

Unusual Landau levels in biased bilayer Bernal graphene

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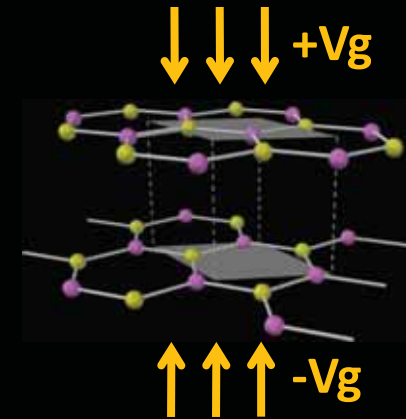
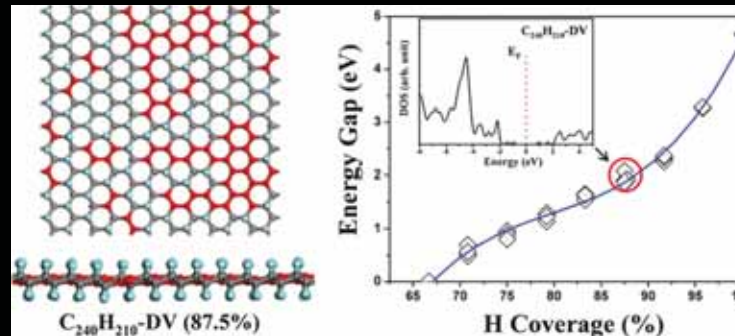


Outline ...

- **Motivation**
 - Why bilayer graphene?*
 - Why magneto physical properties?*
- **Comprehensive tight-binding method**
- **Landau levels and Landau wave functions**
 - Breaking of inter-valley degeneracy
 - Coupling of Landau levels*
- **Magneto-optical properties**
- **Summary**

Why bilayer graphene ?

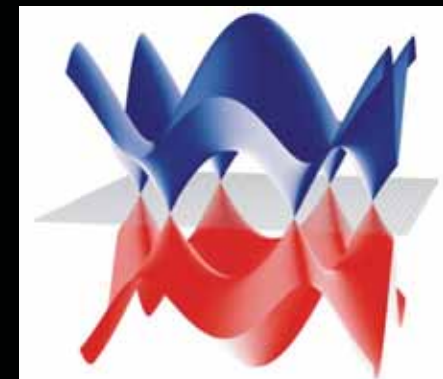
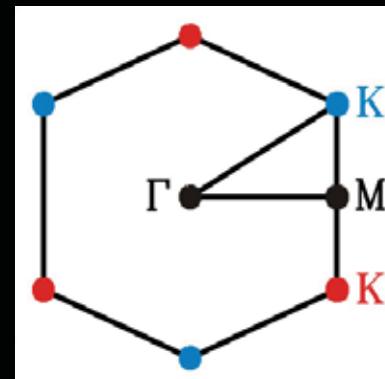
missing of energy gap → Adsorption of H atoms on the C dangling bonds
 Cutting G sheet into 1-D ribbons with finite width
 Applying a bias E field on top and bottom of G bilayer
 (doping level & energy gap)



Why magneto physical properties ?

2-D electronic states → discrete LLs:
 a better way to study electronic properties

Multiple degeneracy: e-h symmetry
 inter-valley symmetry



Atomic structure of bilayer graphene

$$\langle A_{1(2)} | H | B_{1(2)} \rangle \longrightarrow \gamma_0 \quad (-3.12 \text{ eV})$$

$$\langle A_{1(2)} | H | A_{2(1)} \rangle \longrightarrow \gamma_1 \quad (0.38 \text{ eV})$$

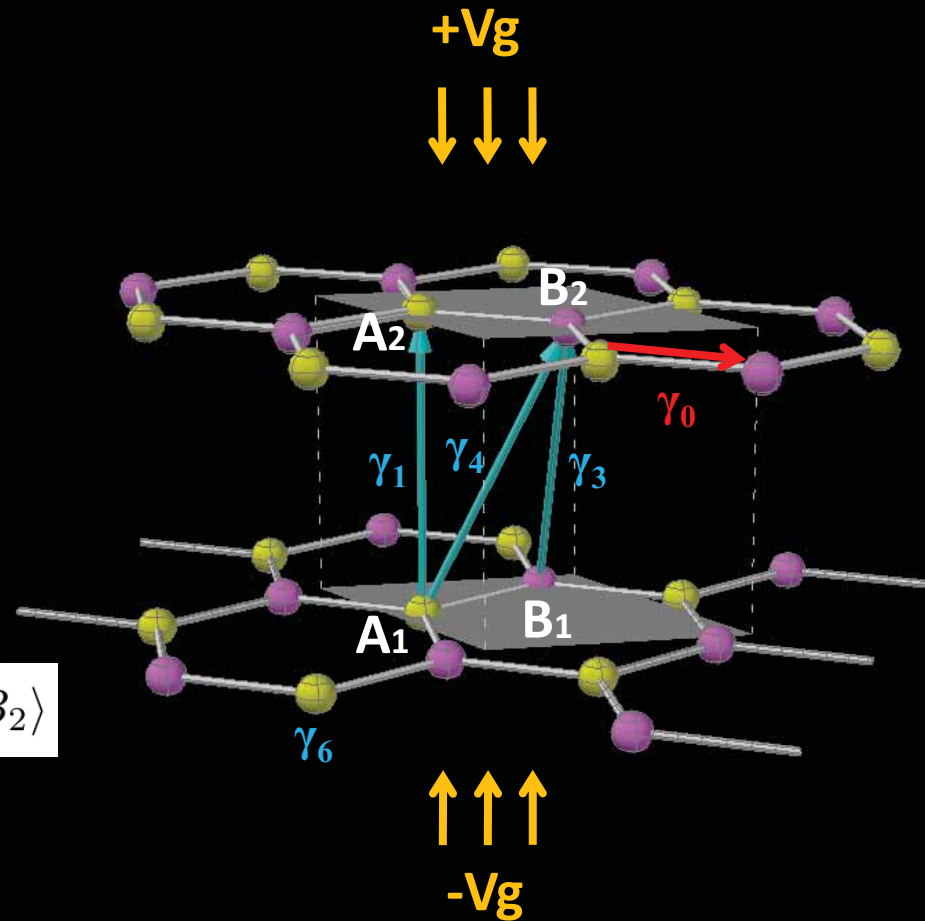
$$\langle B_{1(2)} | H | B_{2(1)} \rangle \longrightarrow \gamma_3 \quad (0.28 \text{ eV})$$

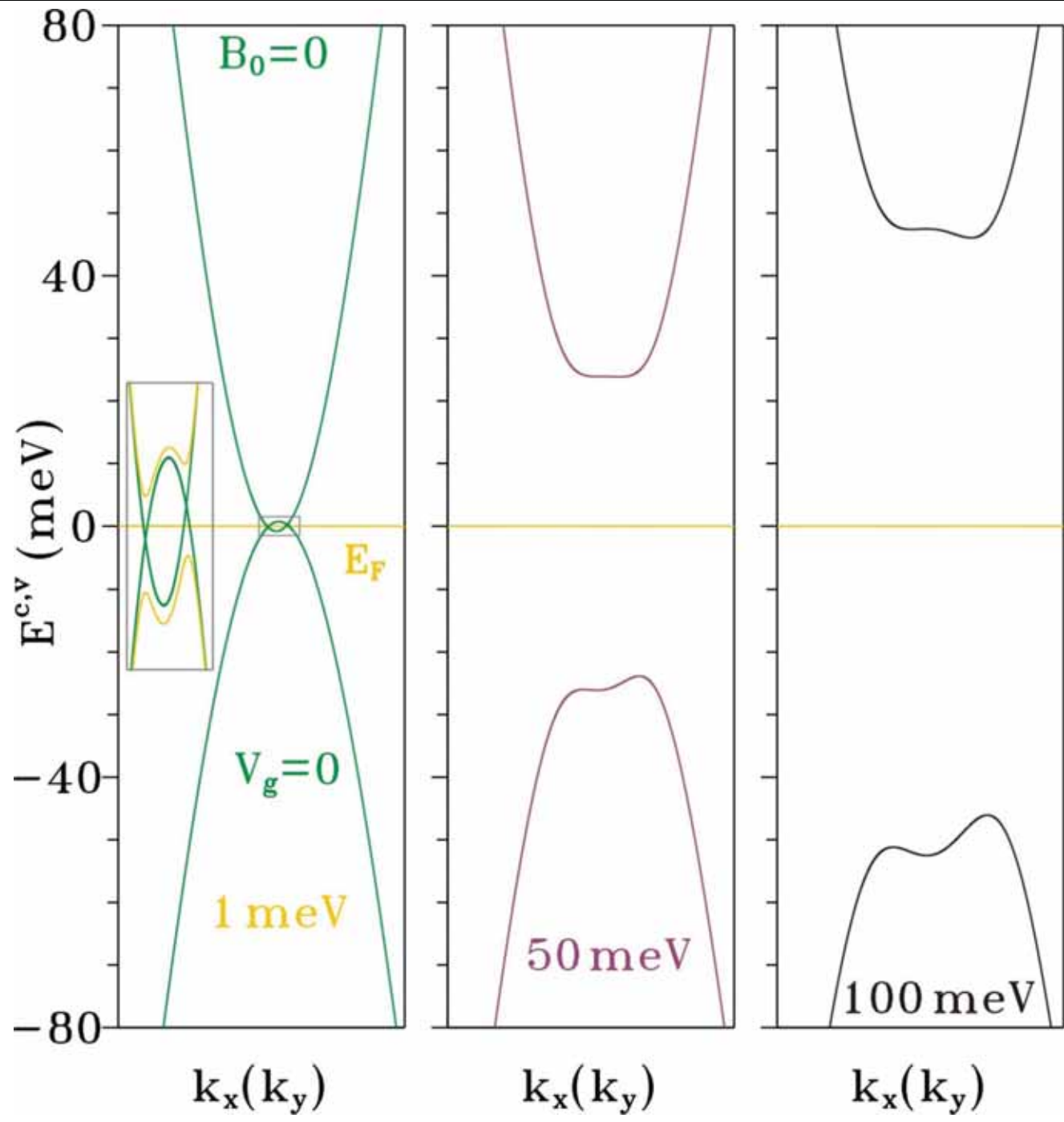
$$\langle A_{1(2)} | H | B_{2(1)} \rangle \longrightarrow \gamma_4 \quad (0.12 \text{ eV})$$

$$\langle A_{1(2)} | H | A_{1(2)} \rangle \longrightarrow \gamma_6 \quad (0.016 \text{ eV})$$

$$\Psi = c_a |A_1\rangle + c_{a'} |A_2\rangle + c_b |B_1\rangle + c_{b'} |B_2\rangle$$

$$H = \begin{pmatrix} H_{B_2B_2} & H_{B_2A_2} & H_{B_2A_1} & H_{B_2B_1} \\ H_{A_2B_2} & H_{A_2A_2} & H_{A_2A_1} & H_{A_2B_1} \\ H_{A_1B_2} & H_{A_1A_2} & H_{A_1A_1} & H_{A_1B_1} \\ H_{B_1B_2} & H_{B_1A_2} & H_{B_1A_1} & H_{B_1B_1} \end{pmatrix}$$





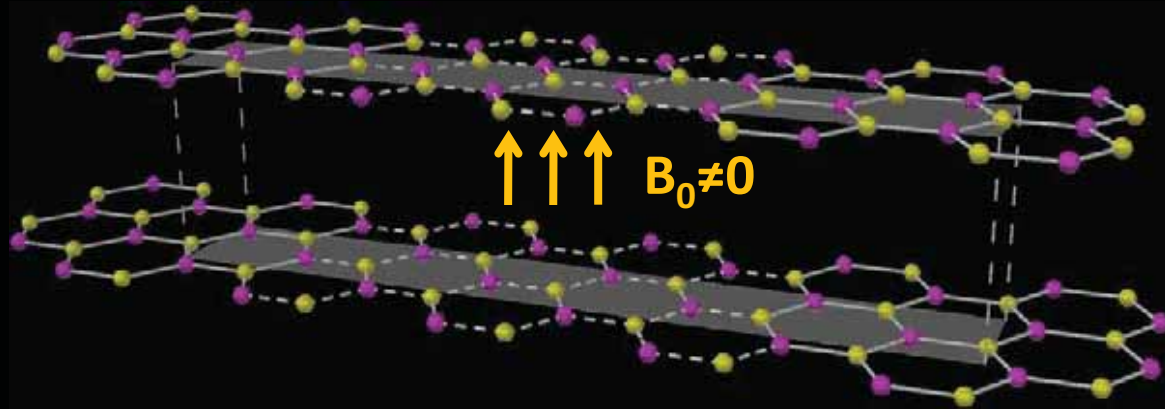
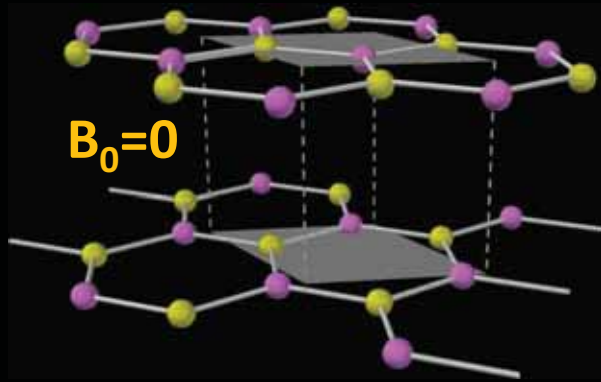
Energy bands:
 $B_0=0$
 $V_g \neq 0$

Mexican-hat shape

Apply a magnetic field perpendicular to the plane...

$$\text{Peierls phase } \Delta G(\vec{R}_{j_l}, \vec{R}_{k_{l'}}) = \int_0^1 (\vec{R}_{k_{l'}} - \vec{R}_{j_l}) \cdot \vec{A}[\vec{R}_{j_l} + \lambda(\vec{R}_{k_{l'}} - \vec{R}_{j_l})] d\lambda$$

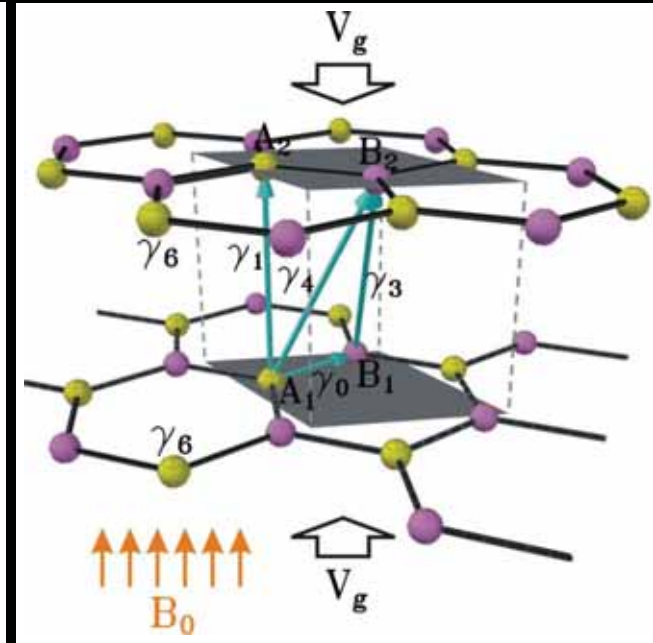
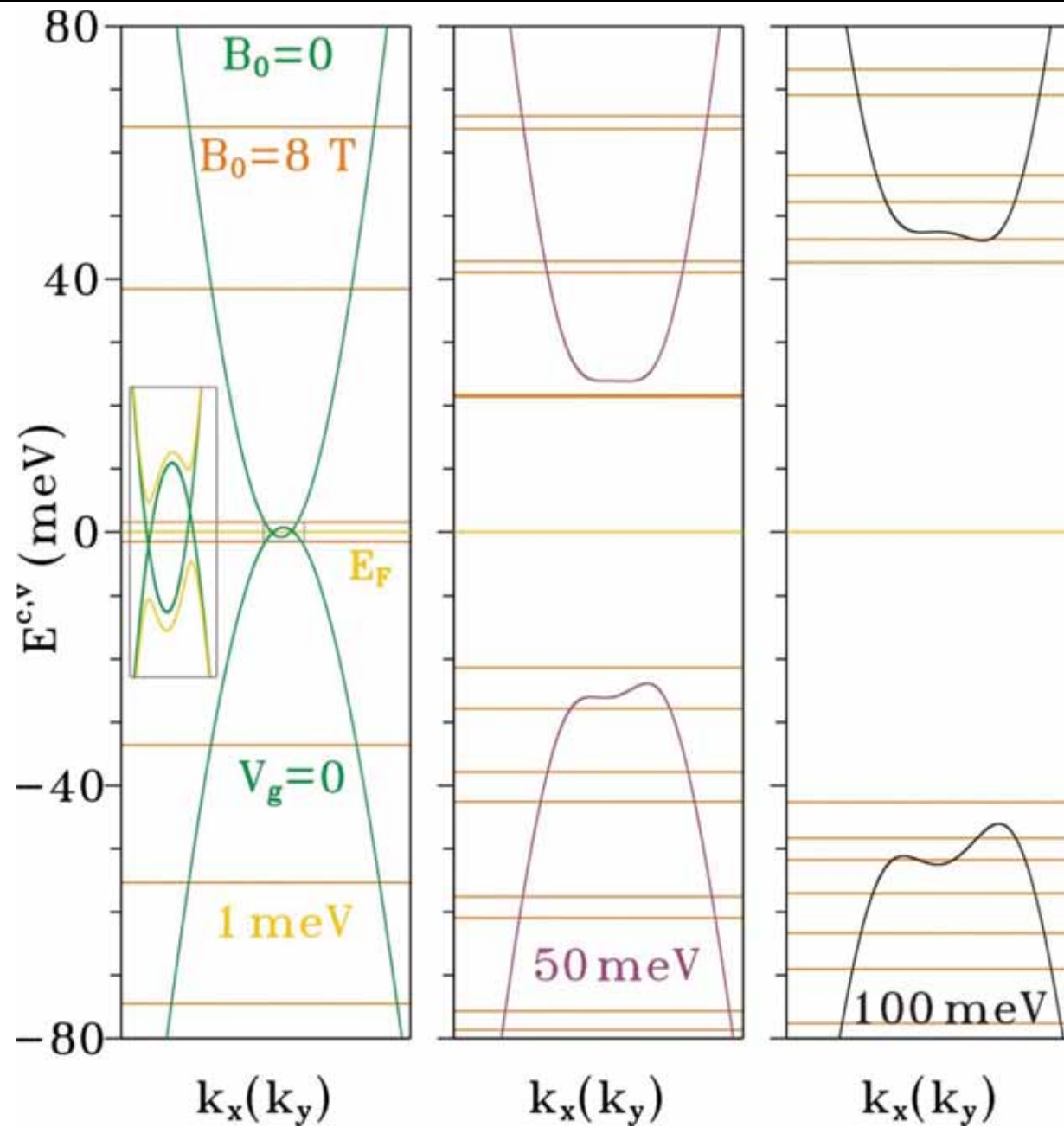
$$\langle k_{l'} | H | j_l \rangle = \gamma_s(\vec{R}_{j_l}, \vec{R}_{k_{l'}}) \sum \exp [i\vec{k} \cdot (\vec{R}_{j_l} - \vec{R}_{k_{l'}}) + i\frac{e}{\hbar} \Delta G(\vec{R}_{j_l}, \vec{R}_{k_{l'}})]$$



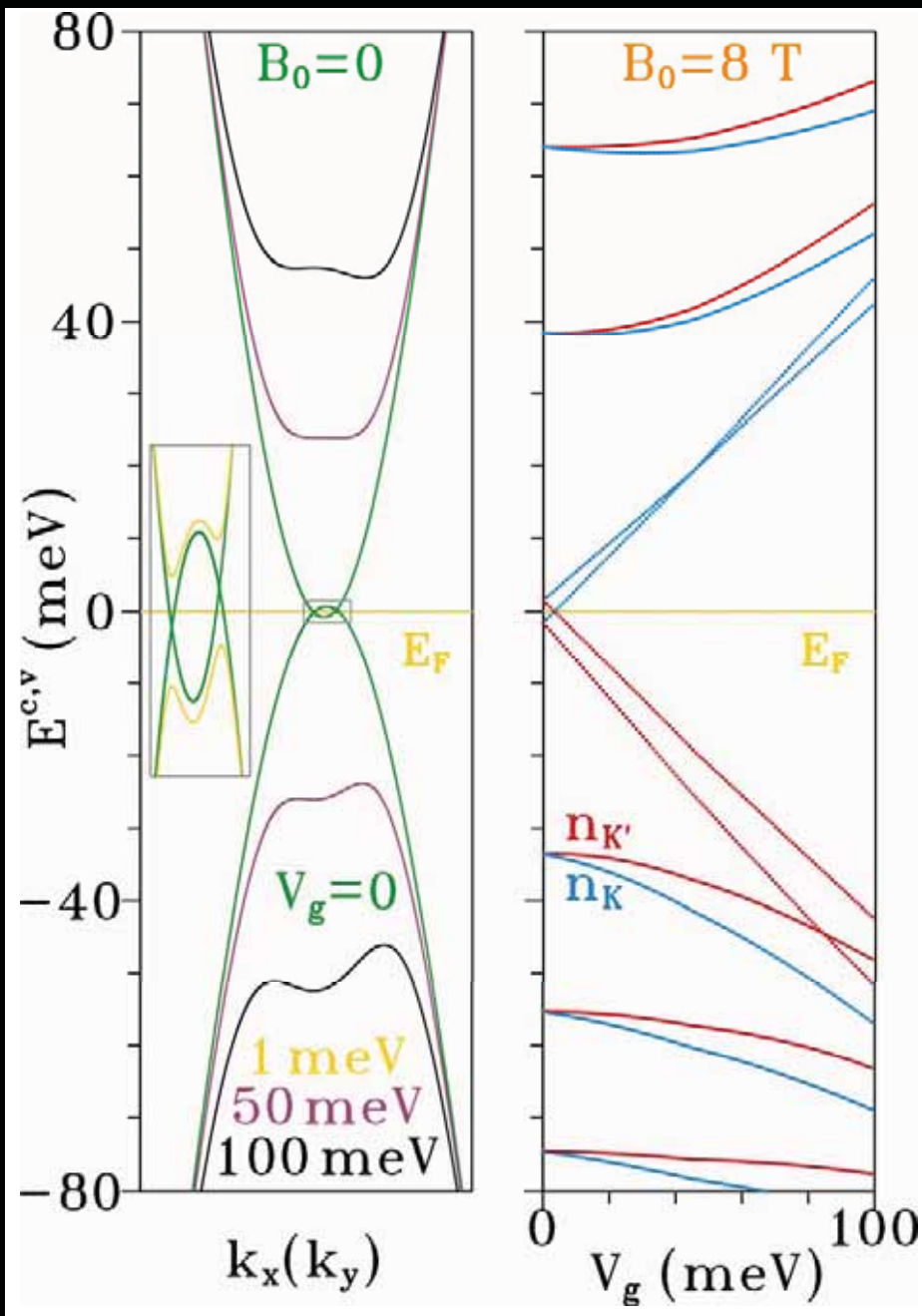
0 T: 4 atoms
8 T: 80,000 atoms

$$\Psi = \sum_{i=1}^{2R_B} c_{a,i} |A_{1,i}\rangle + c_{a',i} |A_{2,i}\rangle + c_{b,i} |B_{1,i}\rangle + c_{b',i} |B_{2,i}\rangle$$

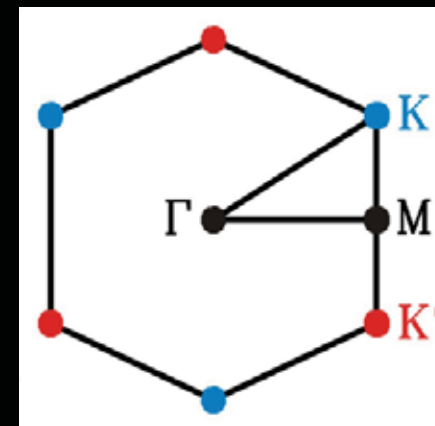
2-D Energy bands $B_0=0$ to Landau levels at $B_0 \neq 0$

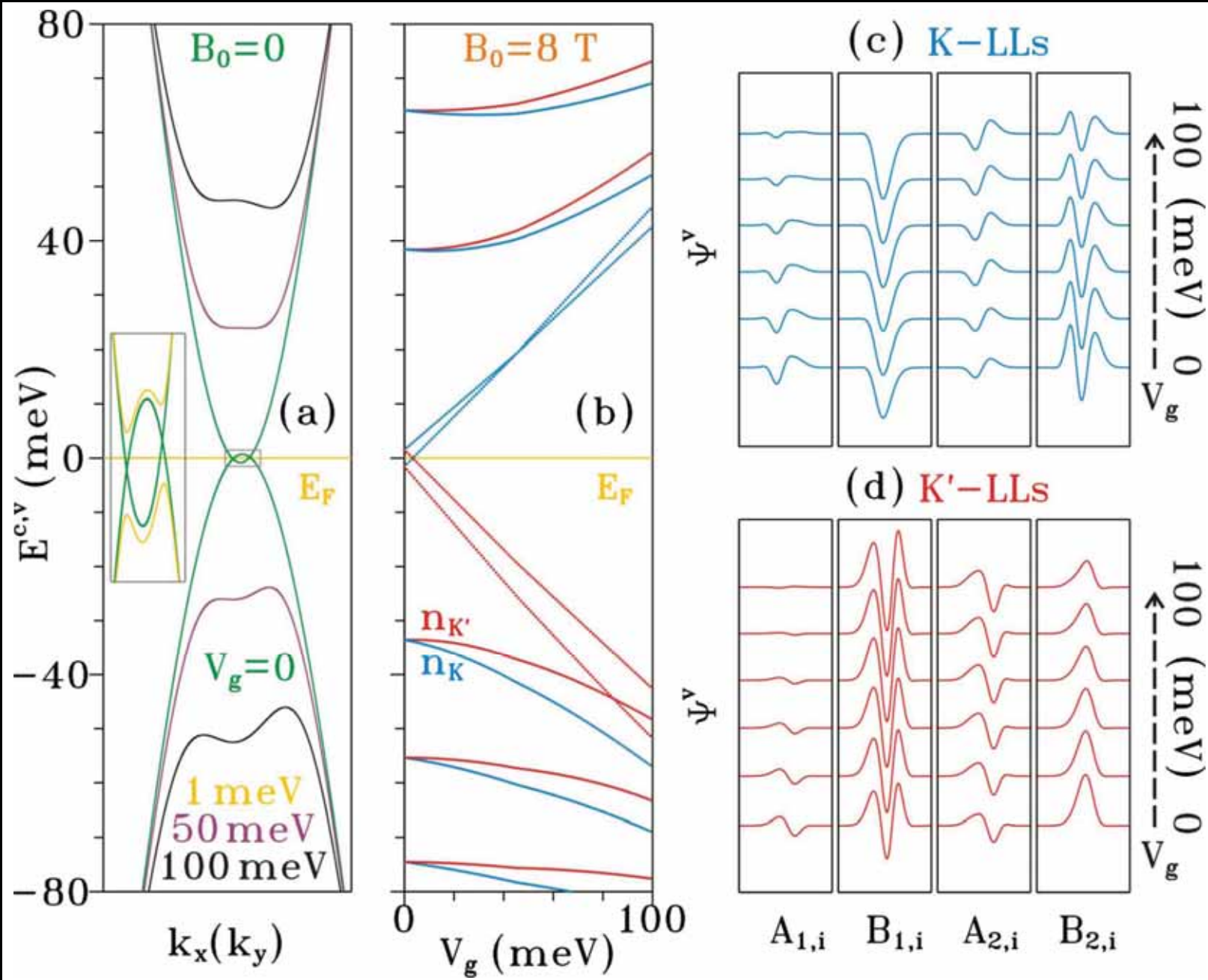


$B \neq 0$: Landau levels

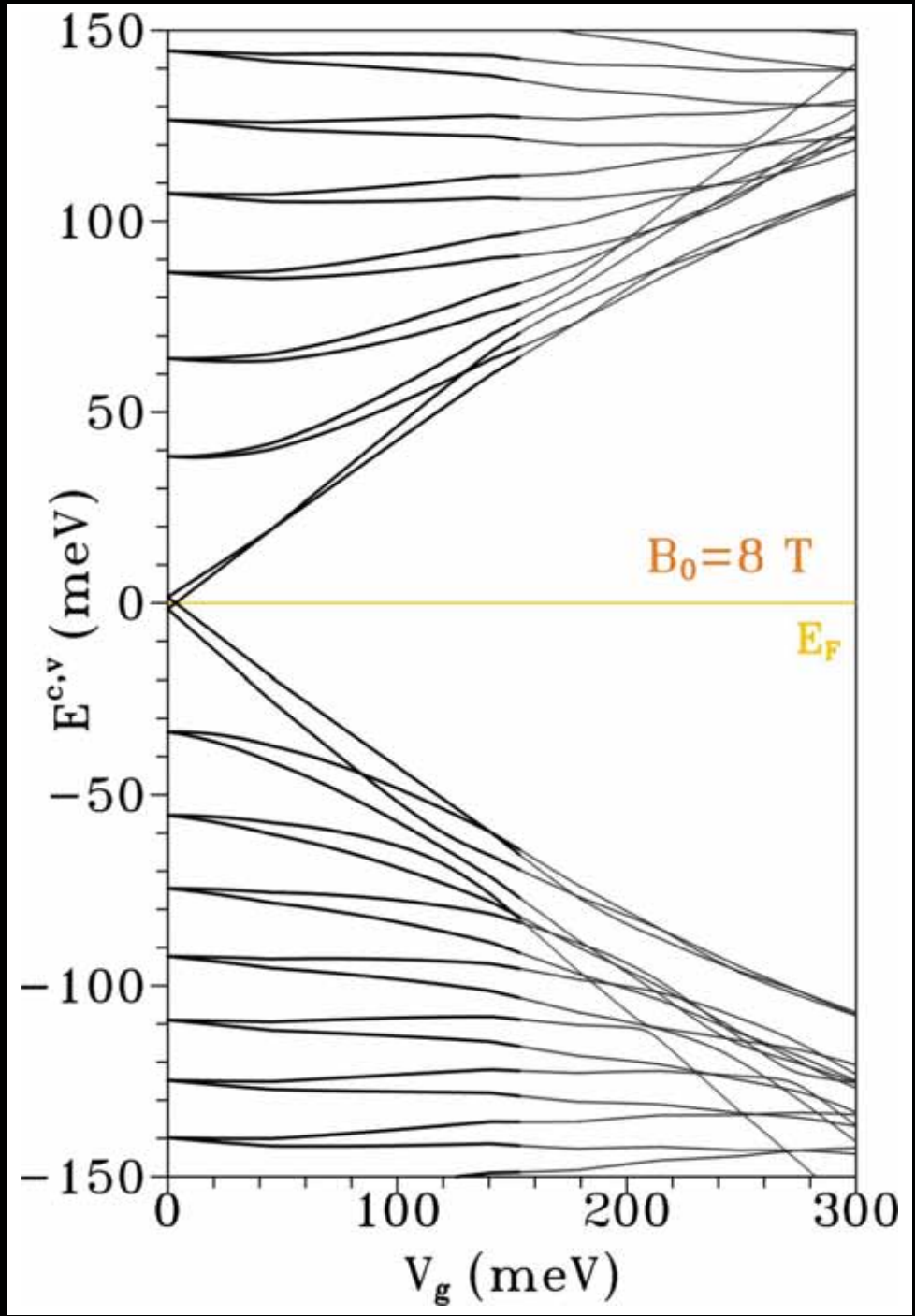


breaking of inter-valley degeneracy



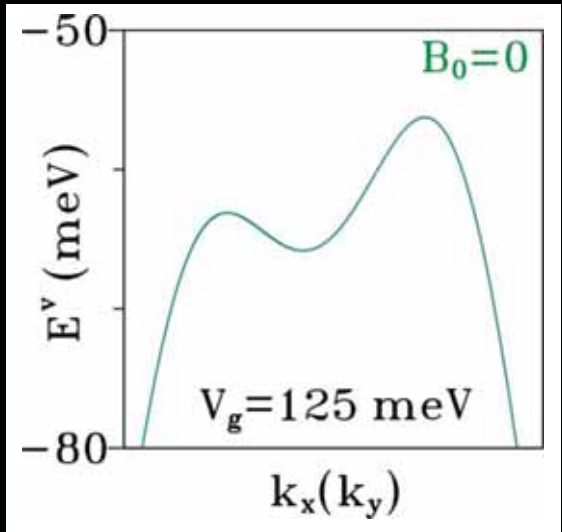


$$\langle \psi_{A_1}^* | -V_g | \psi_{A_1} \rangle + \langle \psi_{B_1}^* | -V_g | \psi_{B_1} \rangle + \langle \psi_{A_2}^* | V_g | \psi_{A_2} \rangle + \langle \psi_{B_2}^* | V_g | \psi_{B_2} \rangle$$



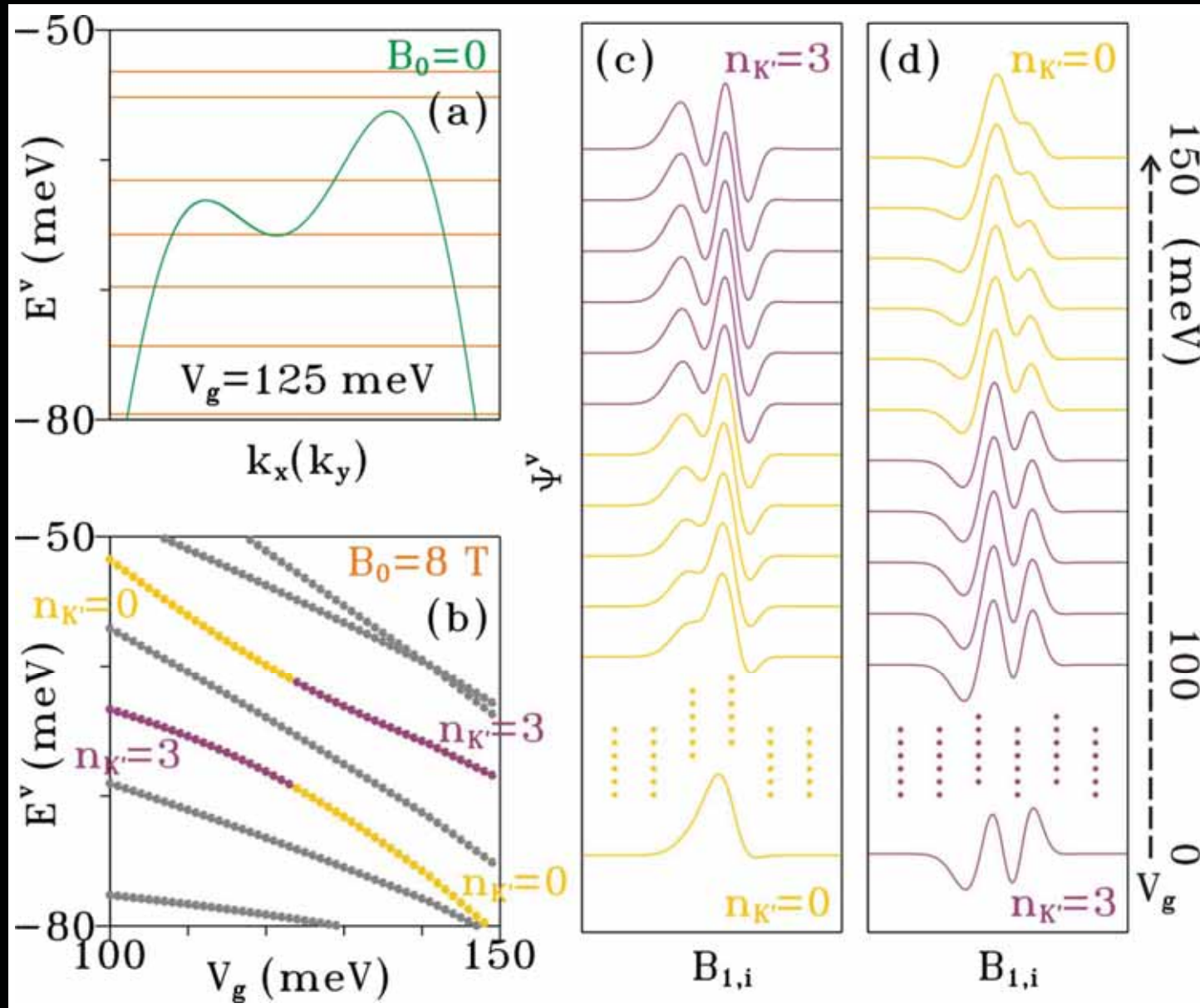
Landau levels for a wider range of energy & bias electric field

Lower LLs are more sensitive to V_g :
 Shift toward higher energy
 Splitting of K & K' states



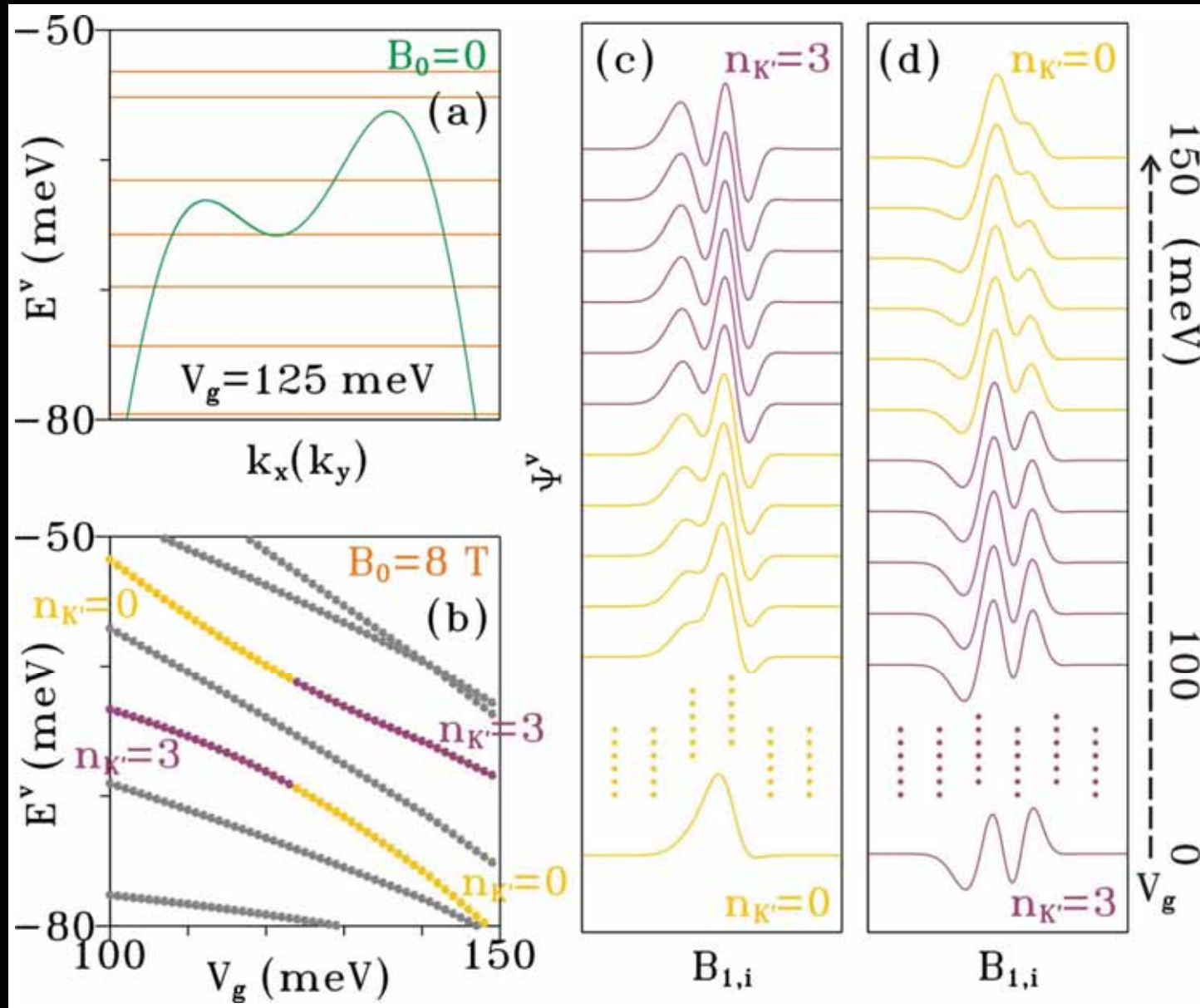
direct crossing
 avoided crossing

Landau-level coupling → mixing of Landau wave functions



Even though WFs are severely mixed, but in energy, there're still well-quantized LLs

Landau-level coupling \rightarrow mixing of Landau wave functions



$$\Delta n = 3I$$

interlayer coupling γ_3

Optical absorption spectrum...

absorption rate: $W(\vec{q} \sim 0, \omega) = \sum_{i,j} P_{i \rightarrow j} - P_{j \rightarrow i}$

$$= \frac{2\pi}{\hbar} \left(\frac{eA_0}{mc}\right)^2 2 \sum_{i,j} |\langle \Phi_j | e^{i\vec{q} \cdot \vec{r}} \hat{e} \cdot \vec{p} | \Phi_i \rangle|^2 \delta(E_j - E_i + \hbar\omega) [f(E_i) - f(E_j)]$$

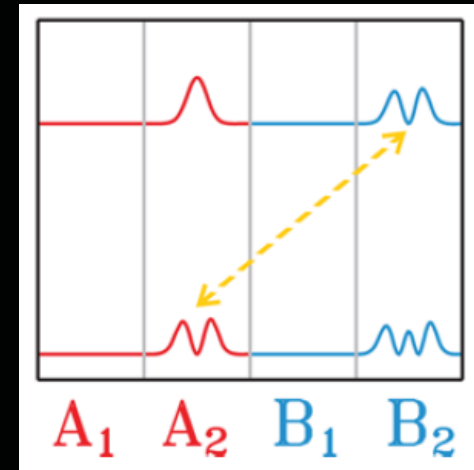
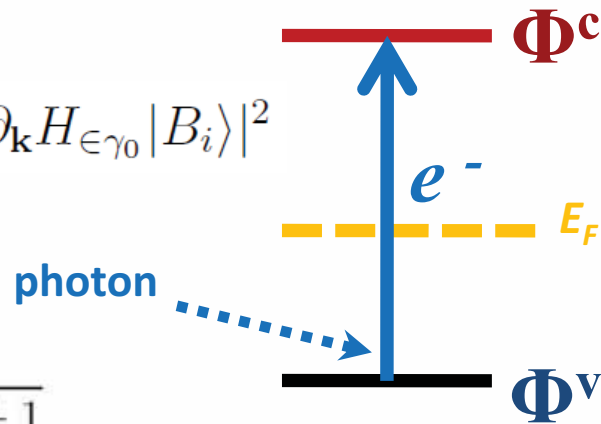
dipole matrix element within the gradient approximation

$$\begin{aligned} \langle \Phi_j | \vec{p} | \Phi_i \rangle &= \langle \Phi_j | \frac{im}{\hbar} [H, \vec{r}] | \Phi_i \rangle \\ &= \frac{im}{\hbar} (\vec{R}_i - \vec{R}_j) \langle \Phi_j | H | \Phi_i \rangle \\ &= \frac{m}{\hbar} \nabla_{\vec{k}} \langle \Phi_j | H | \Phi_j \rangle \end{aligned}$$

$$|\langle A_i | \partial_{\mathbf{k}} H_{\in \gamma_0} | B_i \rangle|^2$$

Fermi-Dirac distribution function:

$$f(E_{i,j}) = \frac{1}{\exp\{[E_{i,j} - \mu(T)]/k_B T\} + 1}$$

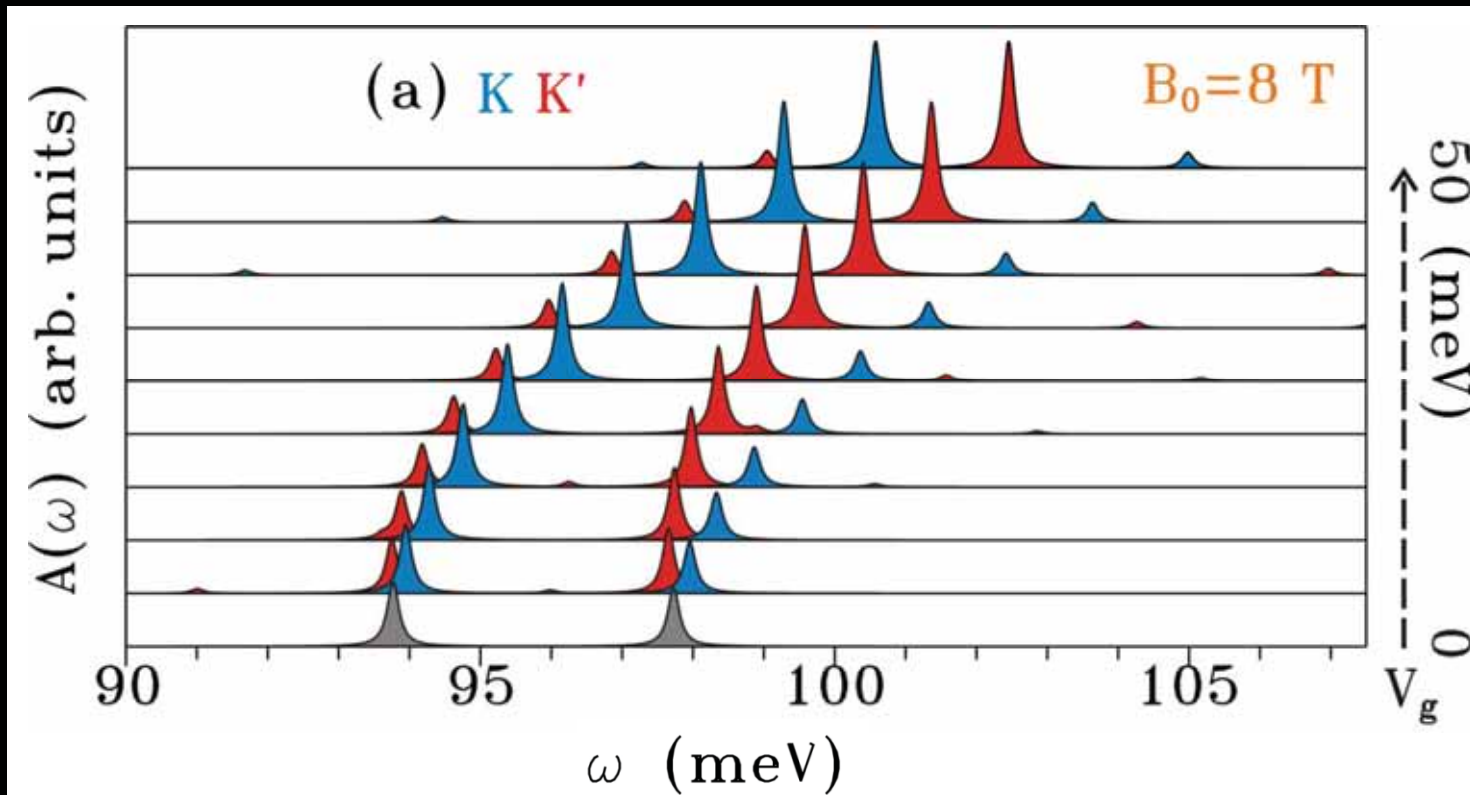


$\gamma_0 \gg \gamma_{1,3,4,6}$: Optical transitions are mainly determined by the terms related to γ_0 , as well as the WF on the both sides.

Sublattices must have the same node number...

to satisfy the orthogonality relation of Hermite polynomial.

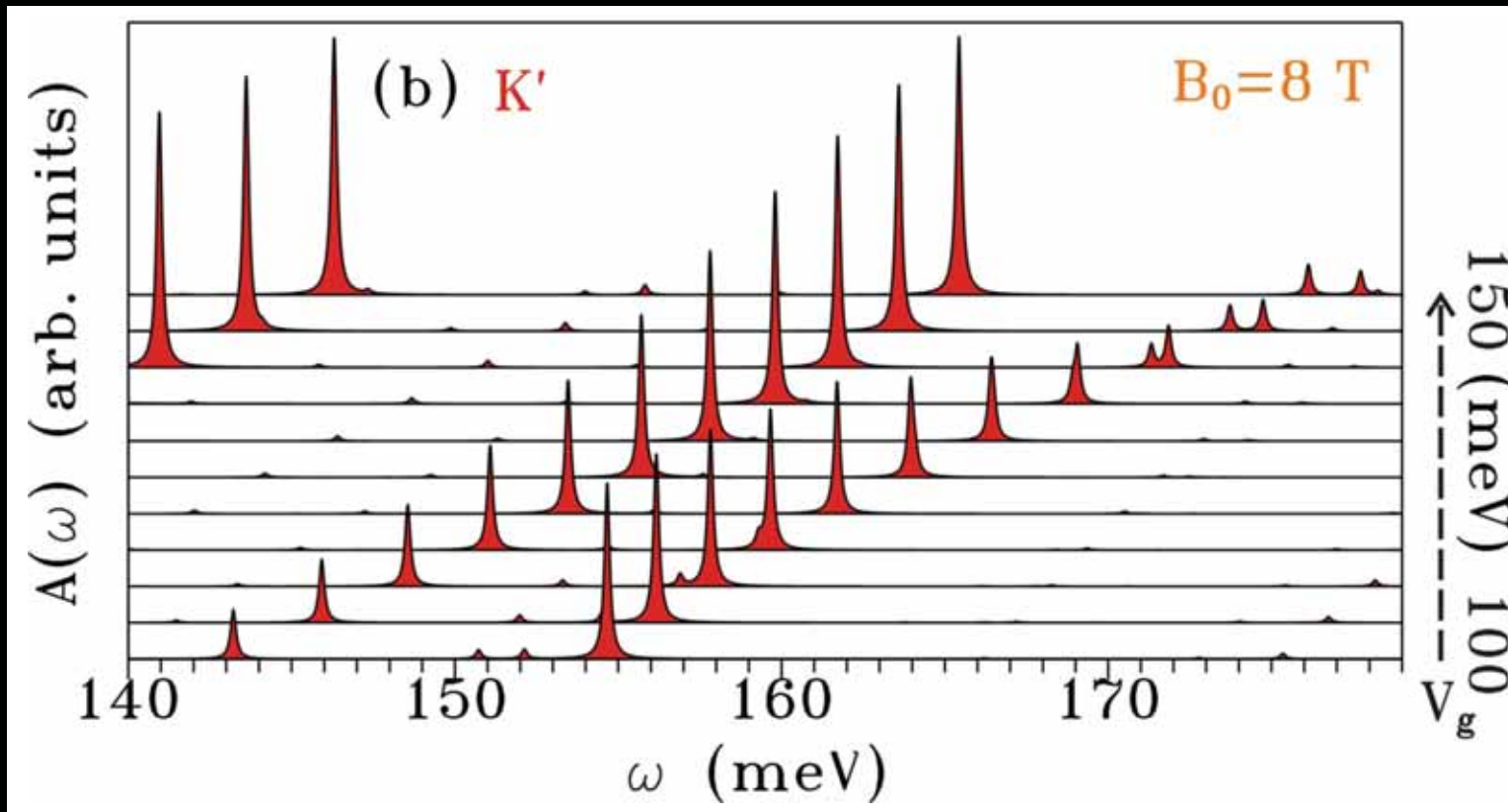
Magneto-optical spectra for breaking of inter-valley degeneracy



Electron-hole asymmetry: $n^v + 1 \rightarrow n^c$ $n^v \rightarrow n^c + 1$

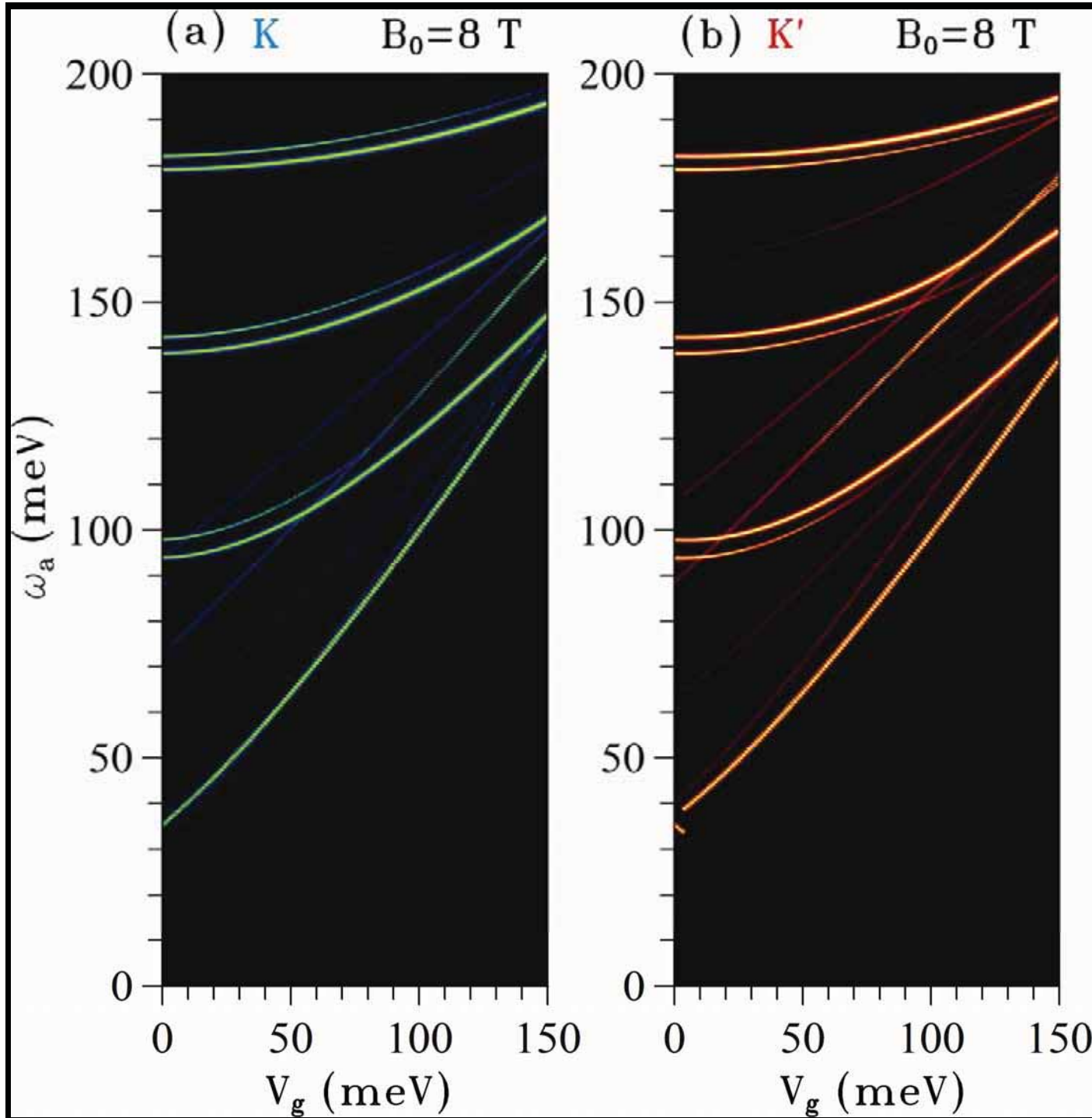
Lift inter-valley degeneracy: energy splitting and intensity changes

Magneto-optical spectra for Landau-level couplings



the changes in peak position & peak intensity, the occurrence of minor peaks.

Selection rules & transition rates by Landau wave functions.



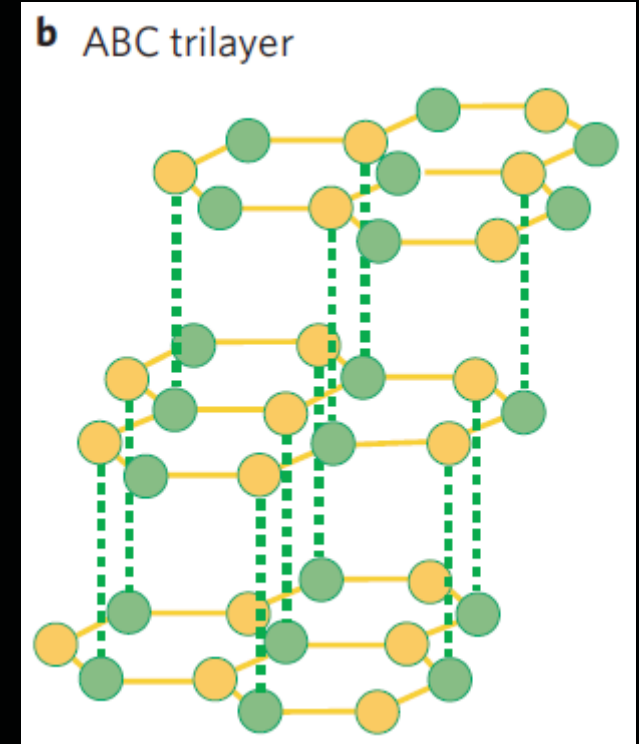
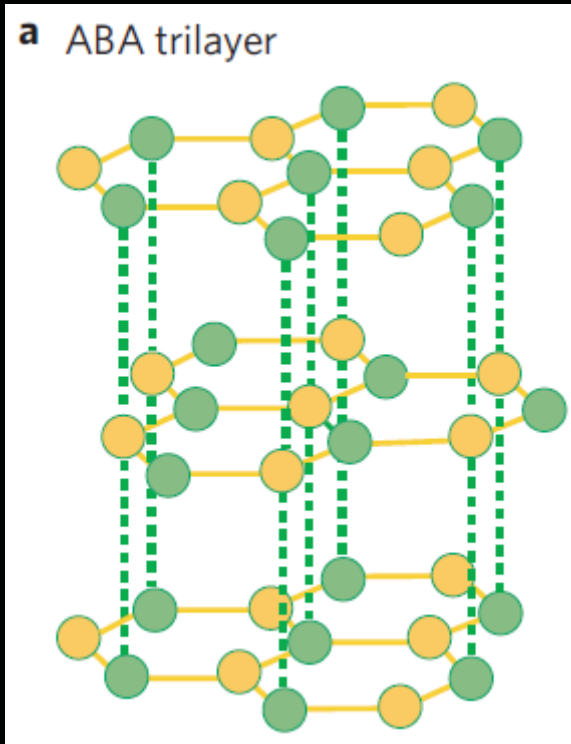
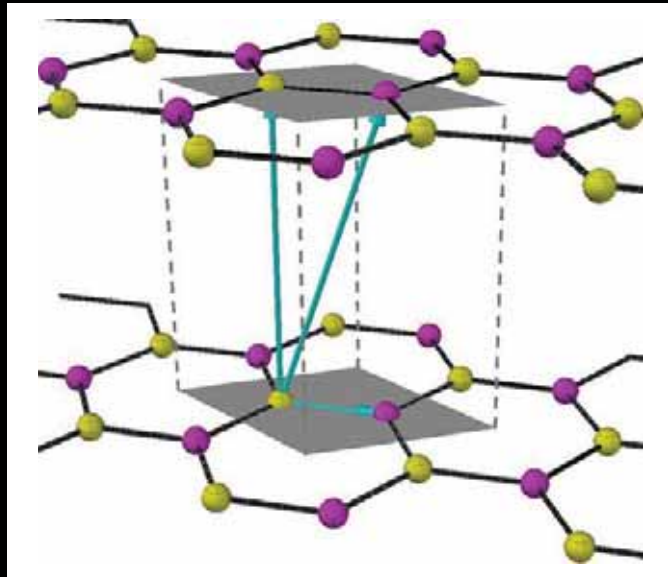
Contour plot of optical spectra: K & K' states

Breaking of inter-valley degeneracy

Landau level couplings

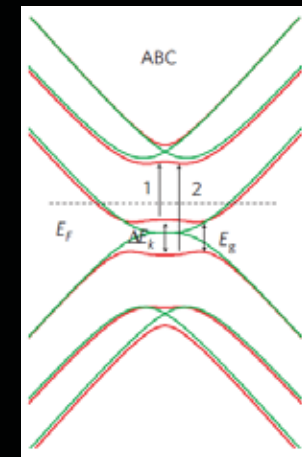
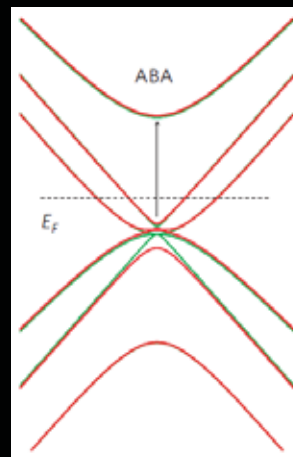
Peak intensities:
Changes in the WFs.

Whether the 2 valley states split and whether the LL couplings appear depend on the stacking order and the layer numbers.



Valley states split:
top & bottom must be asymmetric

LL couplings:
depends on stacking order



Summary ...

We use a complete tight-binding model to study the Landau level structures in biased bilayer graphene. The changes of weight in Landau wave function enable us to characterize the breaking of inter-valley degeneracy and the appearance of unusual Landau levels, which have their direct counterpart in magneto-optical spectra.

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