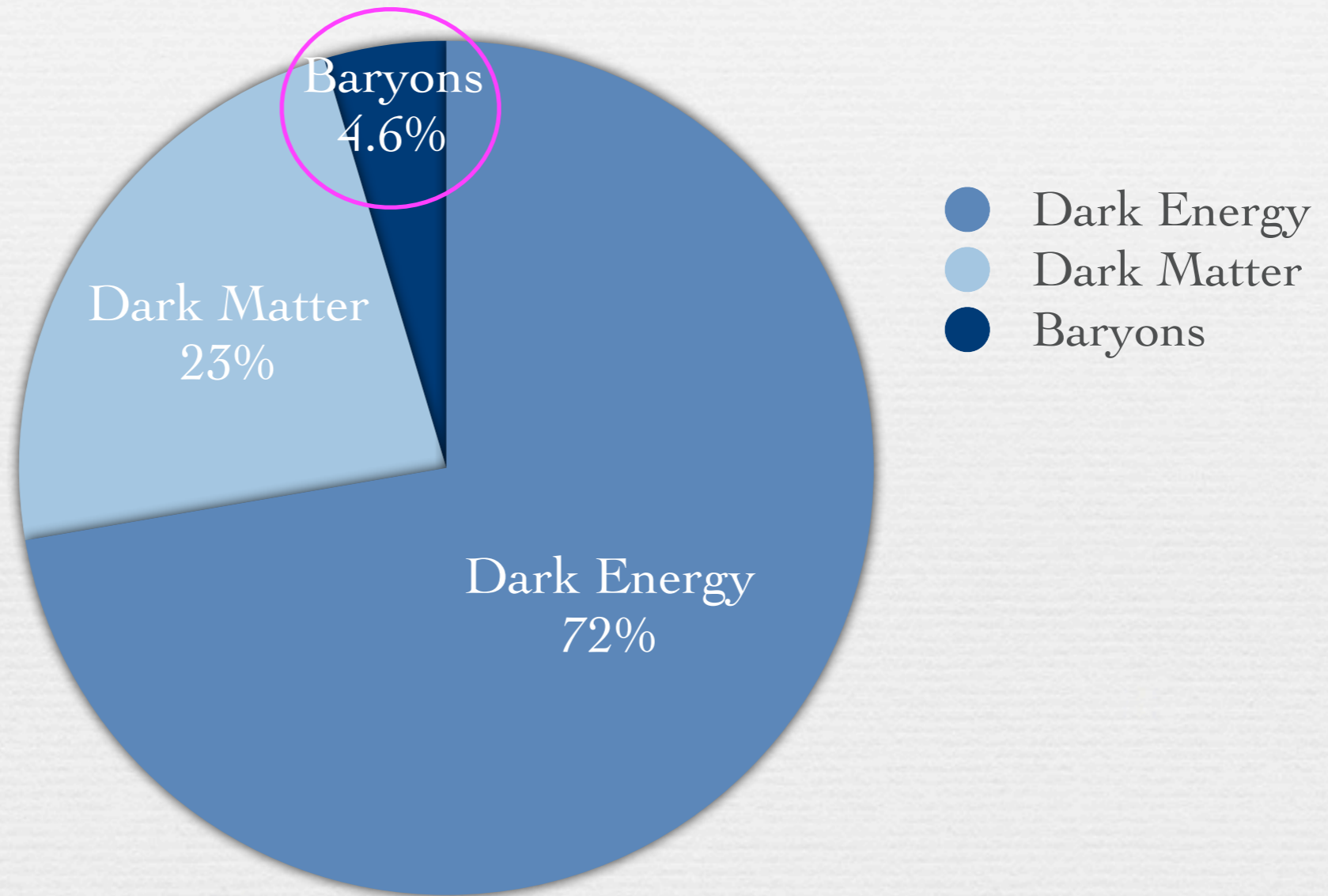
- 95% of the Universe is made of dark object.

- It should be stressed that there remains a mystery in the visible sector as well.

Where did antibaryons go?

Motivation

- Energy budget of the Universe



- 95% of the Universe is made of dark object.

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Where did antibaryons go?

BAU

• Baryon Asymmetry of the Universe (BAU)

$$\frac{n_B}{n_\gamma} = \frac{n_b - n_{\bar{b}}}{n_\gamma} = (4.7 - 6.5) \times 10^{-10} \quad (95\% \text{ C.L.}) \quad (\text{PDG 08})$$

Sakharov's conditions ('67)

(1) Baryon (B) violation

○ sphaleron process

(2) C violation

○ chiral gauge interaction

CP violation

✗ KM phase is not sufficient

(3) out of equilibrium

✗ phase transition is not 1st order

SM case

Extensions of the SM

MSSM, 2HDM, singlet-extended MSSM etc.

Today, we discuss the MSSM baryogenesis (BG).

Tension in the MSSM BG

- From the EWPT point of view, the light Higgs boson is generically favored.
- The LEP data put a strong constraint on the light Higgs boson.

On the other hand,

- In the literatures, $\frac{v_c}{T_c} \gtrsim 1$ is used as the practical criterion for the strong 1st order PT.

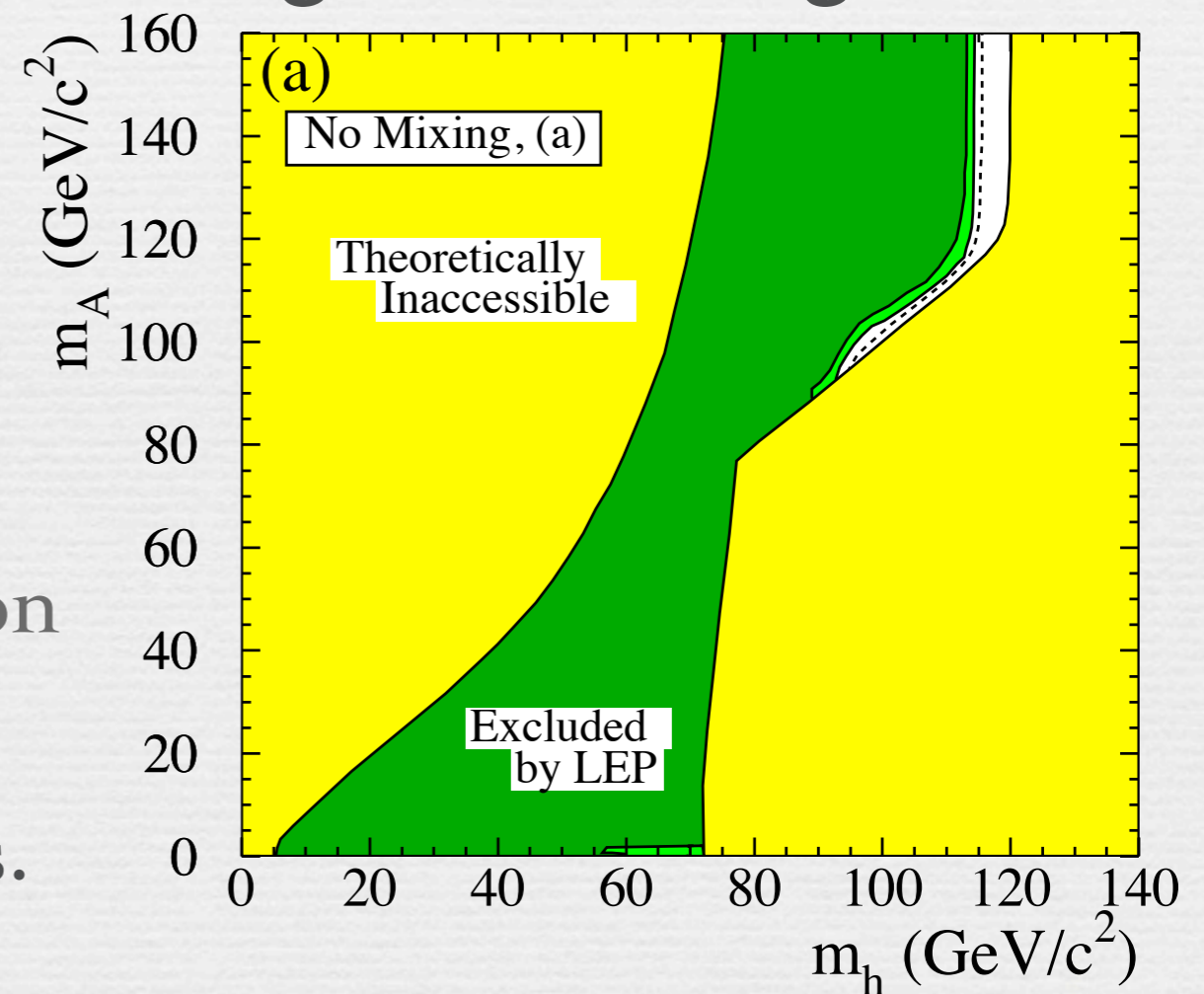
But,

- $v_c/T_c > 0.9$ region or $v_c/T_c > 1.1$ region are completely different.

⇒ Need to reduce the uncertainties.

- We analyze the EWPT in the light Higgs boson scenario ($m_h < 114.4$ GeV) based on the refined sphaleron decoupling cond.

e.g. no-mixing case



Effective potential

- To discuss the PT we use the effective potential.
- The gauge bosons and 3rd generation of quarks/squarks are taken into account.

$$V_{\text{eff}}(\Phi_d, \Phi_u) = V_0(\Phi_d, \Phi_u) + \Delta V(\Phi_d, \Phi_u; T),$$

Tree:
$$V_0(\Phi_d, \Phi_u) = m_1^2 \Phi_d^\dagger \Phi_d + m_2^2 \Phi_u^\dagger \Phi_u - (m_3^2 \epsilon_{ij} \Phi_d^i \Phi_u^j + \text{h.c.})$$
$$+ \frac{g_2^2 + g_1^2}{8} (\Phi_d^\dagger \Phi_d - \Phi_u^\dagger \Phi_u)^2 + \frac{g_2^2}{2} (\Phi_d^\dagger \Phi_u)(\Phi_u^\dagger \Phi_d),$$

1-loop:
$$\Delta V(\Phi_d, \Phi_u; T) = \sum_A c_A \left[F_0(\bar{m}_A^2) + \frac{T^4}{2\pi^2} I_{B,F} \left(\frac{\bar{m}_A^2}{T^2} \right) \right]$$

$$F_0(m^2) = \frac{m^4}{64\pi^2} \left(\ln \frac{m^2}{M^2} - \frac{3}{2} \right), \quad I_{B,F}(a^2) = \int_0^\infty dx x^2 \ln \left(1 \mp e^{-\sqrt{x^2+a^2}} \right)$$

Fitting function:

$$\tilde{I}_{B,F}(a^2) = e^{-a} \sum_{n=0}^N c_n^{b,f} a^n, \quad |\tilde{I}_{B,F} - I_{B,F}| < 10^{-6} \quad (N = 40).$$

- The fitting function is used in our numerical analysis.

Light Higgs boson scenario (LHS)

$$m_h < 114.4 \text{ GeV}$$

$$\text{LEP data } \xi = \frac{g_{\phi ZZ}}{g_{h^{\text{SM}} ZZ}}$$

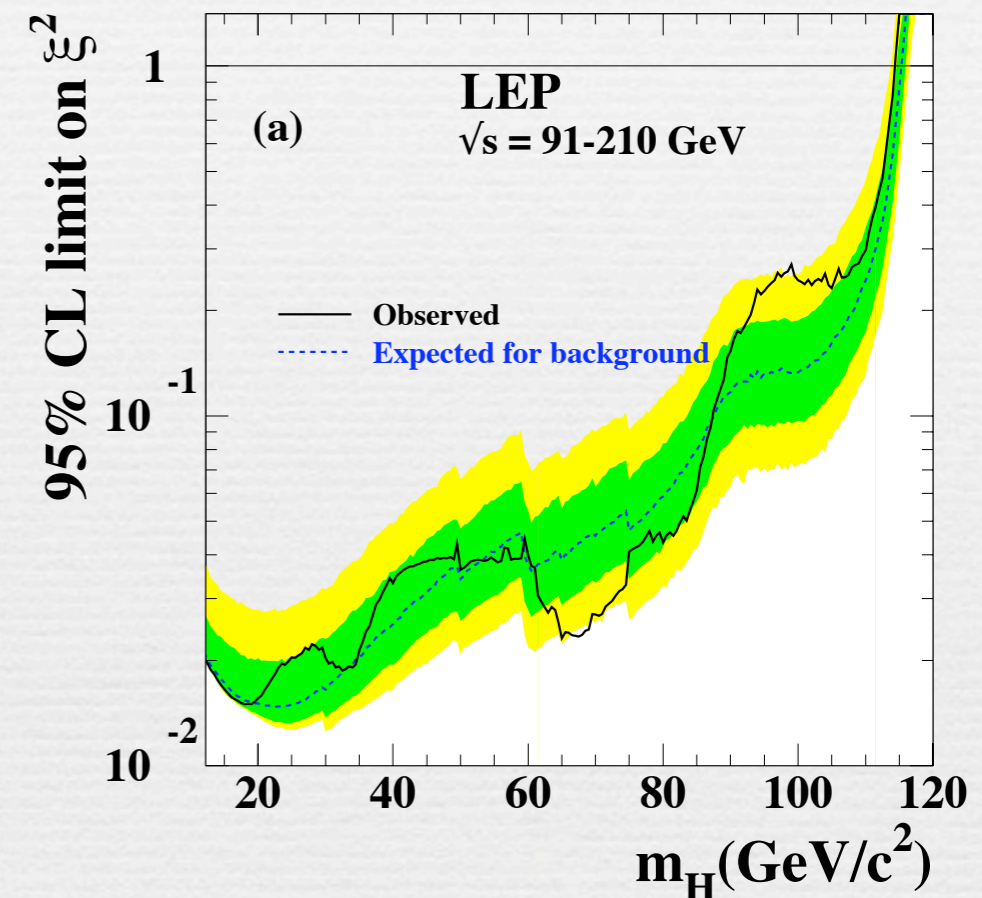
Higgs fields in the rotated basis:

$$(i\tau^2 \Phi'_d{}^* \ \Phi'_u)^T = O(\beta) (i\tau^2 \Phi_d{}^* \ \Phi_u)^T$$

$$\Phi'_d = \begin{pmatrix} \frac{1}{\sqrt{2}}(v_0 + h'_d + iG^0) \\ G^- \end{pmatrix}, \quad \Phi'_u = \begin{pmatrix} H^+ \\ \frac{1}{\sqrt{2}}(h'_u + iA) \end{pmatrix}$$

where $h'_d = h^{\text{SM}}$

$$\begin{pmatrix} h'_d \\ h'_u \end{pmatrix} = \begin{pmatrix} \cos(\beta - \alpha) & \sin(\beta - \alpha) \\ -\sin(\beta - \alpha) & \cos(\beta - \alpha) \end{pmatrix} \begin{pmatrix} H \\ h \end{pmatrix}$$



Higgs couplings:

$$\frac{g_{hZZ}}{g_{h^{\text{SM}} ZZ}} = \sin(\beta - \alpha), \quad \frac{g_{HZZ}}{g_{h^{\text{SM}} ZZ}} = \cos(\beta - \alpha),$$

$$\frac{g_{hZA}}{g_{h^{\text{SM}} ZG^0}} = \cos(\beta - \alpha), \quad \frac{g_{HZA}}{g_{h^{\text{SM}} ZG^0}} = \sin(\beta - \alpha).$$

LHS: $\cos(\beta - \alpha) \rightarrow 1$, $H \rightarrow h^{\text{SM}}$, $m_A \sim m_Z$

Decoupling lim: $\sin(\beta - \alpha) \rightarrow 1$, $h \rightarrow h^{\text{SM}}$, $m_A \gg m_Z$.

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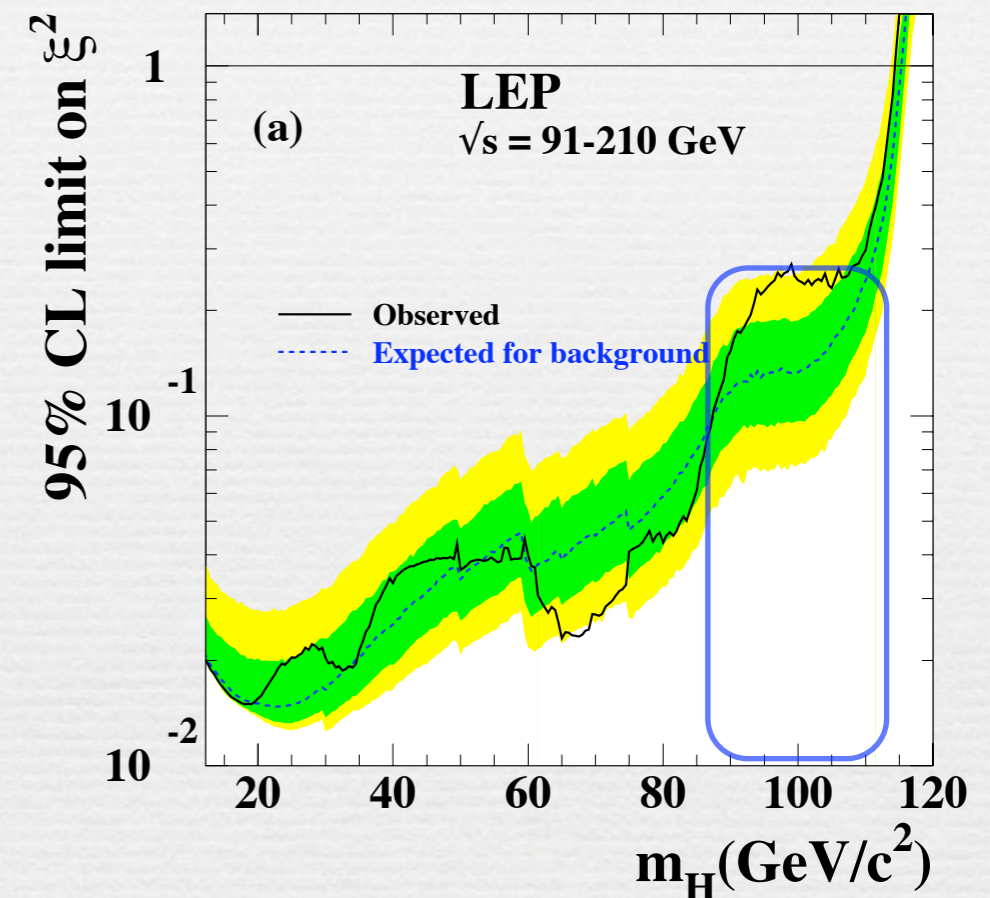
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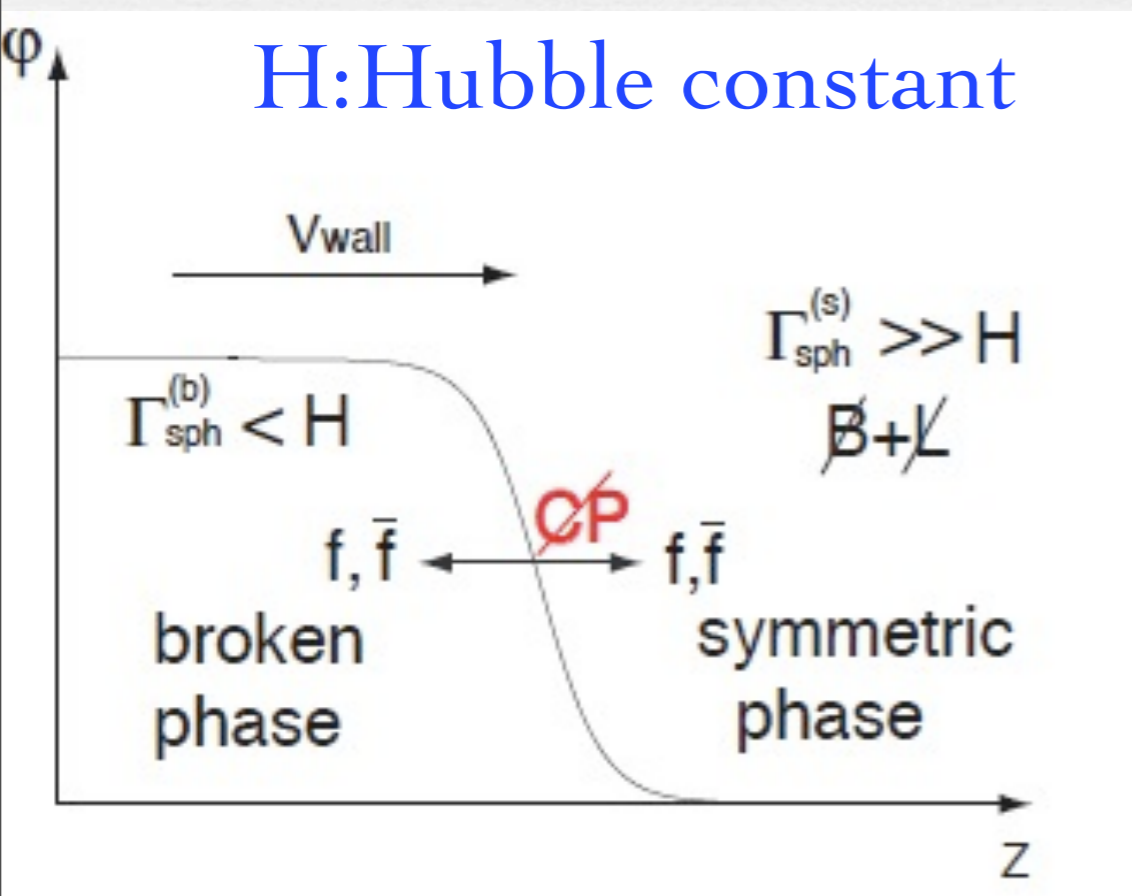
Decoupling lim: $\sin(\beta - \alpha) \rightarrow 1$, $h \rightarrow h^{\text{SM}}$, $m_A \gg m_Z$.

Mechanism of EWBG

[Kuzmin, Rubakov, Shaposhnikov, PLB155,36 ('85)]

- The EWPT must be a first order with expanding bubble walls.

outline



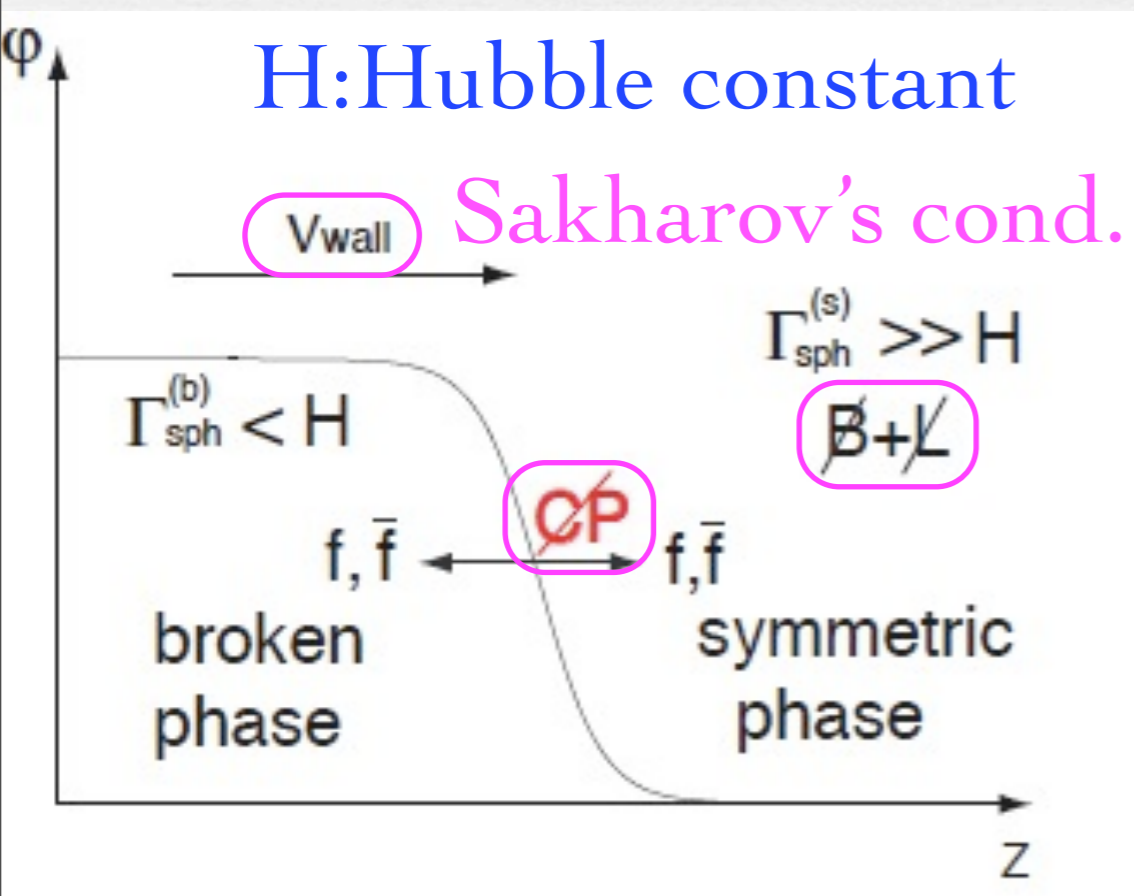
- Due to CP violation at the phase boundary, asymmetries of particle number densities occur.
- They diffuse into symmetric phase.
- Left-handed particle number densities are converted into B via sphaleron process.
- Sphaleron process is decoupled after the PT.
- B is frozen.

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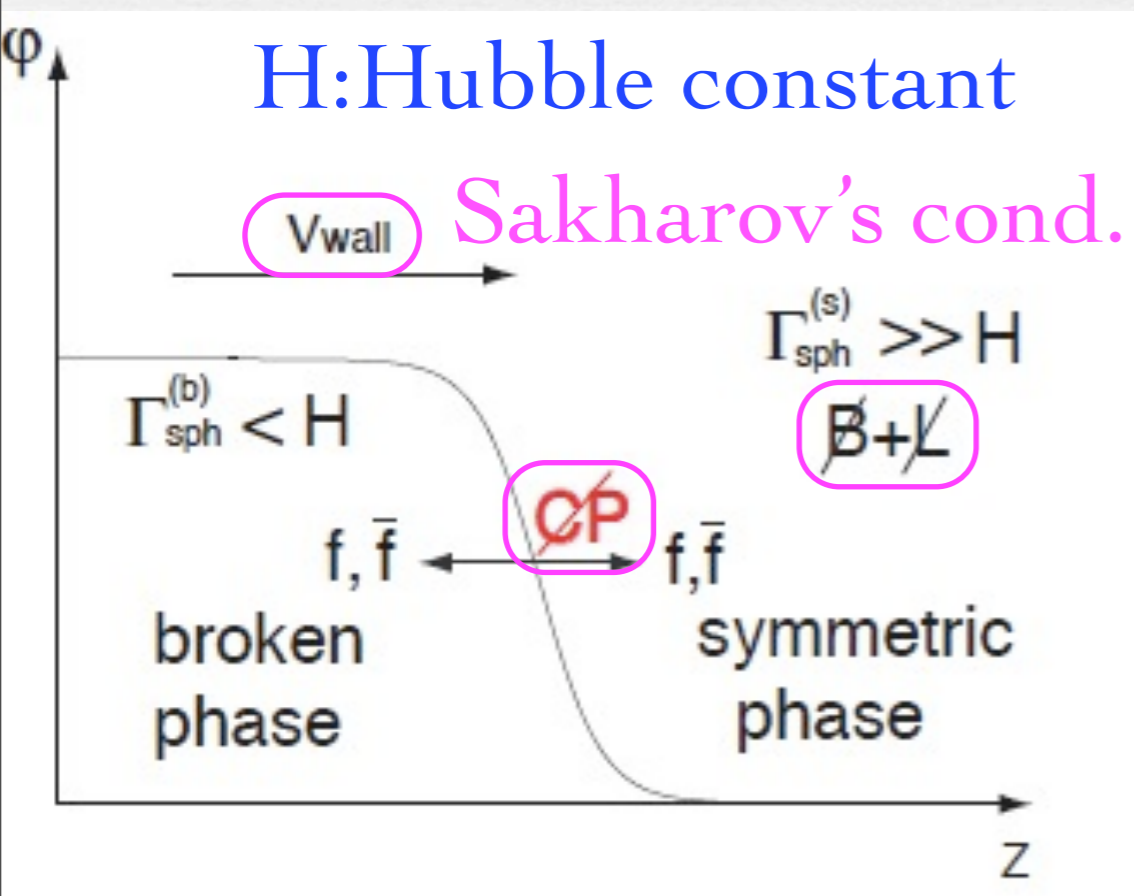
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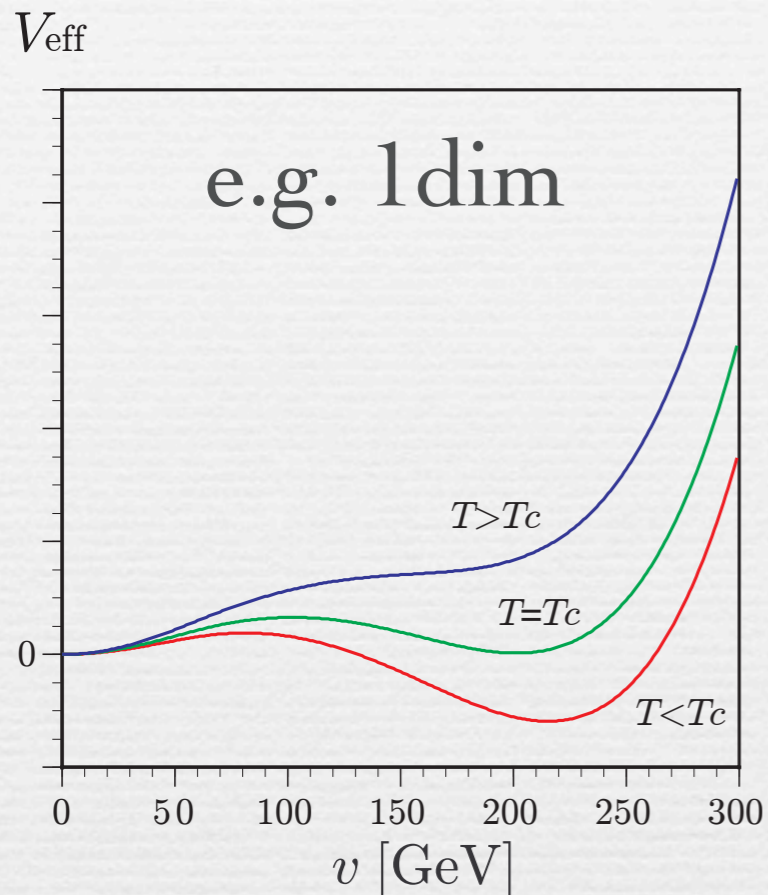
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Phase transition

- The PT must be 1st order to realize out of equilibrium.



- order parameters = Higgs VEVs
- At T_c , the potential has two degenerate minima.

- Light stop plays a crucial role in strengthening the 1st order PT.

[Carena, Quiros, Wagner, PLB380 ('96) 81]

To be consistent with the LEP bound on m_H and ρ -parameter, we should take

$$m_{\tilde{q}}^2 \gg m_{\tilde{t}_R}^2, X_t^2, \quad X_t = A_t - \mu \cot \beta$$

$$\text{lighter stop mass} \Rightarrow \bar{m}_{\tilde{t}_1}^2 = m_{\tilde{t}_R}^2 + \frac{y_t^2 \sin^2 \beta}{2} \left(1 - \frac{X_t^2}{m_{\tilde{q}}^2} \right) v^2 + \mathcal{O}(g^2)$$

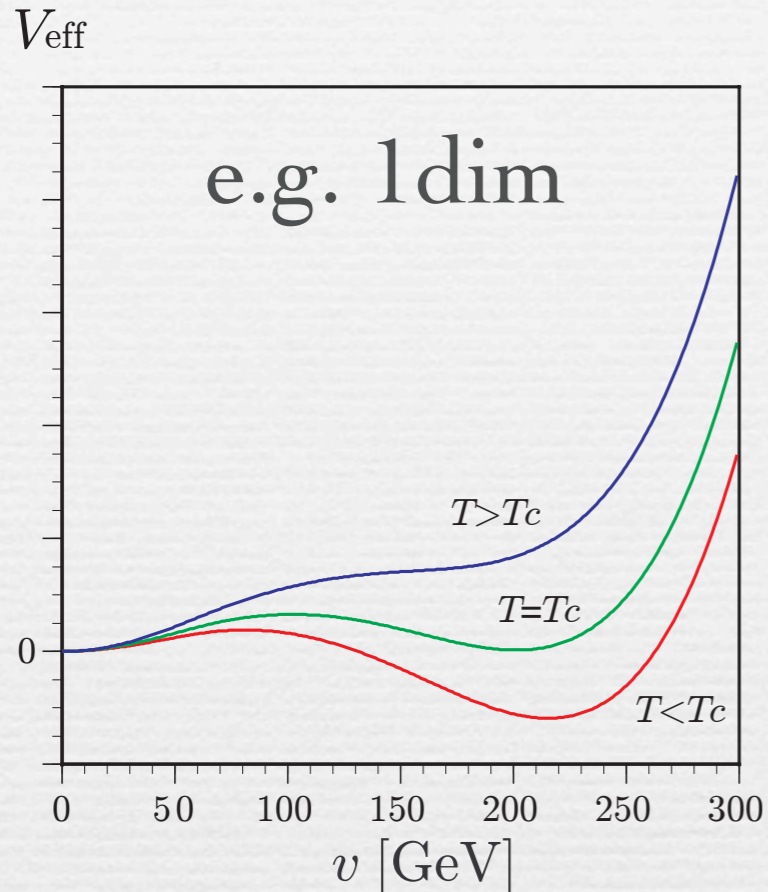
c.f. High T expansion:

$$V_{\text{eff}} \ni - (E_{\text{SM}} + E_{\tilde{t}_1}) T v^3 \quad E_{\tilde{t}_1} \simeq \frac{y_t^3 \sin^3 \beta}{4\sqrt{2}\pi} \left(1 - \frac{X_t^2}{m_{\tilde{q}}^2} \right)^{3/2} \sim 0.06 \text{ (} X_t=0 \text{)}$$

~ 0.01
large effect!

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$\rightarrow 0 \text{ (} m_{\tilde{t}_1} < m_t \text{)}$

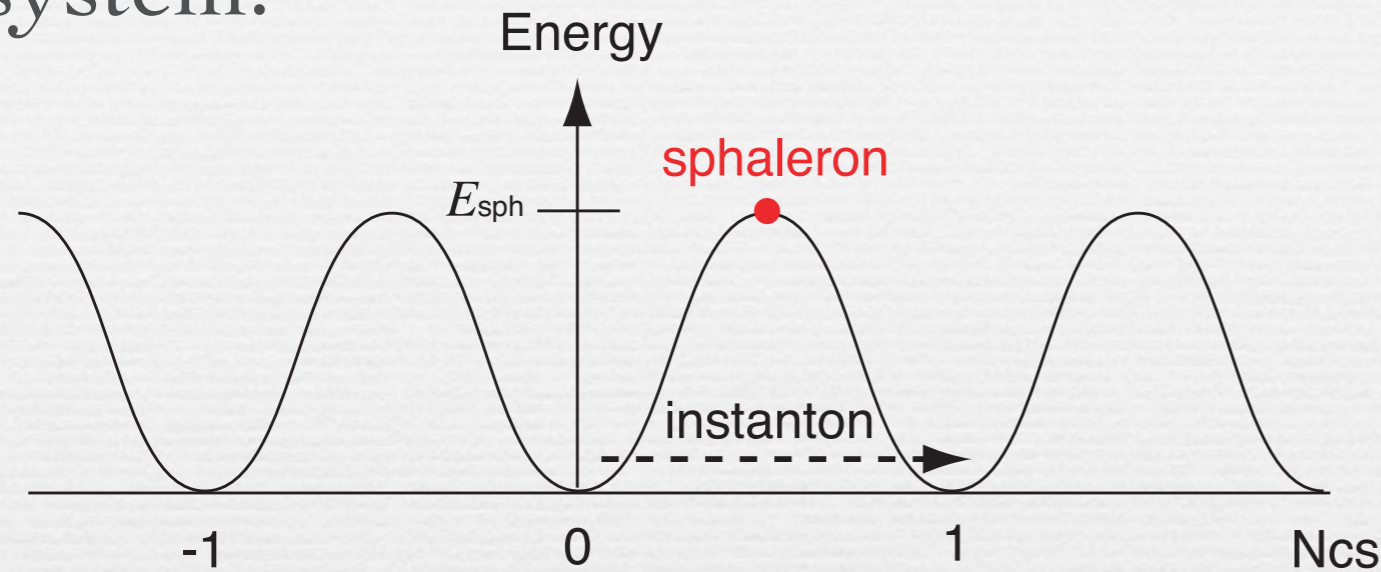
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Sphaleron

- A static saddle point solution w/ finite energy of the gauge-Higgs system. [N.S. Manton, PRD28 ('83) 2019]



$$\Delta B \neq 0$$

Instanton: quantum tunneling

Sphaleron: thermal fluctuation

B violation:

$$\Delta B = 3\Delta N_{CS} \quad N_{CS} = \frac{g_2^2}{32\pi^2} \int d^3x \epsilon_{ijk} \text{Tr} \left[F_{ij} A_k - \frac{2}{3} g_2 A_i A_j A_k \right]$$

Vacuum transition rates:

In the symmetric phase : $\sim \kappa(\alpha_W T)^4$, $\alpha_W = g_2^2/(4\pi)$

In the broken phase : $\sim T^3 e^{-E_{sph}/T}$.

At $T = 0$: $\sim e^{-S_{\text{instanton}}} = e^{-8\pi^2/g_2^2} \ll 1$ **no proton decay**

- **B violating process is active at finite T but is suppressed at $T=0$.**

Sphaleron decoupling condition

- To avoid the washout of the generated BAU, the sphaleron process must be decoupled after the PT.

Hubble constant

$$\boxed{\frac{1}{B} \frac{dB}{dt}} \simeq \frac{13 \cdot 3}{4 \cdot 32\pi^2} \frac{\omega_-}{\alpha_W^3} \kappa \mathcal{N}_{\text{tr}} \mathcal{N}_{\text{rot}} e^{-E_{\text{sph}}/T} < \boxed{H(T)} \simeq 1.66 \sqrt{g_*} T^2 / m_{\text{P}}$$

E_{sph} : sphaleron energy, $\mathcal{N}_{\text{tr}}, \mathcal{N}_{\text{rot}}$: zero mode factors

ω_- : negative mode, $\kappa = \mathcal{O}(1)$.

If we denote $E_{\text{sph}} = 4\pi v \mathcal{E} / g_2$

$$\frac{v}{T} > \frac{g_2}{4\pi \mathcal{E}} \left[42.97 + \ln(\kappa \mathcal{N}_{\text{tr}} \mathcal{N}_{\text{rot}}) + \ln \left(\frac{\omega_-}{m_W} \right) - \frac{1}{2} \ln \left(\frac{g_*}{106.75} \right) - 2 \ln \left(\frac{T}{100 \text{ GeV}} \right) \right]$$

In the SM:

$$\mathcal{E} = 2.00, \quad \mathcal{N}_{\text{tr}} \mathcal{N}_{\text{rot}} = 80.13, \quad \omega_-^2 = 2.3 m_W^2, \quad \kappa = 1, \quad T = 100 \text{ GeV}, \quad \lambda = g_2^2.$$

$$\frac{v}{T} > 0.026 \times (42.97 + \underbrace{4.38}_{10\% \text{ correction}} + 0.416) = 1.24$$

MSSM case

■ Effects of T and zero modes

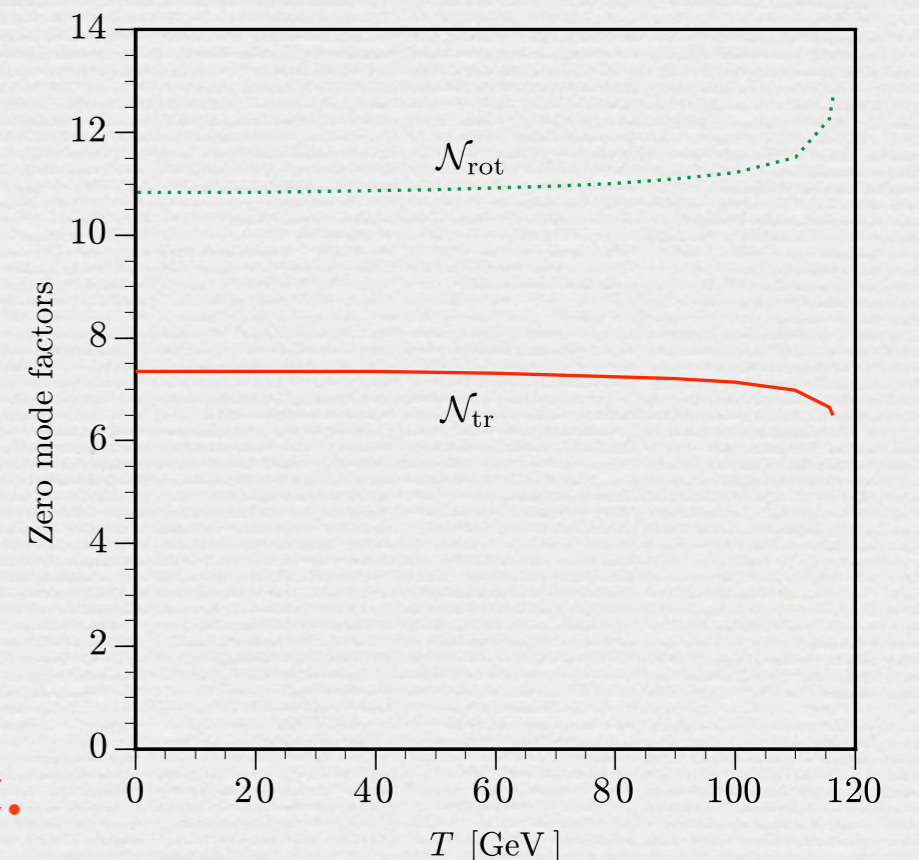
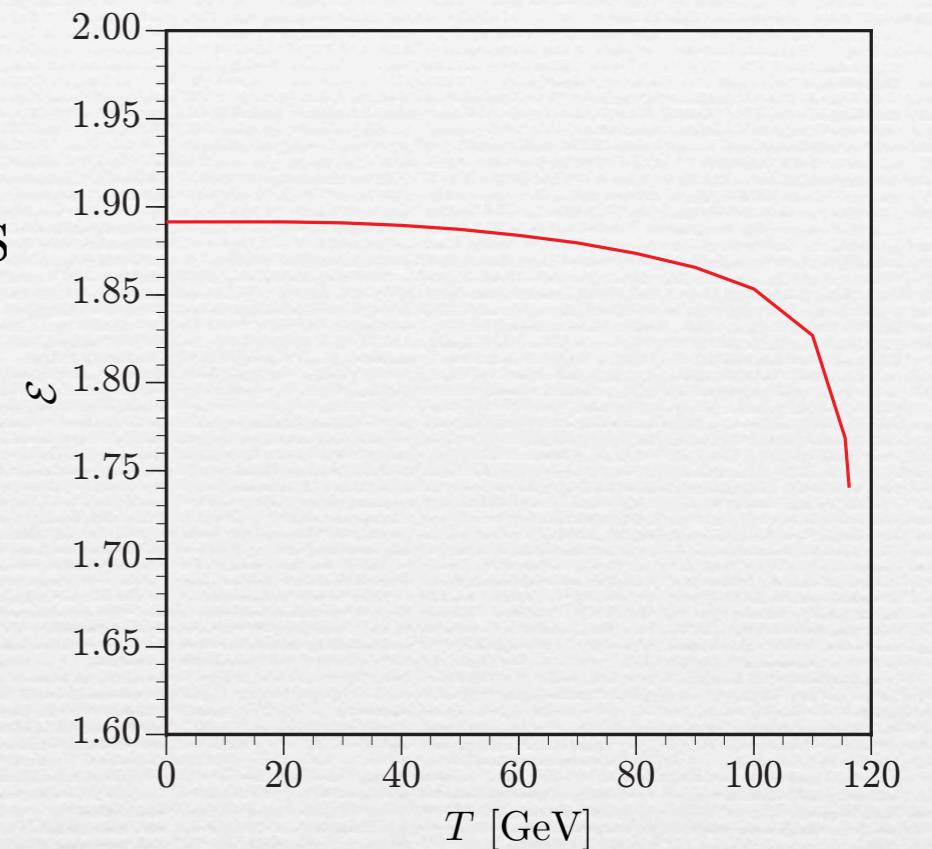
I: based on $V_{\text{eff}}(T = 0)$ *without* the zero modes

II: based on $V_{\text{eff}}(T = 0)$ *with* the zero modes

III: based on $V_{\text{eff}}(T \neq 0)$ *with* the zero modes

For the typical parameter set

	I	II	III
\mathcal{E}	1.89	1.89	1.77
\mathcal{N}_{tr}	—	7.36	6.65
\mathcal{N}_{rot}	—	10.84	12.27
$v_N/T_N >$	1.17	1.29	1.38



- Zero mode factors cannot be neglected.
- T -dependence must be taken into account.

Experimental constraints

- Higgs bounds@ [PLB565, 61 (2003)]

$$g_{H_i Z Z}^2 \times \text{Br}(H_i \rightarrow f \bar{f}) < \mathcal{F}_{H_i Z}(m_{H_i}),$$

$$g_{H_i H_j Z}^2 \times \text{Br}(H_i \rightarrow f \bar{f}) \times \text{Br}(H_j \rightarrow f \bar{f}) < \mathcal{F}_{H_i H_j}(m_{H_i} + m_{H_j}),$$

where $f = b, \tau$. $\mathcal{F}_{H_i Z}$ and $\mathcal{F}_{H_i H_j}$ are the 95% C.L. upper limits

- Lower bounds for SUSY particles:

e.g. chargino mass > 94 GeV

- ρ -parameter:

$$\Delta\rho \equiv \frac{\Pi_{ZZ}^T(0)}{m_Z^2} - \frac{\Pi_{WW}^T(0)}{m_W^2} < 0.002$$

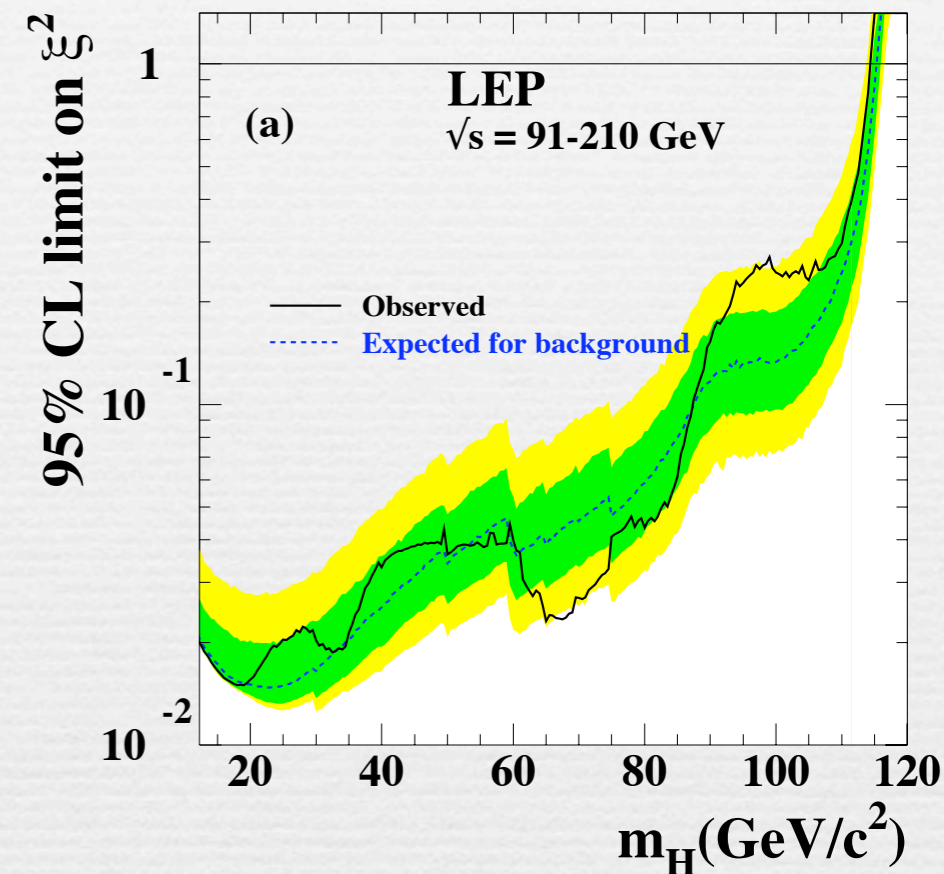
- B physics observables:

$$\text{Br}(B_u \rightarrow \tau \nu_\tau)_{\text{exp}} = 1.41_{-0.42}^{+0.43} \times 10^{-4},$$

$$\text{Br}(\bar{B} \rightarrow X_s \gamma)_{\text{exp}} = (3.52 \pm 0.23 \pm 0.09) \times 10^{-4},$$

$$\text{Br}(B_s \rightarrow \mu^+ \mu^-)_{\text{exp}} < 0.23 \times 10^{-7}.$$

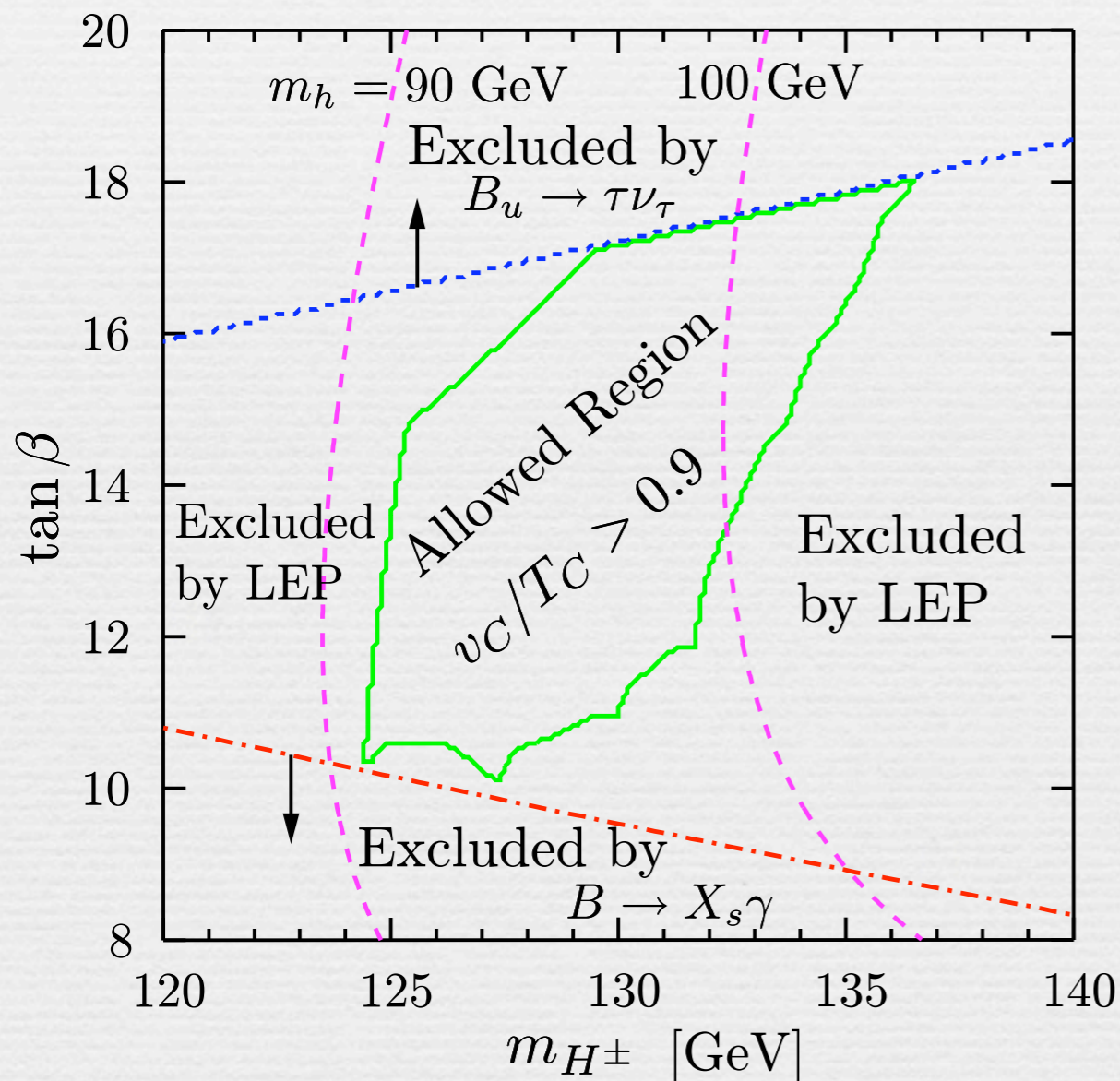
e.g. Higgsstrahlung



Allowed region

- The allowed region is highly constrained by the experimental data.

$$m_{\tilde{q}} = 1200 \text{ GeV}, m_{\tilde{t}_R} \simeq 0, A_t = A_b = -300 \text{ GeV}.$$



Maximal v/T :

$$\tan \beta = 10.1, m_{H^\pm} = 127.4 \text{ GeV}$$

$$\frac{v_C}{T_C} = \frac{107.10 \text{ GeV}}{116.27 \text{ GeV}} = 0.92$$

The sphaleron process is not decoupled at T_c .

Loophole: **supercooling**

\Rightarrow The PT begins to proceed with bubble wall at below T_c .

We need to know the critical bubbles.

Critical bubble

- For the EWPT to proceed, the radius of bubble must be larger than some critical size.

Higgs fields:
$$\Phi_d = \frac{1}{\sqrt{2}} \begin{pmatrix} \rho_d \\ 0 \end{pmatrix}, \quad \Phi_u = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ \rho_u \end{pmatrix},$$

Energy functional:

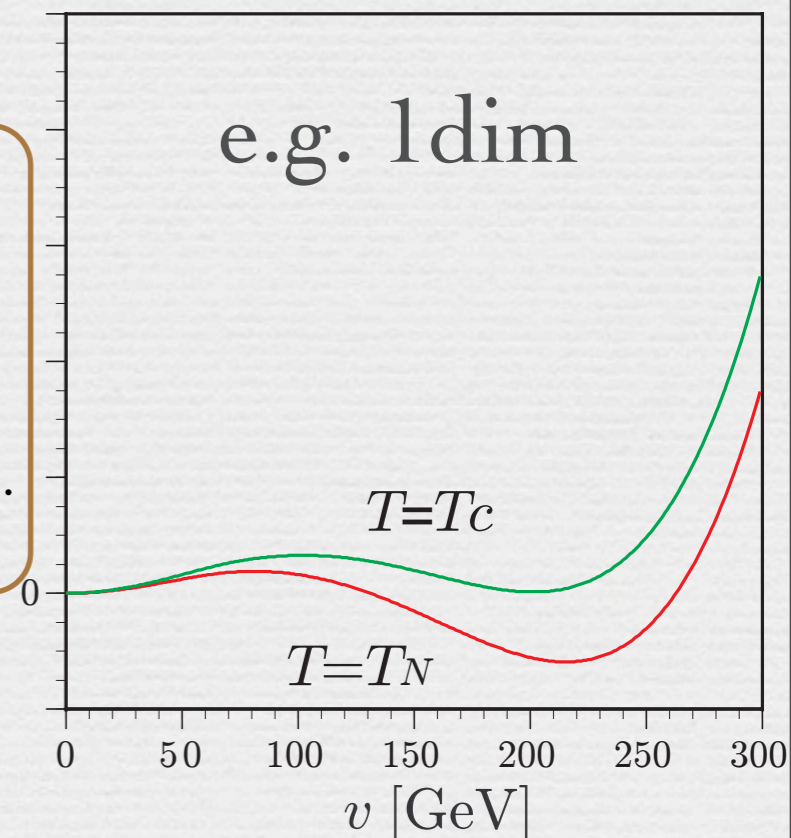
$$E = 4\pi \int_0^\infty dr r^2 \left[\frac{1}{2} \left\{ \left(\frac{d\rho_d}{dr} \right)^2 + \left(\frac{d\rho_u}{dr} \right)^2 \right\} + V_{\text{eff}}(\rho_d, \rho_u; T) \right] \quad r = \sqrt{x^2}$$

Equation of motion (EOM):

$$\begin{aligned} -\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{d\rho_d}{dr} \right) + \frac{\partial V_{\text{eff}}}{\partial \rho_d} &= 0, & \lim_{r \rightarrow \infty} \rho_d(r) &= 0, & \lim_{r \rightarrow \infty} \rho_u(r) &= 0, \\ -\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{d\rho_u}{dr} \right) + \frac{\partial V_{\text{eff}}}{\partial \rho_u} &= 0. & \left. \frac{d\rho_d(r)}{dr} \right|_{r=0} &= 0, & \left. \frac{d\rho_u(r)}{dr} \right|_{r=0} &= 0. \end{aligned}$$

b.c.

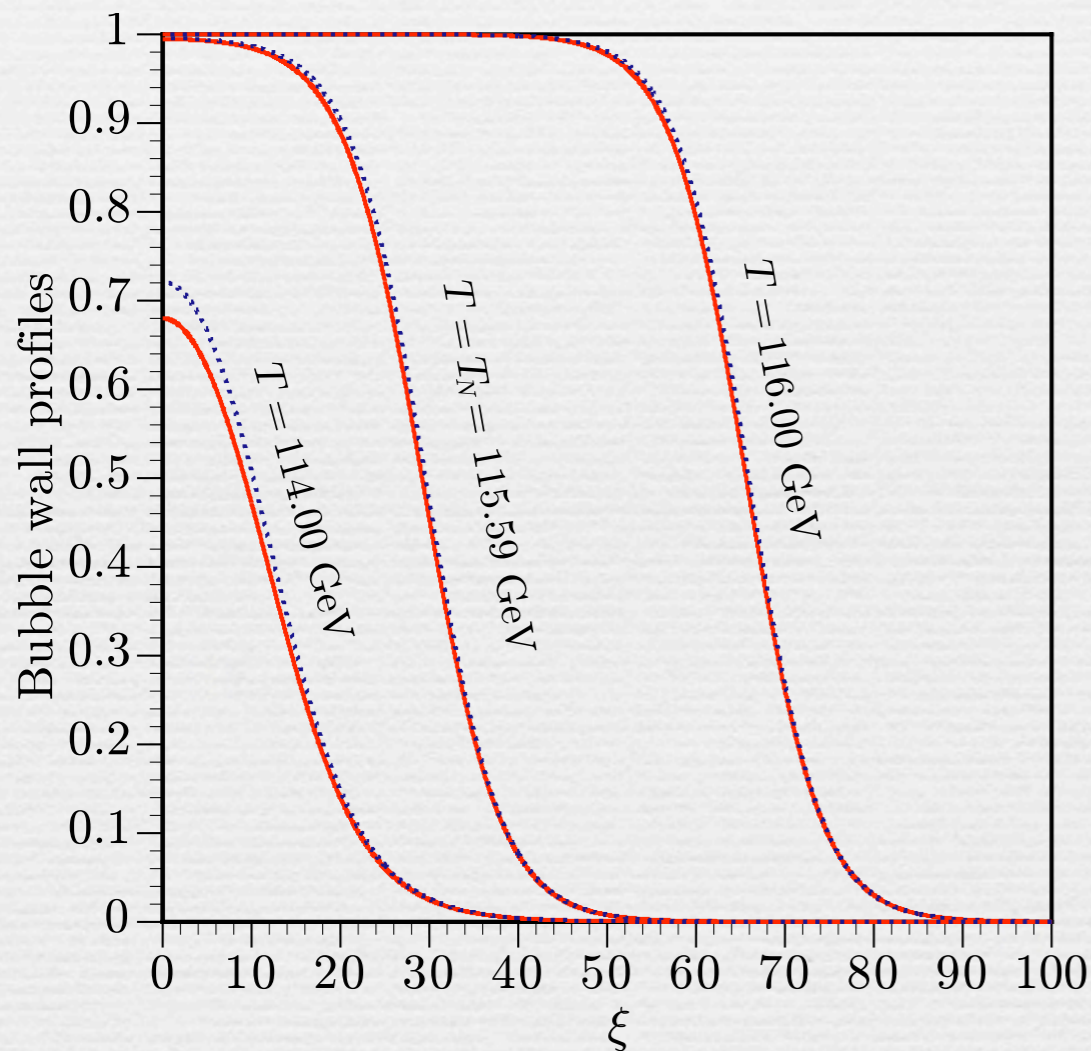
- The EOMs are numerically solved by relaxation methods.



Bubble nucleation

Nucleation rate: $\Gamma_N(T) \simeq T^4 \left(\frac{E_{\text{cb}}(T)}{2\pi T} \right)^{3/2} e^{-E_{\text{cb}}(T)/T}$ [A.D. Linde, NPB216 ('82) 421]

Nucleation T : $\Gamma_N(T_N) H(T_N)^{-3} = H(T_N)$



$$\xi = vr, \quad h_1(\xi) = \frac{\rho_d(r)}{v \cos \beta}, \quad h_2(\xi) = \frac{\rho_u(r)}{v \sin \beta}$$

Numerical results:

$$\frac{v_N}{T_N} = \frac{116.73}{115.59} = 1.01$$

10% enhancement! But,

Sphaleron decoupling cond. @ T_N :

$$\mathcal{E} = 1.77, \quad \mathcal{N}_{\text{tr}} = 6.65, \quad \mathcal{N}_{\text{rot}} = 12.27$$

$$\frac{v_N}{T_N} > 1.38$$

- The sphaleron process is not decoupled at T_N either.

More examples

$$A_b = A_t = -300 \text{ GeV}, m_{\tilde{t}_R} = 10^{-4} \text{ GeV}, m_{\tilde{b}_R} = 1000 \text{ GeV},$$

$$\mu = 100 \text{ GeV}, M_2 = 500 \text{ GeV},$$

$m_{\tilde{q}} \text{ (GeV)}$	1200	1300	1400	1500
$\tan \beta$	10.11	9.87	9.75	9.57
$m_{H^\pm} \text{ (GeV)}$	127.40	127.40	127.50	127.50
v_C/T_C	$\frac{107.096}{116.274} = 0.921$	$\frac{107.512}{116.496} = 0.923$	$\frac{107.769}{116.770} = 0.923$	$\frac{107.915}{117.045} = 0.922$
$\tan \beta_C$	13.803	13.640	13.597	13.455
v_N/T_N	$\frac{116.727}{115.585} = 1.010$	$\frac{117.155}{115.798} = 1.012$	$\frac{117.404}{116.067} = 1.012$	$\frac{117.531}{116.339} = 1.010$
$\tan \beta_N$	13.676	13.503	13.453	13.307
E_{cb}/T_N	150.386	150.379	150.370	150.360
\mathcal{E}	1.769	1.770	1.770	1.771
\mathcal{N}_{tr}	6.652	6.658	6.662	6.667
\mathcal{N}_{rot}	12.266	12.253	12.240	12.229
$v_N/T_N >$	1.383	1.382	1.382	1.380

Typically, $v_N/T_N > 1.38$ is needed for sphaleron decoupling.

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$$A_b = A_t = -300 \text{ GeV}, m_{\tilde{t}_R} = 10^{-4} \text{ GeV}, m_{\tilde{b}_R} = 1000 \text{ GeV},$$

$$\mu = 100 \text{ GeV}, M_2 = 500 \text{ GeV},$$

$m_{\tilde{q}}$ (GeV)	1200	1300	1400	1500
$\tan \beta$	10.11	9.87	9.75	9.57
m_{H^\pm} (GeV)	127.40	127.40	127.50	127.50
v_C/T_C	$\frac{107.096}{116.274} = 0.921$	$\frac{107.512}{116.496} = 0.923$	$\frac{107.769}{116.770} = 0.923$	$\frac{107.915}{117.045} = 0.922$
$\tan \beta_C$	13.803	13.640	13.597	13.455
v_N/T_N	$\frac{116.727}{115.585} = 1.010$	$\frac{117.155}{115.798} = 1.012$	$\frac{117.404}{116.067} = 1.012$	$\frac{117.531}{116.339} = 1.010$
$\tan \beta_N$	13.676	13.503	13.453	13.307
E_{cb}/T_N	150.386	150.379	150.370	150.360
\mathcal{E}	1.769	1.770	1.770	1.771
\mathcal{N}_{tr}	6.652	6.658	6.662	6.667
\mathcal{N}_{rot}	12.266	12.253	12.240	12.229
$v_N/T_N >$	1.383	1.382	1.382	1.380

Typically, $v_N/T_N > 1.38$ is needed for sphaleron decoupling.

More examples

$$A_b = A_t = -300 \text{ GeV}, m_{\tilde{t}_R} = 10^{-4} \text{ GeV}, m_{\tilde{b}_R} = 1000 \text{ GeV},$$

$$\mu = 100 \text{ GeV}, M_2 = 500 \text{ GeV},$$

$m_{\tilde{q}}$ (GeV)	1200	1300	1400	1500
$\tan \beta$	10.11	9.87	9.75	9.57
m_{H^\pm} (GeV)	127.40	127.40	127.50	127.50
v_C/T_C	$\frac{107.096}{116.274} = 0.921$	$\frac{107.512}{116.496} = 0.923$	$\frac{107.769}{116.770} = 0.923$	$\frac{107.915}{117.045} = 0.922$
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Typically, $v_N/T_N > 1.38$ is needed for sphaleron decoupling.

Loopholes

Our negative results might be circumvented.

- $T_N \Rightarrow$ **onset** of the PT. We should know a temperature at which the PT **ends**. The sphaleron decoupling condition should be imposed at such a temperature.
- Higher order (2-loop) contributions must be taken into account. [J.R. Espinosa, NPB475, ('06) 273]

\Rightarrow **The sphaleron decoupling cond. might be relaxed.**

- The potential can be extended in such a way that stop also has a nontrivial VEV. (Charge-Color-Breaking vacuum)

\Rightarrow **MSSM BG is viable**. [Canena et al, NPB812, ('09) 243]

[N.B.] EW vacuum: metastable, CCB vacuum: global minimum

If the refined sphaleron decoupling cond. is used, is it still viable?

Summary

- We have analyzed the strength of the 1st order EWPT in the light Higgs boson scenario of the MSSM.
- v/T at T_N can be enhanced by about 10% compared to that at T_c .
- The sphaleron decoupling condition at T_N is typically given by $v/T > 1.38$.
- The sphaleron process is not decoupled at both T_c and T_N .

Backup

Decoupling limit

$$|A_t| = |A_b| = |\mu|/\tan\beta, \quad m_{\tilde{t}_R} = 10^{-4} \text{ GeV}, \quad m_{\tilde{b}_R} = 1000 \text{ GeV},$$

$$|\mu| = 100 \text{ GeV}, \quad M_1 = 100 \text{ GeV}, \quad M_2 = 500 \text{ GeV}.$$

$m_{\tilde{q}}$ (GeV)	1700	1800	1900	2000
$\tan\beta$	42.62	15.10	10.97	9.35
m_{H^\pm} (GeV)	1000.00	1000.00	1000.00	1000.00
v_C/T_C	$\frac{111.461}{116.993} = 0.953$	$\frac{111.460}{117.007} = 0.953$	$\frac{111.483}{116.994} = 0.953$	$\frac{111.440}{117.060} = 0.952$
$\tan\beta_C$	42.966	15.171	11.022	9.394
v_N/T_N	$\frac{121.454}{116.221} = 1.045$	$\frac{121.452}{116.236} = 1.045$	$\frac{121.478}{116.222} = 1.045$	$\frac{121.424}{116.288} = 1.044$
$\tan\beta_N$	42.955	15.168	11.019	9.392
$E_{\text{cb}}(T_N)/T_N$	150.366	150.370	150.364	150.360
\mathcal{E}	1.773	1.773	1.773	1.773
\mathcal{N}_{tr}	6.677	6.677	6.678	6.678
\mathcal{N}_{rot}	12.211	12.210	12.210	12.209
$v_N/T_N >$	1.379	1.379	1.379	1.379

The sphaleron process is not decoupled in this case either.