

**PPP8**

**RARE ELECTROWEAK DECAYS  
OF  $K$  AND  $B$  MESONS**

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## Introduction

- ☉ Standard Model (SM) successful.
- ☉ Calculation for strong QCD processes not well established.
- ☉ Electroweak processes simple at the quark level.
- ☉ Calculations complicated for hadronic states.
- ☉ Quarks are tightly bound inside hadrons.
- ☉ Non-perturbative nature of QCD forces.
- ☉ Interactions of bound quarks in hadrons not well understood.

## Electroweak rare decay processes

$$M_1 \rightarrow M_2 \ell \bar{\ell} \quad \ell = \nu, e^-$$

### Quark Level

$$q_1 \rightarrow q_2 \ell \bar{\ell} \quad \ell = \nu, e^-$$

**Neutrino Pair Emission**  $q_1 \rightarrow q_2 \nu \bar{\nu}$ 

$$\Gamma(q_1 \rightarrow q_2 \nu \bar{\nu}) = \frac{3\alpha^2 G_F^2}{12288\pi^3 \sin^4 \theta_W} m_1^5 R(y) \left| \sum_j \lambda_j f(x_j) \right|^2$$

$$x_j = m_j^2 / M_W^2, \quad y = m_2^2 / m_1^2, \quad \lambda_j = V_{j2}^* V_{j1}$$

$$R(y) = 1 - 8y + 8y^3 - y^4 - 12y^2 \ln y$$

$$f(x) = f_B(x) + f_Z(x)$$

$$f_B(x) = \frac{4x}{(x-1)^2} (x-1 - \ln x)$$

$$f_Z(x) = \frac{x}{2(x-1)^2} [(-6+x)(x-1) + (2+3x) \ln x]$$

**Charged Lepton Pair Emission**  $q_1 \rightarrow q_2 \ell^+ \ell^-$

$$\Gamma(q_1 \rightarrow q_2 \ell^+ \ell^-) = \frac{\alpha^2 G_F^2}{192 \pi^3} m_1^5 \left\{ (H_1 + H_2) \left| \sum_j \lambda_j f^V(x_j) \right|^2 \right. \\ \left. + (H_1 - H_2) \left| \sum_j \lambda_j f^A(x_j) \right|^2 + H_3 \left| \sum_j \lambda_j f^T(x_j) \right|^2 \right. \\ \left. + H_4 \operatorname{Re} \left( \sum_j \lambda_j f^V(x_j) \right) \left( \sum_j \lambda_j f^T(x_j) \right)^* \right\}$$

Dynamical form factors:

$$f^V(x) = -\frac{1}{32 \pi \sin^2 \theta_W} f_B(x) + \frac{4 \sin^2 \theta_W - 1}{8 \pi \sin^2 \theta_W} f_Z(x) + \frac{1}{\pi} f_{\gamma 1}(x)$$

$$f^A(x) = \frac{1}{32 \pi \sin^2 \theta_W} f_B(x) + \frac{1}{8 \pi \sin^2 \theta_W} f_Z(x)$$

$$f^T(x) = \frac{1}{\pi} f_{\gamma 2}(x)$$

$$f_{\gamma 1}(x) = \frac{1}{72(1-x)^4} \left[ x^2(-25+19x)(1-x) - 2(8-32x+54x^2-30x^3+3x^4) \ln x \right]$$

$$f_{\gamma 2}(x) = \frac{x}{24(1-x)^4} \left[ (-7+5x+8x^2)(x-1) + 6x(2-3x) \ln x \right]$$

***Kinematic Factors:***

$$H_1 = m_1^{-8} \{-2I(2) + (\Sigma + 2m_\ell^2)I(1) + (D^2 - 4m_\ell^2 \Sigma)I(0) + 2m_\ell^2 D^2 I(-1)\}$$

$$H_2 = 6m_1^{-8} m_\ell^2 [\Sigma I(0) - I(1)]$$

$$H_3 = m_1^{-8} \{-\Sigma I(1) - [\Sigma(\Sigma + 2m_\ell^2) + 12m_1^2 m_2^2]I(0) \\ + [2\Sigma^2 (\Sigma - m_\ell^2) - 24m_\ell^2 m_1^2 m_2^2]I(-1) + 4m_\ell^2 \Sigma D^2 I(-2)\}$$

$$H_4 = 6m_1^{-8} \{\Sigma I(1) + [-D^2 + 2m_\ell^2 \Sigma]I(0) - 2m_\ell^2 D^2 I(-1)\}$$

$$\Sigma = m_1^2 + m_2^2$$

$$D = m_1^2 - m_2^2$$

$$I(n) = \int_{s_-}^{s_+} ds s^n \sqrt{(1 - 4m_\ell^2 / s)(s^2 - 2\Sigma s + D^2)}$$

$$s_+ = (m_1 - m_2)^2$$

$$s_- = 4m_\ell^2$$

**CKM mixing matrix**

$$V_{ud} = 0.97418 \pm 0.00027$$

$$V_{us} = 0.22552 \pm 0.0019$$

$$V_{ub} = 0.00393 \pm 0.00036$$

$$V_{cd} = 0.230 \pm 0.011$$

$$V_{cs} = 0.9723 \pm 0.0026$$

$$V_{cb} = 0.0412 \pm 0.0011$$

$$V_{td} = 0.0081 \pm 0.0006$$

$$V_{ts} = 0.0387 \pm 0.0023$$

$$V_{tb} = 0.999134 \pm 0.000045$$

## *Dynamical Form Factors*

### **Current quark masses**

$$m_u = 2.4 \pm 0.9 \text{ MeV}$$

$$m_c = 1.27^{+0.07}_{-0.11} \text{ GeV}$$

$$m_t = 171.2 \pm 2.1 \text{ GeV}$$



**Neutrino Pair Production**

$m_j$	$x_j$	$f(x_j)$
0.2400E-02	0.89110480E-09	0.54816915E-07
0.1270E+01	0.24952480E-03	0.59628445E-02
0.1712E+03	0.45343361E+01	0.62775693E+01

**Lepton Pair Production**

$m_j$	$f^V(x_j)$	$f^A(x_j)$	$f^T(x_j)$
0.2400E-02	0.14740260E+01	0.31787867E-09	-0.82731128E-10
0.1270E+01	0.58651859E+00	0.88889967E-04	-0.23267408E-04
0.1712E+03	-0.27343687E+00	0.72681391E+00	-0.11448136E+00

*Bound Quarks in Hadronic States*

- ☯ Consider  $M_1$  changes into meson  $M_2$  electroweak process
- ☯ Bound  $q_1$  in  $M_1$  changes into bound  $q_2$  in  $M_2$
- ☯ Bound quarks in hadrons behave like free particles
- ☯ Both  $q_1$  and  $q_2$  treated as free quarks
- ☯ All QCD effects adequately taken care of by assigning effective masses to the quarks
- ☯ Empirical determination of effective quark masses
- ☯ Effective quark masses determined from weak decay processes

*Processes considered*

$$K^- \rightarrow \pi^- \nu \bar{\nu}$$

$$B^- \rightarrow \pi^- \nu \bar{\nu}$$

$$B^- \rightarrow K^- \nu \bar{\nu}$$

$$K^- \rightarrow \pi^- e^+ e^-$$

$$B^- \rightarrow \pi^- e^+ e^-$$

$$B^- \rightarrow K^- e^+ e^-$$

*Effective quark masses*

- ☯  $d$ -quark mass in  $\pi^-$  is taken to be negligible.
- ☯  $s$ -quark mass in  $K^- = 339.2 \text{ MeV}$
- ☯  $b$ -quark mass in  $B^- = 2779 \text{ MeV}$

*Kinematic Factors*☯ **R Factor**

<b>Transition</b>	<b><math>m_1</math> (MeV)</b>	<b><math>m_2</math> (MeV)</b>	<b>R</b>
$K^- \rightarrow \pi^-$	<b>339.2</b>	<b>0</b>	<b>1.0</b>
$B^- \rightarrow \pi^-$	<b>2779</b>	<b>0</b>	<b>1.0</b>
$B^- \rightarrow K^-$	<b>2779</b>	<b>339.2</b>	<b>0.892045</b>

\*\*\*\*\* H FACTORS \*\*\*\*\*

M1= 339.2 M2= 1.0E-10

H(1)	H(2)	H(3)	H(4)
0.500000D+00	0.000000D+00	-0.144006D+02	-0.200000D+01

M1= 2229.0 M2= 1.0E-10

H(1)	H(2)	H(3)	H(4)
0.500000D+00	0.000000D+00	-0.144006D+02	-0.200000D+01

M1= 2229.0 M2= 339.2

H(1)	H(2)	H(3)	H(4)
0.419536D+00	0.000000D+00	-0.155275D+02	-0.167814D+01

Magnitudes of  $\lambda_j f(x_j)$  for the decay  $q_1 \rightarrow q_2 \nu \bar{\nu}$

Transition	$ \lambda_u f(x_u) $	$ \lambda_c f(x_c) $	$ \lambda_t f(x_t) $
$s \rightarrow d$	1.2043E-8	1.3333E-3	1.9680E-3
$b \rightarrow d$	1.5979E-10	5.6504E-5	5.2098E-2
$b \rightarrow s$	4.8584E-11	2.3887E-4	2.4275E-1



Magnitudes of  $\lambda_j f^\alpha(x_j)$ ,  $\alpha = V, A, T$  for the decay  $q_1 \rightarrow q_2 \ell^+ \ell^-$

Transition	$j$	$ \lambda_j f^V(x_j) $	$ \lambda_j f^A(x_j) $	$ \lambda_j f^T(x_j) $
$s \rightarrow d$	$u$	3.2384E-1	6.9838E-11	1.8176E-11
	$c$	1.3115E-1	1.9876E-5	5.2026E-6
	$t$	8.57225E-5	2.2786E-4	3.5890E-5
$b \rightarrow d$	$u$	4.2968E-3	9.2662E-13	2.4116E-13
	$c$	5.5579E-3	8.4232E-7	2.2048E-7
	$t$	2.2693E-3	6.0318E-3	9.5008E-4
$b \rightarrow s$	$u$	1.3064E-3	2.8174E-13	7.3325E-14
	$c$	2.3496E-2	3.5609E-6	9.3209E-7
	$t$	1.0574E-2	2.8106E-2	4.4270E-3

## Results

### *Neutrino Pair Emission*

$$\underline{K^- \rightarrow \pi^- \nu \bar{\nu}}$$

- ☉ Dominant contributions: internal  $c$ - and  $t$ -quarks.
- ☉ Because relative phases of CKM matrix elements are not known, only a range for the estimated branching ratio is obtained:

$$BR_{(th)}(K^- \rightarrow \pi^- \nu \bar{\nu}) = (0.39 - 10.55) \times 10^{-10}$$

- ☉ Compared to experimental value of

$$BR_{(exp)}(K^- \rightarrow \pi^- \nu \bar{\nu}) = (1.5_{-0.9}^{+1.3}) \times 10^{-10}$$

$$\underline{B^- \rightarrow \pi^- \nu \bar{\nu}, \quad B^- \rightarrow K^- \nu \bar{\nu}}$$

☯ Internal  $t$ -quark contribution dominates.

$$BR_{(th)}(B^- \rightarrow \pi^- \nu \bar{\nu}) = 1.283 \times 10^{-6}$$

$$BR_{(th)}(B^- \rightarrow K^- \nu \bar{\nu}) = 2.485 \times 10^{-5}$$

☯ Experimental upper bounds:

$$BR_{(exp)}(B^- \rightarrow \pi^- \nu \bar{\nu}) < 1.0 \times 10^{-4}$$

$$BR_{(exp)}(B^- \rightarrow K^- \nu \bar{\nu}) < 1.4 \times 10^{-5}$$

### *Charged Lepton Pair Emission*

Analysis considerably more involved.

$$\underline{K^- \rightarrow \pi^- e^+ e^-}$$

Dominant contributions:

$$\lambda_u f^V(x_u) \text{ and } \lambda_c f^V(x_c)$$

$$\underline{B^- \rightarrow \pi^- e^+ e^-}, \underline{B^- \rightarrow K^- e^+ e^-}$$

Dominant contributions:

$$\lambda_u f^V(x_u), \lambda_c f^V(x_c), \lambda_t f^V(x_t),$$

$$\lambda_t f^A(x_t) \text{ and } \lambda_t f^T(x_t)$$

Charged lepton pair emissions

Decay Mode	$BR_{(th)}$	$BR_{(exp)}$
$K^- \rightarrow \pi^- e^+ e^-$	1.954E-6	$(2.88 \pm 0.13)E-7$
$B^- \rightarrow \pi^- e^+ e^-$	$(0.690 - 1.362)E-8$	$< 1.8E-7$
$B^- \rightarrow K^- e^+ e^-$	$(1.186 - 1.266)E-7$	$(4.9 \pm 1.0)E-7$

## Conclusion

- ☯ A simple model to describe electroweak interactions of quarks inside hadrons.
- ☯ A quark tightly bound inside a hadron behaves as free particle with an effective mass.
- ☯ This simple model allows calculation of electroweak interaction processes at the quark level directly to hadronic processes.

- ☯ Experimental data on the charged weak decays of mesons are used to obtain estimates of effective masses of quarks in the hadrons.
- ☯ The model is applied to the rare electroweak decays of mesons.
- ☯ Reasonable agreement is obtained when comparing to the experimental values.