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Band-structure of graphene above the vacuum level

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... motivation



In the case of a potential well in z direction and a uniform potential in xyplane, two perpendicular motions decouple.

Therefore, states bound in z direction, may have arbitrarily large energy due to the

lateral motion of an electron.



Does this change and how if a laterally periodic potential is applied ?

Obviously, the in-plane and the perpendicular motions do couple in the presence of a periodic lateral potential.

We consider a model system with the potential

$$V(\mathbf{r}) = V(\mathbf{r}_{\parallel}) \times \delta(z), \tag{1}$$

where $V(\mathbf{r}_{\parallel})$ is a 2D periodic function. By Fourier transform

$$V(\mathbf{r}_{\parallel}) = \sum_{\mathbf{G}} V_{\mathbf{G}} e^{i\mathbf{G}\cdot\mathbf{r}_{\parallel}}.$$
 (2)

 V_0 is chosen negative to ensure the existence of a bound-state in *z* direction.

We are solving the single-particle Schrödinger equation

$$\left[-\frac{1}{2}\Delta + V(\mathbf{r}_{\parallel})\delta(z)\right]\psi(z,\mathbf{r}_{\parallel}) = E\psi(z,\mathbf{r}_{\parallel}).$$
(3)

Solutions localized near and continuous across the z = 0 plane are

$$\psi(z, \mathbf{r}_{\parallel}) = \sum_{\mathbf{G}} a_{\mathbf{G}} e^{-\sqrt{(\mathbf{G} + \mathbf{k})^2 - 2E} |z|} e^{i(\mathbf{G} + \mathbf{k}) \cdot \mathbf{r}_{\parallel}}$$
(4)

where $a_{\mathbf{G}}$ satisfy the system of linear equations

$$\sum_{\mathbf{G}'} V(\mathbf{G} - \mathbf{G}') a_{\mathbf{G}'} = -\sqrt{(\mathbf{G} + \mathbf{k})^2 - 2E} a_{\mathbf{G}}.$$
 (5)

Eqs. (5) constitute a non-linear eigenvalue problem with respect to the eigenvalue E.

A little of philosophy

Stand-alone vs. repeated geometry





Stand-alone 2D-periodic layer.

Repeated geometry representation.

Bound-states above the vacuum level are resolved inside the continuous spectrum. Price: Non-linear eigenvalue problem. Continuous spectrum and bound-states are intermixed above the vacuum level. Conventional (linear) eigenvalue problem.

$$\sum_{\mathbf{G}'} V(\mathbf{G} - \mathbf{G}') a_{\mathbf{G}'} = -\sqrt{(\mathbf{G} + \mathbf{k})^2 - 2E} a_{\mathbf{G}}.$$
 (6)

First, if $E > (\mathbf{G} + \mathbf{k})^2/2$ with an arbitrary \mathbf{G} then the eigenvalue problem involves complex matrix elements. We, however, prove the following

Lemma

The problem (6) admits no resonance solutions, i.e., all the eigenvalues E are real.

Proof.

$$\sum_{\mathbf{GG}'} a_{\mathbf{G}'}^* V(\mathbf{G} - \mathbf{G}') a_{\mathbf{G}'} = -\sum_{\mathbf{G}} \sqrt{(\mathbf{G} + \mathbf{k})^2 - 2E} |a_{\mathbf{G}}|^2.$$
(7)

We are numerically solving the non-linear eigenvalue problem

$$\sum_{\mathbf{G}'} B_{\mathbf{G}\mathbf{G}'}(E) a_{\mathbf{G}'} = a_{\mathbf{G}},\tag{8}$$

where

$$B_{\mathbf{GG}'}(E) = -\frac{V(\mathbf{G} - \mathbf{G}')}{\sqrt{(\mathbf{G} + \mathbf{k})^2 - 2E}}.$$
(9)



The eigenvalue of matrix B as a function of E.

Model band-structure



Model band-structure of a single periodic layer (red points) and layers in repeated-geometry (blue lines).

Projection on the ab initio graphene problem



z-dependence of the wave-functions of the states above the vacuum level obtained in the repeated geometry calculation.

- Our analysis shows that the graphene band-structure is affected drastically by the possibility for an electron above the vacuum level to escape from the layer;
- Whether or not at a particular k point a state above the vacuum level is bounded to the graphene layer is determined by the symmetry of that point.